

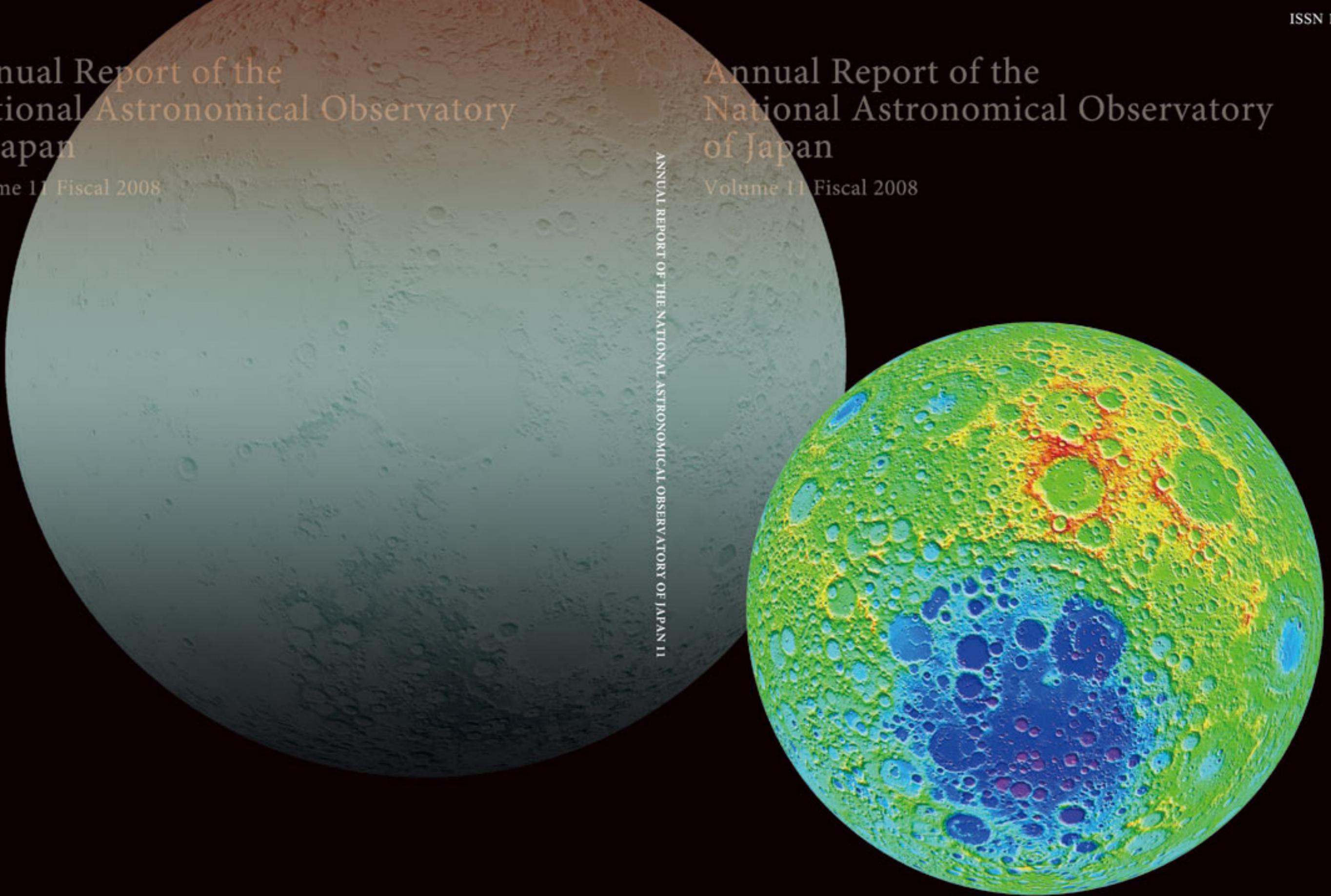
Annual Report of the  
National Astronomical Observatory  
of Japan

Volume 11 Fiscal 2008

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ANNUAL REPORT OF THE NATIONAL ASTRONOMICAL OBSERVATORY OF JAPAN 11



## Cover Caption

Lunar topographic map derived from KAGUYA laser altimeter LALT (NAOJ/JAXA/SELENE).  
South Pole-Aitken basin is spread over the lunar far side.

## Postscript

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# Annual Report of the National Astronomical Observatory of Japan

Fiscal **2008**

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**Annual Report of the  
National Astronomical Observatory of Japan  
Fiscal 2008**



**Shoken MIYAMA**  
**Director General of NAOJ**

## P R E F A C E

We are pleased to present the Annual Report of the National Astronomical Observatory of Japan (NAOJ) for fiscal 2008.

This year marked the 20th anniversary of the establishment of NAOJ, and we held a commemorative ceremony in November 2008. For the last 20 years, many observation facilities (such as Subaru Telescope) have been completed and they are now achieving remarkable results.

Over the last two decades, there were various findings in astronomy that brought us new insights and contributed to astronomical advancement. In 1995, researchers of Switzerland discovered a planet circling a star outside the solar system incidentally when observing a primary star. It turned out that this planet is orbiting around Pegasus 51 in an about four-day period and has a mass roughly a half of Jupiter's mass. Though the observed planetary system is quite different from our solar system, this discovery marked the beginning of similar observations which led to discoveries of more than 300 planets.

After this discovery, astronomers began to focus on three major targets:

- (1) Direct imaging of planets outside the solar system;
- (2) Exploration of the origin of diverse solar systems and reconstruction of conventional planet formation theory; and
- (3) Search for life in planets outside the solar system.

In connection with the target (1), a Canada-US

research team announced that they identified three potential planets around a primary star using the Gemini telescope and W.M. Keck Observatory in Hawaii, and found out that these potential planets are undoubtedly orbiting around the primary star by comparing two images taken this time and a few years ago. However, it seems to be difficult to estimate the mass of the potential planets from pictures and thus it is still controversial whether the mass of the potential planet is less than 13 times the Jupiter's mass, which is regarded as an essential qualification for a "planet" (the definition of planet defined by AIU does not apply to planets outside the solar system). Subaru Telescope will also start observation with advanced AO (Adaptive Optics) and HiCIAO (High Contrast Instrument for the Subaru next generation Adaptive Optics: a new stellar coronagraph) aiming for direct observation of planets outside the solar system.

For the target (2), researchers are actively working on theoretical studies to investigate the diversity of planets and their formation process. In theoretical simulation, more detailed calculation has become possible. The GRAPE system was a special-purpose computer system for astrophysical many-body simulations at the beginning, but it has evolved to a new general-purpose system called GRAPE-DR. It is expected that this new system will enable us to study planet formation under a wide range of initial conditions by performing calculation of N bodies around the Sun.

Meanwhile, initial conditions of planetary formation may be clearly defined by ALMA (Atacama Large Millimeter/submillimeter Array) under construction in Chile. ALMA is capable of analyzing details of a gas disk (primordial solar system disk) formed around a young star in a nearby star-forming region. By performing numerical simulation using the initial conditions derived from the structure of a gas disk, we may be able to figure out the formation process of planets in our solar system and the diversity of planetary systems.

It has been five years since the start of the ALMA Project, which will be completed over eight years. At the construction site in Chile, test observations have been conducted with antennas delivered from Japan, and the Japanese antennas proved their high performance and capability. As for other ALMA subsystems, we successfully developed a submillimeter receiver called "Band 10" covering the highest frequency band in ALMA, and confirmed that it meets extremely high requirements of ALMA. This achievement proved that the receiver developed by NAOJ has the world's highest capability in this frequency band.

The target (3) "Search for life in planets outside the solar system" is a medium-term objective in the future. Even if Subaru Telescope could obtain direct images of some planets outside the solar system with its advanced equipment, such images would be limited to those of gas-giant planets with the size of Jupiter because small rocky stars like the earth are difficult to be observed due to its darkness. On the other hand, indirect observation of habitable planets may become possible by introducing a high-resolution spectroscope. Generally speaking, in case of a Sun-like star, habitable planets (where a large amount of liquid water exists) are supposed to have orbits in a distance equivalent to the distance from the Sun to the Earth. If they are in such a distant region, current spectroscopes are unable to detect oscillations of the central star. In case of an M-type star (whose mass is less than a half of the Sun) with surface temperature much lower than that of the Sun, its habitable zone would be much closer to the central star. If there is an Earth-like planet within such a close habitable zone, oscillations of the central star would be larger due to closer distance and smaller mass ratio between the central star and the planet. Given this, if there is a habitable planet around an M-type star, it may be indirectly observed by Subaru Telescope.

Furthermore, if any future 30-meter telescope (such as TMT) can successfully observe a habitable planet blocking the front of the central star, we may be able to analyze atmospheric components of the planet: oxygen molecules suggest the existence of plants, and methane suggests the existence of animals. We expect that great mysteries of mankind will be solved within the next 10 years.

To address these big challenges, there is nothing more important than nurturing young researchers. NAOJ is opening positions to researchers every year. Researches are employed through open recruitment from the public or through recommendation from each project, and currently about 50 researchers are working at NAOJ. And we are also providing research guidance for many graduate students (currently 50 students are accepted) mainly from the Graduate University for Advanced Studies. We hope some of these students will be astronomers who work on the astronomical hottest themes that are now earnestly explored by researchers all over the world, such as planetary formation, planetary exploration, and search for life in the universe.

Compared to last year, this year's report contains a lot more achievements in observational astronomy, theoretical astronomy, equipment development, and engineering development focused on research targets other than those listed above. Astronomy has continuously made enormous and rapid progress. We appreciate your continued support for the growth of astronomy and our activities.



Shoken MIYAMA  
Director General of NAOJ

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# Vega's Rotational Velocity Determined from Spectral Line Profiles

TAKEDA, Yoichi, KAWANOMOTO, Satoshi, OHISHI, Naoko  
(NAOJ)

Although Vega (A0V) has long been well known for its sharp-line nature with a low projected rotational velocity ( $v_e \sin i \sim 20 \text{ km s}^{-1}$ ), we now believe that it is a rapidly rotating star seen nearly pole-on. This evidence came from studying the characteristic flat-bottom profiles of spectral lines, which were first discovered by [1] based on the ultra-high  $S/N$  spectrum and successively analyzed in detail by [2] and [3] to determine  $(v_e, i)$  separately, since they may contain information on rotation-induced gravity darkening. Furthermore, this fact has recently been confirmed also by interferometric observations ([4], [5]), by which the rotation-induced darkening on the stellar disk may be directly studied.

Unfortunately, from a quantitative point of view, consensus has not yet been reached regarding how fast Vega is actually rotating, since the  $(v_e, i)$  results derived so far do not necessarily agree, which diverge into low scale  $v_e$  ( $\sim 160\text{--}170 \text{ km s}^{-1}$ ; [3]) and high scale  $v_e$  ( $\sim 250\text{--}270 \text{ km s}^{-1}$ ; [4, 5]). In order to shed light on this confusing situation while establishing its absolute rotational velocity independently by ourselves, we decided to carry out an extensive line profile study based on the very high- $S/N$  ( $\sim 1000\text{--}3000$ ) and high-resolution ( $R \sim 100000$ ) spectrum data we have recently published ([6]).

Regarding the modeling of a rapidly-rotating star, we made the following assumptions:

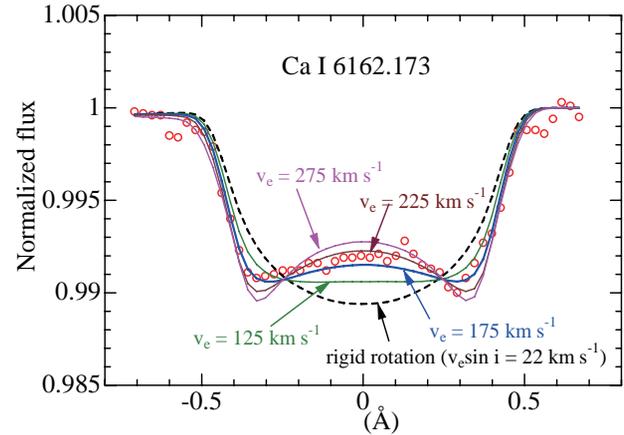
- (1) Point-mass potential (Roche model).
- (2) Axially-symmetric uniform rotation.
- (3) Stellar shape of equipotential surface.
- (4) Local  $T_{\text{eff}}$  depends upon  $g_{\text{eff}}$  as  $T_{\text{eff}} \propto g_{\text{eff}}^\beta$ , ( $\beta \simeq 0.25$ ).

A total of six model parameters are involved in our modeling:  $M$  (mass),  $[X/H]$  (metallicity),  $v_e$  (equatorial rotation velocity),  $i$  (inclination angle),  $R_p$  (polar radius), and  $T_{\text{eff,p}}$  (polar effective temperature). However, we could reduce the number of degrees of freedom to one (only  $v_e$ ) by adequately assuming  $M$  ( $2.3 M_\odot$ ) and  $[X/H]$  ( $-0.5$  dex), by the constraint of  $v_e \sin i = 22 \text{ km s}^{-1}$  (a reasonable value seen from recent determinations) and by the requirement of absolute spectral energy distribution (SED), resulting in a sequence of candidate models [ $i(v_e)$ ,  $R_p(v_e)$ ,  $T_{\text{eff,p}}(v_e)$ ] parameterized by  $v_e$ .

Then, by comparing the observed profiles of carefully selected sufficiently weak lines with the theoretical profiles (computed for various  $v_e$ ; see, e.g., Figure 1), we determined the best-fit  $v_e$  value for each line. Finally, the best solution was concluded to be  $v_e = 175 \text{ km s}^{-1}$ , which further gives  $i = 7^\circ.2$ ,  $R_p = 2.52 R_\odot$ , and  $T_{\text{eff,p}} = 9867 \text{ K}$ . Comparing these with the results of the previous studies, we see that our  $v_e$  solution is near to Hill et al.'s

spectroscopic result ( $\sim 160 \text{ km s}^{-1}$ ; [3]), while it is significantly discrepant from the conclusions of recent interferometric determinations ( $\sim 270 \text{ km s}^{-1}$ ; [4, 5]).

The complete procedures of our analysis and in-depth discussions of the results are described in [7].



**Figure 1:** Example of line-profile fitting. Symbols: observed profile; lines: computed profiles with different  $v_e$ .

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# Observations of 1999 YC and the Breakup of the Geminid Stream Parent

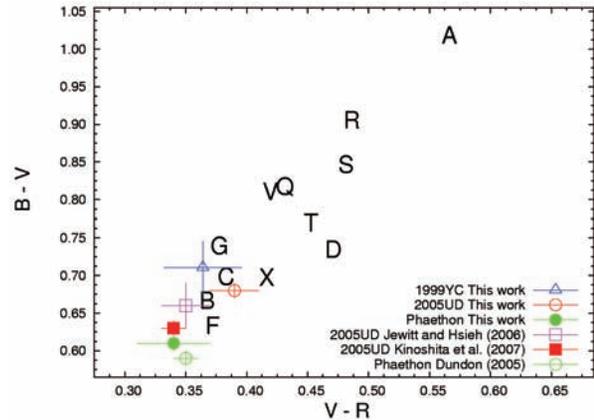
KASUGA, Toshihiro  
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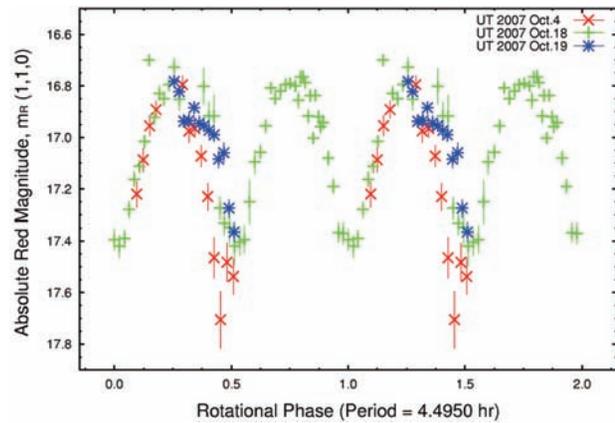
Apollo asteroid 1999 YC may share a dynamical association with the Phaethon-Geminid stream complex. Here, we present photometric observations taken to determine the physical properties of 1999 YC.

The object shows a nearly neutral reflection spectrum, similar to but slightly redder than related objects 3200 Phaethon and 2005 UD (Fig. 1). Assuming an albedo equal to 3200 Phaethon's we find that the diameter is  $1.4 \pm 0.1$  km. Time-resolved broad-band photometry yields a double-peaked rotational period of  $4.4950 \pm 0.0010$  hr while the range of the lightcurve indicates an elongated shape having a projected axis ratio near 1.9:1 (Fig. 2). Surface brightness models provide no evidence of lasting mass loss of the kind seen in active short period cometary nuclei [1].

An upper limit to the mass loss is set at about 0.001 kg/s, corresponding to an upper limit on the fraction of the surface that could be sublimating water ice of 0.001. If sustained over the 1000 yr age of the Geminid stream, the total mass loss from 1999 YC ( $3 \times 10^7$  kg) would be small compared to the reported stream mass ( $1 \times 10^{12} - 1 \times 10^{13}$  kg), suggesting that the stream is the product of catastrophic, rather than steady-state, breakup of the parent object [1].



**Figure 1:** Color plots of  $V - R$  vs.  $B - V$  for PGC and near-Earth asteroids within various Tholen taxonomic classes.



**Figure 2:** Rotational phase vs. absolute red magnitude variation of 1999 YC observed on UT 2007 October 4, 18, and 19.  $mR(1, 1, 0)$  is phased to the double-peaked rotational period of  $P_{rot} = 4.4950 \pm 0.0010$  h.

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[1] Kasuga, T., Jewitt, D.: 2008, *AJ*, **136**, 881.

# Systematic infrared 2.5–5 $\mu\text{m}$ spectroscopy of nearby ultraluminous infrared galaxies: Buried AGNs as a function of galaxy infrared luminosity

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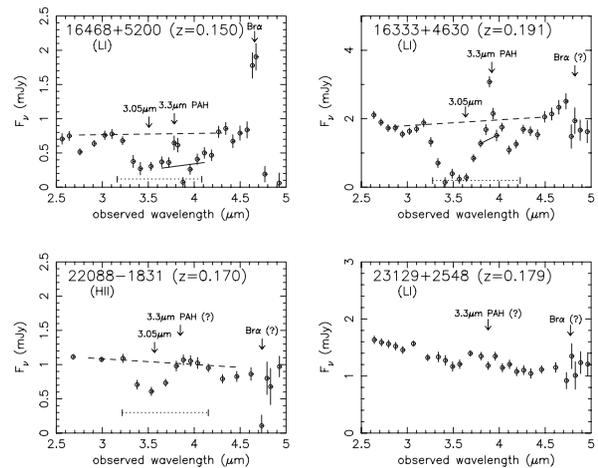
NAKAGAWA, Takao, OHYAMA, Yoichi, SHIRAHATA, Mai, WADA, Takehiko  
(ISAS/JAXA)

Ultraluminous infrared galaxies (ULIRGs) radiate the bulk of their large luminosities ( $>10^{12}L_{\odot}$ ) as infrared emission. This means that (1) very luminous energy sources are present hidden behind dust, (2) the bulk of the energetic radiation from the hidden energy sources is once absorbed by the surrounding dust, and (2) the heated dust emits strong infrared radiation. The energy source can be nuclear fusion inside rapidly formed stars (starburst) and/or active mass accretion onto a central supermassive blackhole (AGN). Distinguishing the hidden energy sources of ULIRGs is indispensable not only to understand the true nature of the ULIRG population, but also to unveil the history of star-formation and supermassive blackhole growth in the dust-obscured portion of the universe.

Since AGN activity is spatially very compact, it can easily be *buried* (= obscured in virtually all directions) in dusty ULIRG's cores. Unlike an AGN surrounded by torus-shaped dusty medium, which is easily detectable through optical spectroscopy, such *buried* AGNs are very difficult to find. However, detecting *buried* AGNs and estimating their energetic role are the most important issue in the current ULIRG study. For this purpose, observations at wavelengths of low dust extinction are necessary, and infrared 3–4  $\mu\text{m}$  (rest-frame) spectroscopy is an effective means. First, PAH (polycyclic aromatic hydrocarbon) emission features, found at 3.3  $\mu\text{m}$  (rest-frame), are detected in starburst galaxies, but not in AGNs, because of the destruction of PAHs. Hence, the presence or absence of the PAH emission can be used to identify the energy sources. Next, in a normal starburst, where the stellar energy sources and dust are spatially well mixed, there is an upper limit for the optical depth of dust absorption features found at 3–4  $\mu\text{m}$ , while it can be arbitrarily large in a buried AGN, where the energy source is more centrally concentrated than the surrounding dust [1].

Using the Subaru telescope, we have so far performed systematic infrared 2.8–4.1  $\mu\text{m}$  spectroscopy of ULIRGs at  $z < 0.15$ , and found luminous buried AGN signatures in a significant fraction of the observed ULIRGs [1]. AKARI satellite has spectroscopic capability at 2.5–5  $\mu\text{m}$ , unaffected by Earth atmosphere, as well as sufficient sensitivity to obtain high quality spectra of ULIRGs at  $z < 0.3$  in a short exposure time. By extending our study to more distant ULIRGs at  $z > 0.15$  (Figure 1), many ULIRGs with higher infrared luminosities are

included. We can thus investigate the detected buried AGN fraction as a function of galaxy infrared luminosity. We found that with increasing galaxy infrared luminosity, buried AGNs become *relatively* important energetically, but the remaining star-formation-originating luminosities are also higher (i.e., a star formation rate is higher and more stars will be formed). Buried AGNs can have a particularly strong feedback to the surrounding gas and dust in host galaxies. Our results may observationally support the widely-proposed AGN feedback scenario as the origin of the galaxy downsizing phenomenon [2], where more massive galaxies with currently larger stellar masses have finished their major star-formation in an early cosmic age on a shorter time scale.



**Figure 1:** Examples of AKARI 2.5–5  $\mu\text{m}$  spectra of ULIRGs at  $z = 0.15–0.3$ . All ULIRGs show low-equivalent-width PAH emission and/or strong dust absorption features, suggesting that buried AGNs are energetically important. The 3.05  $\mu\text{m}$  absorption feature is due to ice-covered dust grains. Dashed lines mean continuum levels. Dotted lines indicate the wavelength range where the 3.05  $\mu\text{m}$  ice absorption feature is strong. HII, L1 mean ULIRGs optically classified as HII-regions and LINERs, respectively.

## References

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# Nobeyama Millimeter Array HCN( $J=1-0$ ) and HCO<sup>+</sup>( $J=1-0$ ) observations of luminous infrared galaxies: The final sample

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Luminous infrared galaxies radiate very large luminosities ( $>10^{11}L_{\odot}$ ) as infrared dust emission, suggesting that powerful energy sources (starburst and/or AGN) are present, surrounded by dust. Unlike AGNs obscured by torus-shaped dusty medium, putative AGNs in LIRG's nuclei are likely to be obscured by a large amount of dust and gas in all directions (= buried), making the identification of AGNs difficult.

In an AGN, X-ray emission ( $E > 2$  keV) is strong and an emission surface brightness is high, compared to a starburst. In the case of a buried AGN at a LIRG's nucleus, we may be able to detect these AGN signatures, through the effects to the surrounding gas and dust. In fact, observationally, the brightness-temperature ratios of HCN( $J=1-0$ ) and HCO<sup>+</sup>( $J=1-0$ ) emission lines at  $\lambda \sim 3.4$  mm ( $\nu \sim 89$  GHz) are different between weakly-obscured AGN-dominated galaxy nuclei and starburst galaxies. Through the systematic observations of LIRGs, using the Nobeyama Millimeter Array (NMA), we have found that LIRGs with luminous buried AGN signatures in infrared spectra tend to show high HCN( $J=1-0$ )/HCO<sup>+</sup>( $J=1-0$ ) brightness-temperature ratios, as seen in weakly-obscured AGNs. However, our sample of buried AGN candidates has been limited to sources with very high infrared luminosities ( $>10^{12}L_{\odot}$ ), so that a scenario without invoking an AGN (e.g., different molecular gas properties in galaxies with different luminosities) has been proposed to explain the high ratios.

In this paper, we have observed four LIRGs (five nuclei) of moderate infrared luminosities ( $<10^{12}L_{\odot}$ ), with or without buried AGN signatures at other wavelengths (Figure 1), and confirmed that those with AGN signatures still show high ratios (Figure 2), reinforcing that the buried AGN scenario (not different molecular gas properties scenario) is at work [1]. The high ratios could be explained by an HCN abundance enhancement in an AGN, due to strong X-rays, or infrared radiative pumping of HCN molecules by hot dust heated by an AGN (= a high emission surface brightness energy source). Sophisticated theoretical models are needed to properly interpret our observational results.

## Reference

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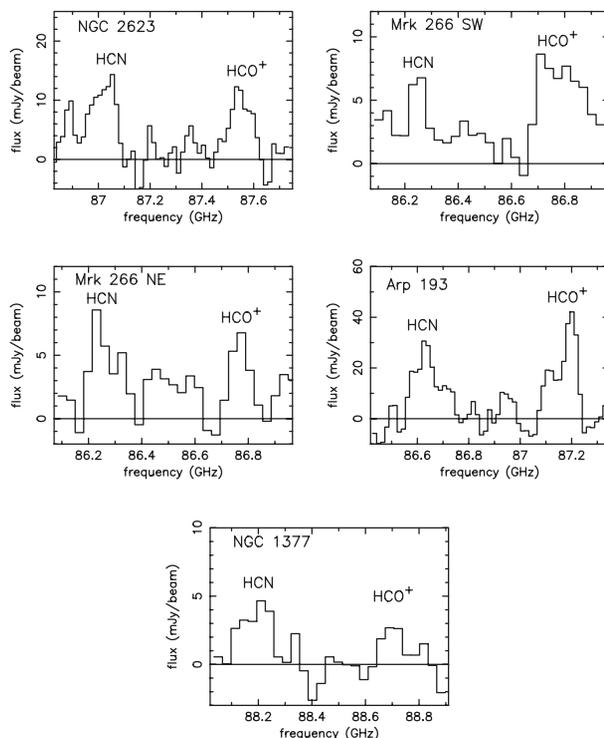


Figure 1: Example spectra of LIRGs around HCN( $J=1-0$ ) and HCO<sup>+</sup>( $J=1-0$ ) lines, obtained with NMA.

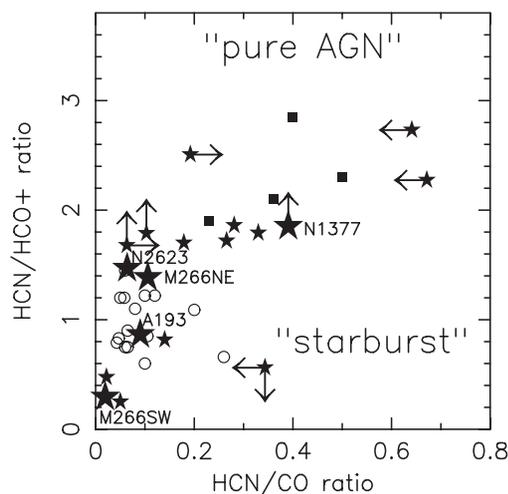


Figure 2: Filled squares: weakly-obscured AGN-dominated galaxy nuclei. Open circles: starburst galaxies. AGNs distribute at the upper part, compared to starbursts. Filled stars: LIRGs. Large filled stars with objects name: LIRGs observed this time. LIRGs with luminous buried AGN signatures at other wavelengths tend to distribute at the upper part occupied by AGN-dominated galaxy nuclei.

# Same-Beam VLBI Observations in SELENE (Kaguya)

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SELENE (Kaguya) was successfully launched on Sep.14, 2007. It consists of a three-axis stabilized Main satellite and two small spin stabilized satellites, called Rstar and Vstar. Rstar relays the Doppler signal between the main satellite and ground station (4-way Doppler) for the world's first direct measurement of the gravity field on the far side of the Moon. The differential VLBI between Rstar and Vstar are used to determine the gravity field of the Moon precisely, especially near the limb region and for the low degree coefficients of the spherical harmonics.

Rstar and Vstar only transmit three carriers (S1: 2212, S2: 2218, S3: 2287 MHz) in S-band and one carrier (8456 MHz) in X-band; the differential phase delay (DPD) between Rstar and Vstar was obtained from the relation between correlation phase and frequency of the four carriers. In this case, the differential correlation phase at each frequency has to be estimated without the  $2\pi$  ambiguity. This imposes strict conditions such as the error of the correlation phase must be lower than 4.3 deg rms. To resolve the  $2\pi$  ambiguity, we have mainly used samebeam differential VLBI, in which Rstar and Vstar are observed simultaneously with the main beam of the receiving antennas.

Figure 1 shows the correlation phase fluctuations of Rstar and Vstar and their differences at S1, S2 and S3. The separation angle was between 0.04 to 0.052 deg. The integral time is 30 seconds and long-time trends, which were mainly attributed to the error of the a priori orbit, were removed by fitting 9th order polynomials. The correlation phase fluctuations of Rstar and Vstar with amplitude of about 10 deg are mainly caused by the atmosphere and ionosphere, they are very similar because the propagation paths from Rstar and Vstar to the ground stations are nearly the same. In addition, the profiles of phase fluctuations at S1, S2 and S3 are also very similar because the difference in receiving frequencies is not large. These fluctuations were reduced to a low level by differencing the correlation phases of Rstar and Vstar as shown in Fig. 1. The differences between correlation phase fluctuations of Rstar and Vstar at S1, S2 and S3 were 0.99, 1.00, 0.83 deg RMS, respectively, which satisfied the condition for obtaining DPD.

Figure 2(a) shows S-band DPD for three baselines, Ishigaki-Mizusawa, Mizusawa-Iriki and Iriki-Ishigaki. The separation angle is between 0.04 to 0.09 deg. The fluctuations of DPD for the three baselines after removing long-time trends by fitting 9th order polynomials are shown in Fig. 2(b). They were very small, about 0.0003

m (1 ps) RMS. In addition, the DPD for each baseline was estimated separately before and after 11 UT, the fluctuations have no jump at 11 UT as shown in Fig. 2(b). The closure DPD among the three baselines, i.e., the sum of DPD of (IS-MZ)+(MZ-IR)+(IR-IS) shown in Fig. 2(a) was also calculated. As shown in Fig. 2(c), it was very small, less than 0.001 m (3 ps). Considering one ambiguity at S-band corresponds to 0.136 m in DPD, it is confirmed that the DPD has been successfully estimated.

We also performed the orbit determination of Rstar and Vstar using the GEODYN II software. The orbit errors were evaluated by overlap analysis, where orbit differences were computed between two overlapping arcs. Orbit consistency was greatly improved by including the VLBI data. Without VLBI data, orbit consistency for Rstar is 18.01 m, and that for Vstar is 75.26 m. With VLBI data, the consistency was 2.79 m for Rstar and 10.48 m for Vstar. In addition, the precise VLBI data have been used to improve the lunar gravity field model.

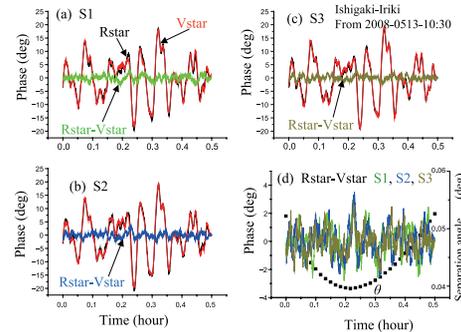


Figure 1: Correlation phase fluctuations of Rstar and Vstar at S1, S2, S3 and difference between Rstar and Vstar.

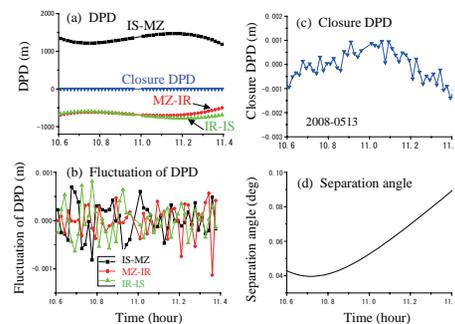


Figure 2: (a) Differential phase delay, (b) their fluctuations, (c) closure phase delay, and (d) separation angle.

## Reference

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# Formation Scenario for Wide and Close Binary Systems

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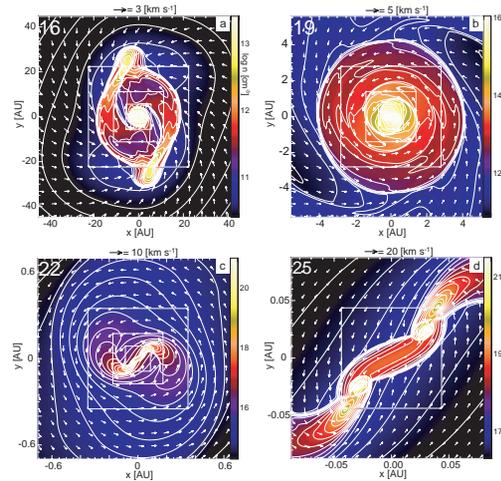
MATSUMOTO, Tomoaki  
(Hosei University)

TOMISAKA, Kohji  
(NAOJ)

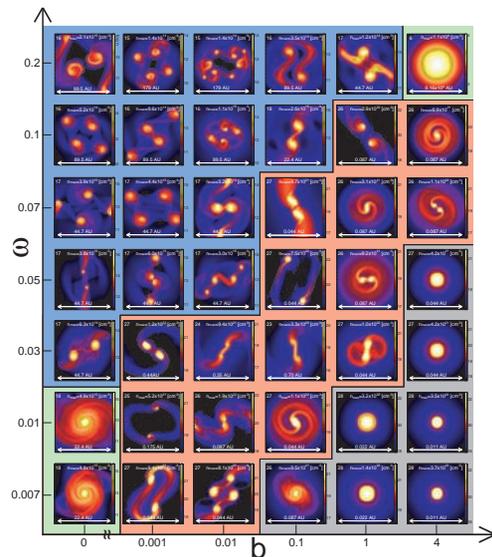
INUTSUKA, Shu-ichiro  
(Kyoto University)

Observations indicate that stars are born as binary or multiple systems. The star formation occurs in the collapsing cloud core. When fragmentation occurs in the collapsing cloud, fragments evolve into binary or multiple stellar systems. However, we cannot directly observe fragmentation process, because it is embedded in a dense cloud core. Thus, we need to investigate the star and binary formation using numerical simulation. In the collapsing cloud, two adiabatic cores hierarchically appear. The gas density increases as the cloud collapses. When the central number density exceeds  $n = 10^{11} \text{ cm}^{-3}$ , the first adiabatic core (hereafter the first core) is formed. Then, further collapse begins inside the first core, and the second core (i.e., the protostar) is formed when the gas density reaches  $n = 10^{21} \text{ cm}^{-3}$ . The evolution of such collapsing cloud was investigated in some past studies. When the molecular cloud has a certain amount of the angular momentum and unmagnetized, fragmentation is possible to occur just after the first core formation [1]. However, only a few studies focused on the evolution of magnetized clouds. In this study, using resistive MHD nested grid simulations, we investigated the evolution of the magnetized disk, and found the condition for fragmentation and binary formation [2]. Figure 1 shows the cloud evolution for a typical model with different scales. In this model, fragmentation occurs twice. The first fragmentation occurs just after the first core formation. After the first core formation, the gas collapsing timescale shortens and the perturbation to induce fragmentation grows. The second fragmentation occurs just before the protostar formation. This is because (i) further collapse is initiated without the formation of spiral arms owing to the removal of the magnetic field, and (ii) the transportation of the angular momentum by the magnetic braking becomes ineffective. Figure 2 shows results of parameter survey, in which the magnetic field and angular velocity of the cloud core is parameterized. In the figure, fragmentation occurs at the epoch of the first core formation in models with blue background, while at the epoch of the second core formation in models with red background. On the other hand, models with gray background evolve to a single star without fragmentation. The figure indicates that the rotation promotes fragmentation, while the magnetic field suppresses it. In addition, fragmentation is possible even in the high-density gas region. The binary separation depends on both the magnetic field and rotation of the cloud. Models showing fragmentation at lower density

(i.e., the epoch of first core formation) evolve to a wide binary system, while the model showing fragmentation at the higher density (i.e., the epoch of the second core formation) evolve to a close binary system.



**Figure 1:** Density distributions (color and contours) with difference scales. Arrows indicate the velocity vector.



**Figure 2:** Results of parameter survey. The density distribution at fragmentation epoch for each model is plotted against the strength of the magnetic field (x-axis) and degree of the rotation (y-axis).

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# Magnetic Field Configuration at the Galactic Center Investigated by Wide Field Near-Infrared Polarimetry

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HOUGH, James H.  
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SATO, Yaeko  
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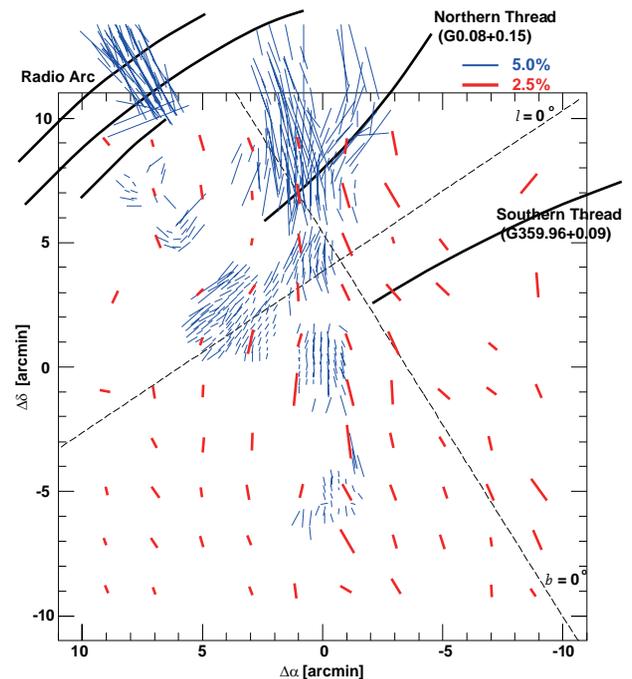
SUGITANI, Koji  
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OKUDA, Haruyuki  
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The magnetic field (MF) configuration at the Galactic center (GC) has been investigated with a wide variety of methods [1]. The discovery of polarized radio emission extending perpendicular to the Galactic plane, such as non-thermal radio filaments and polarized plumes, has provided early evidence for a substantial poloidal component of the MF at the GC. However, recent far-infrared (FIR) and sub-millimeter (sub-mm) polarimetric observations point out that the MF is generally parallel to the Galactic plane.

Previous near-infrared (NIR) polarization measurements of the GC were discussed in terms of selective absorption by the intervening interstellar dust grains in the Galactic disk. To our knowledge, no one has studied the MF configuration at the GC with NIR polarimetry. In this work, we conducted NIR polarimetric observations toward the GC, and demonstrate that NIR polarization of point sources can provide information on the MF structure at GC [2].

We carried out NIR polarimetric observations with the NIR polarimeter SIRPOL on the IRSF telescope. Comparing the Stokes parameters between high extinction stars and relatively low extinction ones, we obtained a polarization originating from magnetically aligned dust grains at the GC (Fig. 1). The distribution of the position angles shows a peak at  $\sim 20^\circ$ , nearly parallel to the Galactic plane, suggesting a toroidal MF configuration. The derived direction of the MF is in good agreement with that obtained from FIR/sub-mm observations, which detect polarized thermal emission from dust in the molecular clouds at the GC [3, 4, 5]. Our results show that by subtracting foreground components, NIR polarimetry allows investigation of the MF structure at the GC.



**Figure 1:**  $K_s$  band polarization map for the GC from our observation (*red bars*). The polarization map at the GC derived from FIR/sub-mm observations (*blue bars*) is also shown [3, 4, 5]. The length of the bars is proportional to the measured degree of polarization, and their orientation is drawn parallel to the inferred magnetic field direction. Some prominent radio filaments are shown as heavy dark lines.

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# Interstellar Extinction Law toward the Galactic Center: V, J, H, and K<sub>S</sub> Bands

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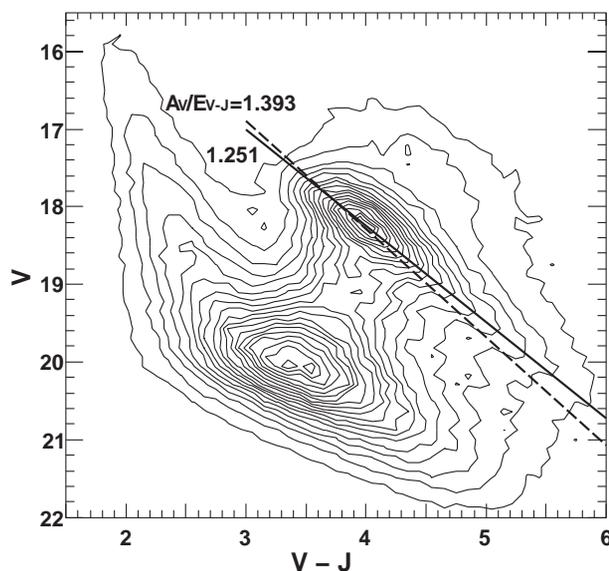
TAMURA, Motohide, KANDORI, Ryo  
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The wavelength dependence of interstellar extinction provides important diagnostic information about the dust grain properties. Interstellar extinction law shows a large range of variability from one line of sight to another, especially in the ultraviolet and optical wavelengths. In comparing the wavelength dependence among different lines of sight, the normalization of the extinction curves by the total extinction  $A_V$ , instead of the usual color excess  $E_{B-V}$ , is vitally important. Toward the Galactic center (GC), assuming only that the center of stellar distribution in the lines of sight is at the same distance from us and that the foreground extinction is patchy, we can directly derive the ratio of the total to selective extinction (e.g.,  $A_V/E_{V-J}$ ), from which we can obtain the extinction ratio normalized by the total extinction (e.g.,  $A_J/A_V$ ). By plotting the apparent magnitude versus the color excess of a group of stars, one obtains a straight line with the slope equal to the ratio of total to selective extinction ratio.

In this work, we have determined the ratios of total to selective extinction directly from observations in the optical V band and near-infrared J band toward the GC [1]. The OGLE Galactic bulge fields were observed with the SIRIUS camera on the IRSF telescope. We constructed color-magnitude diagrams, and measured the regression of the mean color of red clump stars in the GC on their mean magnitude. We obtain  $A_V/E_{V-J} = 1.251 \pm 0.014$  (see Fig. 1) and  $A_V/E_{V-J} = 0.225 \pm 0.007$ . From these ratios, we have derived  $A_J/A_V = 0.188 \pm 0.005$ ; if we combine  $A_J/A_V$  with the near-infrared extinction ratios obtained by Nishiyama et al. [2] for more reddened fields near the GC, we get  $A_V : A_J : A_H : A_{K_S} = 1 : 0.188 : 0.108 : 0.062$ , which implies steeply declining extinction toward the longer wavelengths.

The resultant  $A_\lambda/A_V$  in the  $JHK_S$  wavelength range is a steeply declining function. The frequently used extinction curve derived by Cardelli et al. (1989 [3]) (for  $R_V = 3.1$ ,  $A_V/E_{V-J} = 1.393$ , see Fig. 1), which is based on Rieke & Lebofsky (1985 [4]), decreases much more slowly toward the longer wavelengths. In particular, it is striking that the  $K_S$  band extinction is  $\approx 1/16$  of the visual extinction  $A_V$ , much smaller than one tenth of  $A_V$  so far employed ([3, 4]).



**Figure 1:**  $V$  versus  $V-J$  color magnitude diagram of the OGLE bulge fields. The *solid* and *dashed* lines show the direction of the interstellar extinction with  $A_V/E_{V-J} = 1.251$  (this work [1]), and 1.393 ( $R_V = 3.1$  [3]), respectively. It is clear that the *solid* line is along the ridge of bulge red clump stars.

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# Evolution of the Anemone Active Region NOAA 10798 and the Related Geo-Effective Flares and CMEs

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(Kyoto University)

OKA, Mitsuo  
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KATAOKA, Ryuho  
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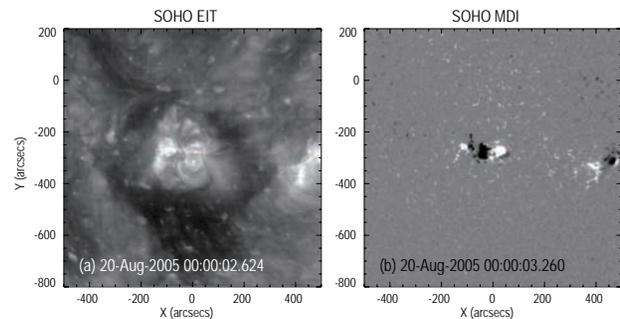
FUJIKI, Ken'ichi  
(Nagoya University)

GOPALSWAMY, Nat  
(NASA GSFC)

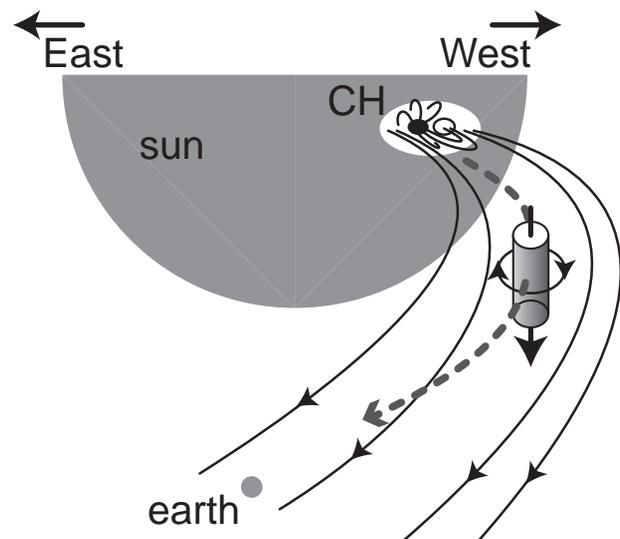
Space weather has attracted a lot of attention in recent times. Space weather research involves various related fields. Therefore, in order to understand what kind of events on the solar surface can generate a large geomagnetic storm, it is necessary to study active phenomena on the solar surface and the propagation in the interplanetary space, in relation to the surrounding magnetic structure. We investigated the evolution of the Active Region (AR) NOAA 10798 and the features on/near the sun and in the interplanetary space, since it generated flares/coronal mass ejections (CMEs) that caused a large geomagnetic storm on 2005 August 24 with the minimum Dst index of  $-216$  nT [1].

The AR emerged in the middle of a small coronal hole (CH), and formed a sea anemone like configuration [2] (Fig. 1).  $H\alpha$  filaments were formed in the AR, which have southward axial field. Two halo CMEs associated with M-class flares occurred on 2005 August 22. The speeds of the CMEs were extremely fast, and recorded about  $1200$  and  $2400$  km s $^{-1}$ , respectively. The second CME was especially fast, and probably caught up and interacted with the first (slower) CME during their travelings toward the earth (Fig. 1). The interplanetary disturbances with strong southward magnetic field of about  $-50$  nT and strong compression of plasma were produced.

The reasons why the CMEs were particularly geo-effective were the high speeds of the two CMEs and their interaction as well as the CMEs traveled directly toward the earth. The high speeds of the CMEs are more notable. We suggest that the fast CMEs are probably a consequence of the eruption inside a CH from an anemone AR. Eruptive activities of anemone ARs are usually low [3]. However, the relation between anemone ARs and fast solar winds have been paid attentions to, and recent papers (e.g. [4]) have also reported the association between fast Halo CMEs and CHs. Our results indicate that flares/CMEs from anemone ARs are also important for space weather researches.



**Figure 1:** Coronal feature of the AR NOAA 10798. (a): An EUV ( $195 \text{ \AA}$ ) image obtained with *SOHO*/EIT. The bright region near the center of the image is the AR. The surrounding dark region is a CH. (b): A magnetogram taken with *SOHO*/MDI.



**Figure 2:** Schematic cartoon of AR NOAA 10798 and related flares/CMEs. The ejected plasma is bent eastward by the surrounding magnetic field with positive magnetic polarity. The ejecta from the anemone AR becomes a magnetic cloud (shown as a cylinder) that has a southward axial magnetic field and is approaching to the earth.

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# Strongly Blueshifted Phenomena Observed with *Hinode*/EIS in the 2006 December 13 Solar Flare

ASAI, Ayumi, HARA, Hirohisa, WATANABE, Tetsuya, IMADA, Shinsuke (NAOJ)

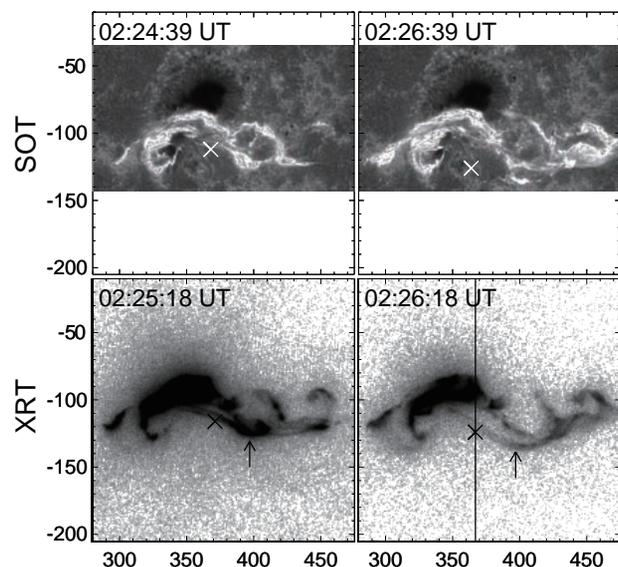
SAKAO, Taro, NARUKAGE, Noriyuki (JAXA)

CULHANE, J. L. (MSSL)

DOSCHEK, G. A. (Naval Research Laboratory)

Solar flares are very spectacular, and are accompanied by a variety of plasma motions. These plasma motions can be observed as phenomena accompanied by line shifts (Doppler shifts) in spectroscopic observations. We examined in detail strongly blueshifted emission lines observed with the EUV Imaging Spectrometer (EIS) on board the *Hinode* satellite that were observed to be associated with the 2006 December 13 solar flare [1]. We found two kinds of blueshifted phenomenon associated with the X3.4 flare.

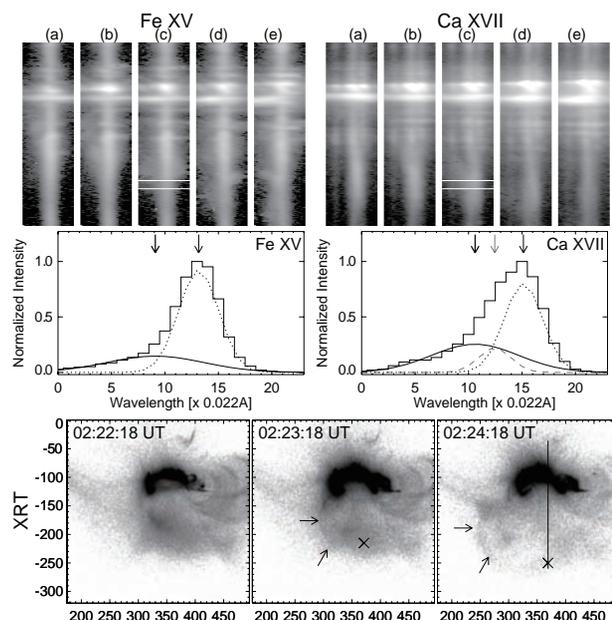
One was related to a plasmoid ejection seen in soft X-rays (SXR), and we call it BS1. It was very bright in all the lines used for the observations (Fig. 1). The Doppler velocity is about  $250 \text{ km s}^{-1}$ . Although they are associated with the plasmoid ejection, there are no corresponding flare ribbons. Therefore, we concluded that this blueshift is the ejected plasma, and is not an evaporation flow.



**Figure 1:** *Top:* Ca II (H-line) images taken with *Hinode*/SOT. *Bottom:* Soft X-ray images obtained with *Hinode*/XRT (negative images). The cross signs (x) point to the positions of the BS1. The arrows in the bottom right panel follows the plasmoid ejection seen in XRT. The vertical line in the bottom right panel shows the position of the EIS slit.

The other, on the other hand, was very faint, and showed spectra that broaden in the wavelength space

(BS2; Fig. 2). The center of the blueshifted component recorded a Doppler velocity of about  $100 \text{ km s}^{-1}$ , the drift velocity is about  $450 \text{ km s}^{-1}$ . These components are observed only in the hottest lines of the raster observation, and therefore, the plasma must be heated to more than 2 MK. The blueshift region corresponds to the propagation of the coronal wavelike (arc-shaped) ejection seen in the SXR images. The SXR ejection is thought to be an MHD fast-mode shock wave with the Alfvén Mach number of about 1.4. This is, therefore, the first spectroscopic observation of an MHD fast-mode shock wave associated with a flare.



**Figure 1:** *Top:* Time sequenced spectra of Fe xv (left) and Ca xvii (right) windows observed with *Hinode*/EIS. *Middle:* Normalized spectra at BS2 in the windows. The solid histograms show the spectra, and the dotted and solid black lines are the fitting results that represent the main and the blueshift components of the line. The peaks of each line are shown with thin and thick arrows. *Bottom:* SXR negative images taken with the *Hinode*/XRT. The arrows show the front of the wavelike ejection. Crosses (x) represent BS2 determined by the Fe xv line. The vertical line in the bottom right panel shows the position of the EIS slit.

## Reference

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# Exploring the substellar temperature regime down to $\sim 550$ K

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KUZUHARA, Masayuki  
(University of Tokyo)

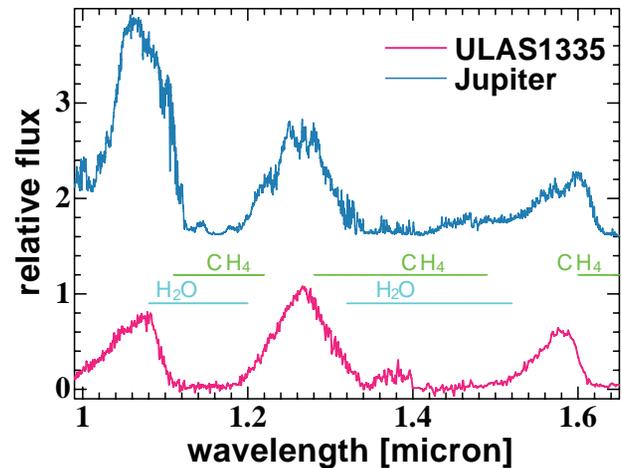
BURNINGHAM, Ben, PINFIELD, David  
(University of Hertfordshire)

and the UKIDSS Cool Dwarf Science Working Group

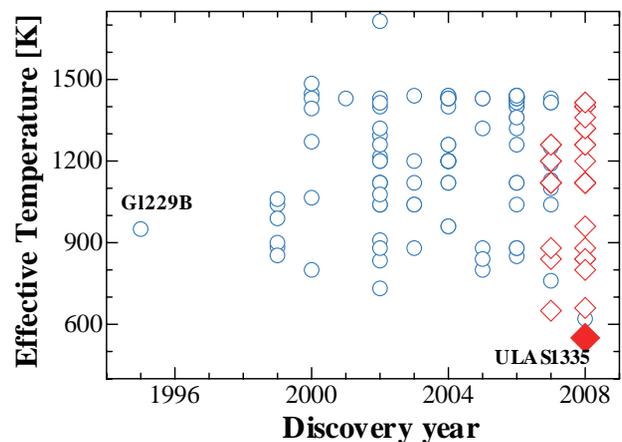
Observations in the 1980–90s have established that the conventional spectral sequence from O to M stars is followed by even cooler dwarf stars — L and T dwarfs (effective temperature  $T_{\text{eff}} = 2100\text{--}1300$  K and  $T_{\text{eff}} \lesssim 1300$  K, respectively). Some L dwarfs and all T dwarfs correspond to “brown dwarf” whose mass is below 0.075 solar masses.

Previous large area surveys (e.g., 2MASS, SDSS) have discovered about 600 L dwarfs and 100 T dwarfs so far. On the other hand, a deeper large area survey is necessary to detect cooler (i.e., intrinsically fainter) brown dwarfs. Ongoing UKIDSS Large Area Survey (LAS) meets the requirement. The UKIDSS/LAS will cover 4000 sq. degrees using the UKIRT Wide Field Camera (WFCAM), with 4 mag deeper detection limit than 2MASS. From the resultant data, the coolest T dwarfs and, further more, a new class of dwarfs cooler than T are expected to be found. Such ultra cool dwarfs are referred to be “Y” dwarfs, and their spectra are expected to be analogous to the spectrum of the Jupiter.

The Cool Dwarf Science Working group, an international collaboration including UK and Japanese astronomers, has developed the selection method using UKIDSS  $YJHK$  colors and the combination of SDSS  $z\text{-}J$ ,  $i\text{-}z$  to select L/T and Y dwarf candidates. Followup spectroscopy using large telescopes such as Subaru and Gemini is adopted to determine their spectral types finally. As a result, we have discovered 72 T dwarfs from the UKIDSS early data so far, indicating that our method is highly effective (e.g., [1, 2]). Six of them are classified as T8.5–T9, coolest types among known T dwarfs. In particular, ULAS J133553.45+113005.2 (ULAS335) has been shown to be the coolest T dwarf ever known (Figures 1, 2), with the effective temperature of 550 K and the mass of about 20 Jupiter masses [3]. After the completion of UKIDSS/LAS, we expect to discover Y dwarfs together with hundreds of T dwarfs, which will shed light on the formation processes of substellar objects, the Initial Mass Function (IMF) and the formation history.



**Figure 1:** Near-infrared spectrum of the coolest brown dwarf ever discovered (ULAS1335) (lower). For comparison spectrum of the Jupiter is shown (above ULAS1335). Broad absorption bands of  $\text{H}_2\text{O}$  and  $\text{CH}_4$  are seen in their spectra.



**Figure 2:** History of the coolest brown dwarf discovered: T dwarfs discovered by UKIDSS ( $\diamond$ ) and previous surveys ( $\circ$ ) are plotted against their discovery year.

## References

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# High-Precision Correction of the Instrumental Polarization for High-Level Polarimetry of the Sun

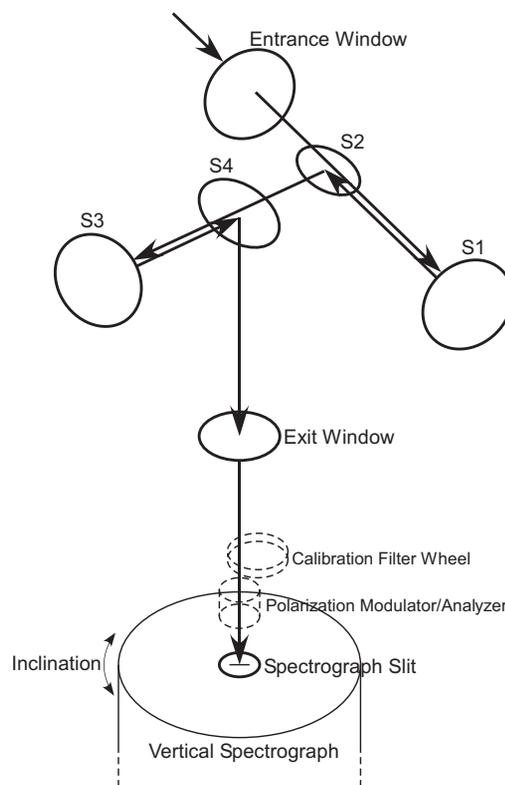
HANAOKA, Yoichiro  
(NAOJ)

The polarization of the light from the Sun shows the surface magnetic field of the Sun. However, generally it is not avoidable for big solar telescopes to have some folding mirrors, which vary the polarization of the incident light. In such cases, the instrumental polarization has to be corrected with a high precision to evaluate the incident polarization of the light from the Sun.

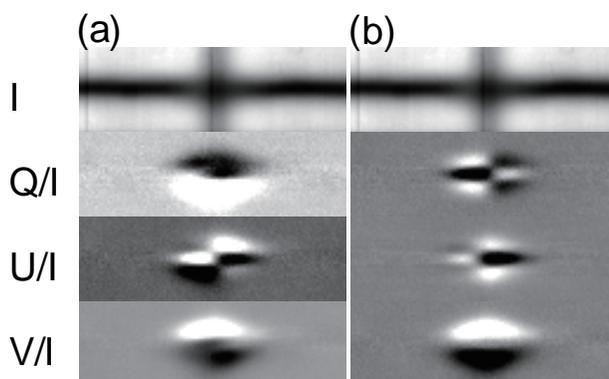
On the other hand, to carry out high precision polarization measurements, it is also required to calibrate the polarimeter itself with a high precision as well. Generally unpolarized light (usually the continuum from the Sun) is necessary for the calibration, but the instrumental polarization changes the unpolarized incident light into polarized light. Therefore, the calibration with unpolarized light is difficult in the system with the instrumental polarization. This fact means that the evaluation of the parameters of the instrumental polarization and that of the calibration parameters for the polarimeter are coupled together.

We installed a polarimeter of NAOJ with ferroelectric liquid crystals into the Domeless Solar Telescope of Hida Observatory, Kyoto University, and developed a method to evaluate the parameters of the instrumental polarization and those of the polarimeters simultaneously [1]. As shown in Figure 1, the Domeless Solar Telescope has two folding mirrors, which produce the instrumental polarization. However, the spectrograph rotation and the function of the turning-over of the telescope attitude cancel the instrumental polarization virtually. We evaluated the polarization parameters with a high precision using these characteristics, and we succeeded to reproduce the measured instrumental polarization by the calculated ones with RMS errors of 0.03 %, which meets the required error level to measure the weak polarization in chromospheric lines. Figure 2 shows a sample set of the polarization spectra with the slit positioned on a sunspot. The raw data show significant crosstalk, but the spectra corrected for the instrumental polarization show almost symmetric polarization profiles.

The method developed in this work is applicable for various solar telescopes with oblique reflections. Furthermore, it will become a guide to consider the method to correct the instrumental polarization for the solar telescope designed in future, where the oblique reflections cannot be avoided.



**Figure 1:** Layout of the light path in the Domeless Solar Telescope. Mirrors S2 and S4 are folding mirrors. The signs of Stokes Q and U signals into the polarimeter can be reversed with a rotation of the spectrograph, and that of Stokes V can be done so with the turning-over of the telescope.



**Figure 2:** A sample set of Stokes Spectra of the Fe 6302.5 line taken at a sunspot. The horizontal and vertical directions are along the slit and the dispersion, respectively. Panels (a) and (b) show the spectra before and after the correction of the instrumental polarization.

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# Sensitivity improvement of gravitational wave detector TAMA300

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The TAMA project of NAOJ has been developing a laser interferometer gravitational wave detector TAMA300. Here, we report a sensitivity improvement of TAMA300.

We have been working on the new vibration isolation system, called “Seismic Attenuation System (SAS)”, since 2005. The sensitivity of TAMA300 at around 100 Hz was previously limited by alignment control noise. In order to improve the sensitivity in the low frequency band, we replaced the previous vibration isolation system to the SAS for all of four test masses. As a result of the installation and tuning works, we could decrease the control bandwidth of the alignment servo system. It leads an update of the sensitivity.

The vibration isolation performance of the SAS essentially relies on low resonant frequencies of the mechanical systems. The SAS has a multiple-stage structure with a total height of about 2.5 meters so as to realize seismic filtration from the low frequency such as 50 mHz. In order to suppress excited motions of the low frequency resonances, the active local controls are applied. All of the local control servos are realized by digital control systems based on LabVIEW. With the intrinsic vibration isolation performance of the SAS and the local control systems, test mass motions are well stabilized in all frequency band. It was confirmed by the optical path length measurement of the 300-m arm cavity that the RMS motion of the cavity length is improved from  $0.84 \mu\text{m}$  to  $1.4 \mu\text{m}$  [1].

After the shaking down and tuning of the SAS, the power recycled operation of the 300-m Fabry–Perot Michelson interferometer has been recovered. The current sensitivity of TAMA300 is shown in Figure ???. The revised wave front sensing technique for the alignment control helps to keep the stability of the sensitivity in the kHz band, while the improved sensitivity at the frequency band below 100 Hz is not compromised. As a result, the sensitivity between 0.1 Hz to 400 Hz was improved. The floor sensitivity is  $h = 1.3 \times 10^{-21} \text{ m}/\sqrt{\text{Hz}}$  at around 1 kHz. The contribution of the alignment control noise is also shown in the figure. It indicates that this noise is not limiting the current sensitivity. This was owing to the reduction of the alignment control bandwidth down to 3 Hz.

Figure 2 shows how distant compact binaries TAMA300 should be able to detect. Each curve means the event at the indicated distance is detected with S/N of 10, assuming the direction of the event and the polarization are optimal for the detector. For all binary masses farther

events should be detected than before. Particularly, the distance for stellar mass binaries are remarkably improved as the low frequency sensitivity contributes to it.

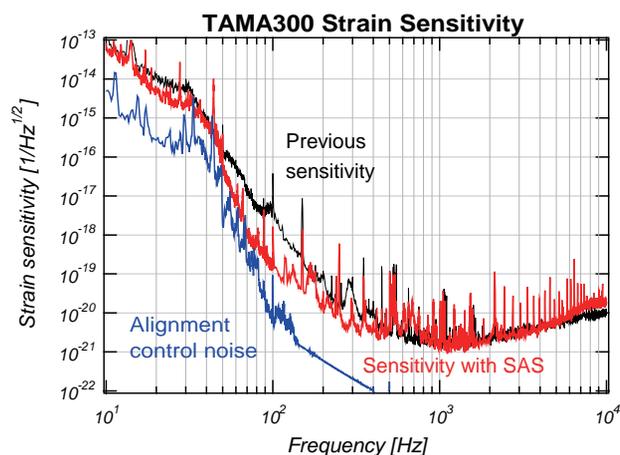


Figure 1: Current sensitivity of TAMA300 (red), being compared with the previous sensitivity (black). The blue line indicates the estimation of the alignment control noise.

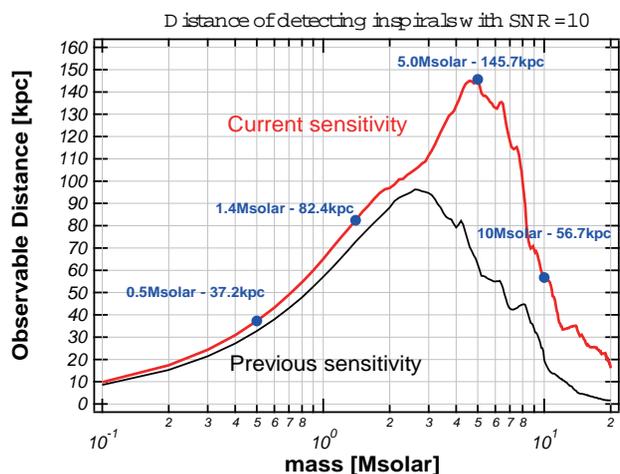


Figure 2: “Observable distance” indicates how distant compact binaries TAMA300 should be able to detect. This shows the event at this distance is detected with S/N of 10, assuming the direction of the event and the polarization are optimal for the detector.

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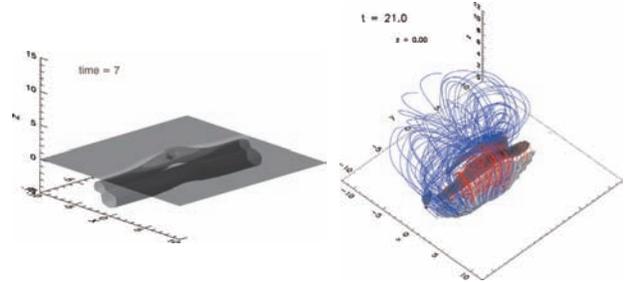
# Theoretical and Observational Studies Using *Hinode* of the Magnetic-Field Emergence into the Solar Atmosphere

MAGARA, Tetsuya  
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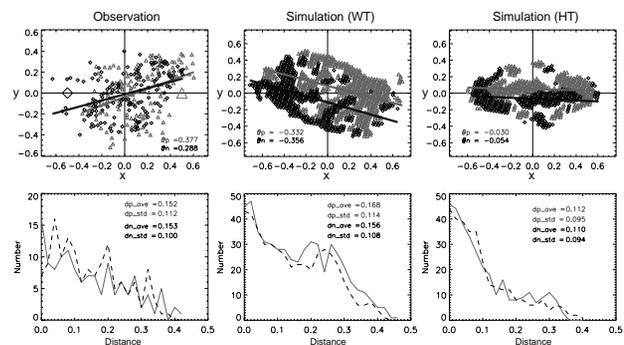
We have performed a three-dimensional MHD simulation for the emergence of a partially split flux tube into the solar atmosphere, which is focused on the mechanism of the activity observed on the Sun such as the merging of magnetic polarity regions in the photosphere and the formation of multiple flux domains in the corona. The simulation reproduces that a small polarity region associated with a rotational flow merges into a main polarity region, and later in this main polarity region is developed a rotational flow in the same direction as the small polarity region. In accordance with the photospheric merging process, multiple flux domains form in the corona and a current layer separating these flux domains travels outward as the flux domains expand. We also investigated the nature of helicity injection in a multi-pole system.

We reported a result of investigating the structure of subsurface magnetic field in an emerging flux region on the Sun, using the data obtained by *Hinode*. We derived several statistical quantities characterizing the distribution of surface magnetic field. These quantities were then used to investigate the structure of invisible subsurface magnetic field, such as the twist of field lines, by comparing the observations with numerical simulations. The result suggests a possible way of investigating the structure of subsurface magnetic field by using observational information on surface magnetic field.

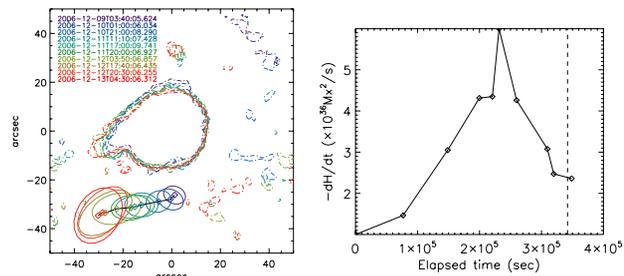
We used *Hinode*'s observational data of photospheric magnetic field to study the evolution of a flare-producing active region. *Hinode* can provide the temporal development of the photospheric field as a vector quantity, which shows that magnetic shear initially developed in an active region (AR10930), then it decreased toward the onset of an X-class flare. The magnetic helicity in this active region first increased rapidly, while it became saturated in the late phase. We explained the physical origin of these features in terms of the emergence of a twisted flux tube into the solar atmosphere. We showed how helicity evolution relates to the occurrence of a flare, suggesting that the helicity evolution obtained from measurements of the photospheric field can be used for flare prediction [1, 2, 3].



**Figure 1:** Emergence of a partially split flux tube. The horizontal plane indicates the photosphere.



**Figure 2:** Scattered diagrams of positive (triangles) and negative (diamonds) polarity regions. Left: Data from *Hinode*. Middle: Simulation of a weakly twisted flux tube. Right: Simulation of a highly twisted flux tube. Note that the observational result has mirror symmetry with the simulations.



**Figure 3:** Left: Relative sunspot motions obtained from an elliptic approximation applied to the boundary of the travelling sunspot. Right: Time variation of the injection rate of magnetic helicity. The vertical dashed line indicates the occurrence time of a flare.

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# Young Brown Dwarfs in the Core of the W3 Main Star-Forming Region

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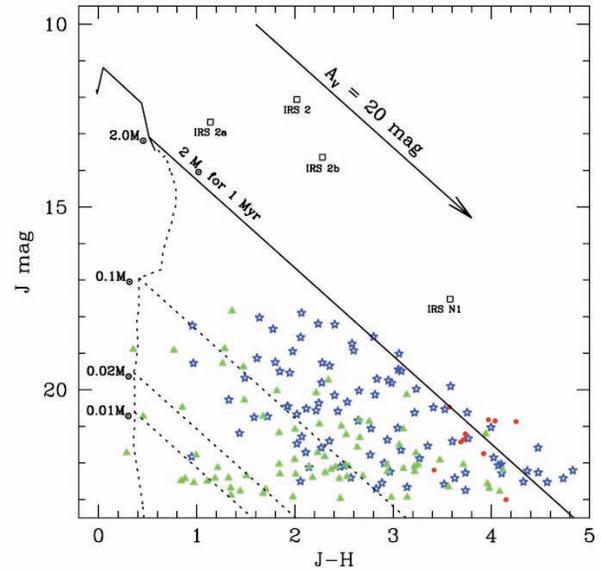
We present the results of deep and high-resolution (FWHM  $\sim 0.35''$ ) *JHK* near-infrared (NIR) observations with the Subaru telescope, to search for very low mass young stellar objects in the W3 Main star-forming region. The NIR survey covers an area of  $\sim 2.6$  arcmin<sup>2</sup> with  $10\sigma$  limiting magnitude exceeding 20 mag in the *JHK* bands. The survey is sensitive enough to provide unprecedented details in W3 IRS 5 and IRS 3a regions and reveals a census of the stellar population down to objects below the hydrogen-burning limit.

We construct *JHK* color-color and *J-H/J* and *H-K/K* color-magnitude diagrams to identify very low luminosity young stellar objects and to estimate their masses (see for example Figure 1). Based on these color-color and color-magnitude diagrams, we identified a rich population of embedded YSO candidates with infrared excesses (Class I and Class II), associated with the W3 Main region. A large number of red sources ( $H-K > 2$ ) have also been detected around W3 Main, which are arranged from the northwest toward the southeast regions. Most of these are concentrated around W3 IRS 5. We argue that these red stars are most probably pre-main-sequence (PMS) stars with intrinsic color excesses.

We find that the slope of the *K*-band luminosity function of W3 Main is lower than the typical values reported for young embedded clusters. Based on the comparison between theoretical evolutionary models of very low-mass PMS objects with the observed color-magnitude diagram, we find there exists a substantial substellar population in the observed region.

The mass function does not show the presence of cutoff and sharp turnover around the substellar limit, at least at the hydrogen-burning limit. Furthermore, the mass function slope indicates that the number ratio of young brown dwarfs and hydrogen-burning stars in the W3 Main is probably higher than those in Trapezium and IC 348.

The presence of mass segregation, in the sense that relatively massive YSOs lie near the cluster center, is seen. The estimated dynamical evolution time indicates that the observed mass segregation in the W3 Main may be the imprint of the star formation process [4].



**Figure 1:** Color-magnitude diagram for the YSO candidates in W3 Main. Class II candidates are indicated by blue stars, filled green triangles represent Class I candidates, and the filled red circles show red sources with  $H-K > 2$  with *J* band counterparts. Known IRS sources are shown with boxes. The solid and dotted curves denote the loci of 1 Myr old PMS stars; derived from the models of [1], [2] and [3].

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# Exoplanet Detection with Polarimetry

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Since the discovery of planets orbiting around normal stars other than the Sun, exoplanet detection is one of the hottest topics in astronomy. Most (>80%) of the exoplanet candidates have been detected via the radial velocity method. The transit method is detecting many candidates (especially with the Kepler and the Corot missions). However, the other successful methods are still limited. Therefore, an independent approach is important to reveal various physical parameters of exoplanets.

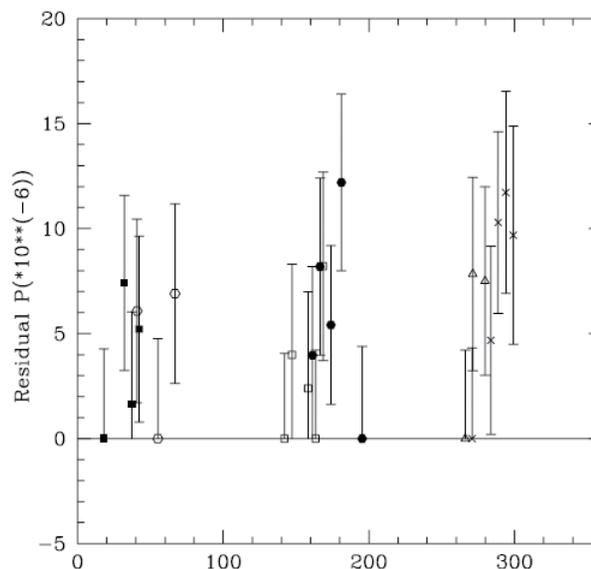
Aged planets shine by reflected light at optical and near-IR wavelengths. Thus, exoplanet detection with reflected light is a promising technique [1, 2].

In particular, high precision polarimetry is considered to be direct detection if the weakly polarized light is extracted from the highly diluted (unpolarized) light from the central star. Close-in planets are expected show as much as a few  $10^{-6}$  polarization. In order to achieve this, a high precision ( $<10^{-6}$ ) polarimeter is necessary. The Planetpol is an optical polarimeter composed of PEM (photoelastic modulator) and APD (avalanche photodiode) that can attain a polarization accuracy of  $<10^{-6}$  [3].

The PlanetPol observations were made during a Saharan dust event over the La Palma observatory that show excess linear polarization in the horizontal direction due to the passage of the starlight through the dust [4]. The polarization reached a maximum value of  $4.8 \times 10^{-5}$ . These results imply that some fraction of the dust grain population aligns with a preferred orientation. We suggest a possible mechanism for vertical orientation resulting from the electric field in the atmosphere.

The PlanetPol observations were also made for 55 Cnc and  $\tau$  Boo in an attempt to detect the partially polarized reflected light from the planets orbiting these two stars [5]. The polarization of this system is very stable, showing no sign of the periodic variations that would be expected if a short-period planet were detected. The measured standard deviation of the nightly averaged Stokes Q/I and U/I parameters is  $2.2 \times 10^{-6}$  (Figure 1). These results contrast markedly with the recent claim of detection of a periodic polarization signal from HD189733 with amplitude  $2 \times 10^{-4}$ , attributed to the planet HD189733b [6]. Our multiple scattering model atmosphere calculations indicate that the large amplitude periodic polarization signal from the HD189733 system cannot be explained in terms of reflected light from the planet HD189733b. If the observations are confirmed it would be important to consider the possible contribution of star spots to the polarization of the system, given that HD189733 is an active star with much larger photometric

variations than  $\tau$  Boo or 55 Cnc.



**Figure 1:** Polarization vs. orbital angle of 55 Cnc e (period= 2.817d). 180 degrees correspond to the minimum illumination. Polarization due to telescope and instrument is subtracted. Different symbols correspond different nights.

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# Achievement of the CLIO target sensitivity at room temperature

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CLIO (Cryogenic Laser Interferometer Observatory) is the first cryogenic interferometric gravitational wave detector in the world. It has 100-m baseline length and is located in an underground mine at Kamioka. It is a prototype interferometer for the future LCGT project as well as TAMA300 detector in Mitaka campus. To demonstrate advantages of the cryogenic interferometer and the Kamioka mine as a quiet site is the purpose.

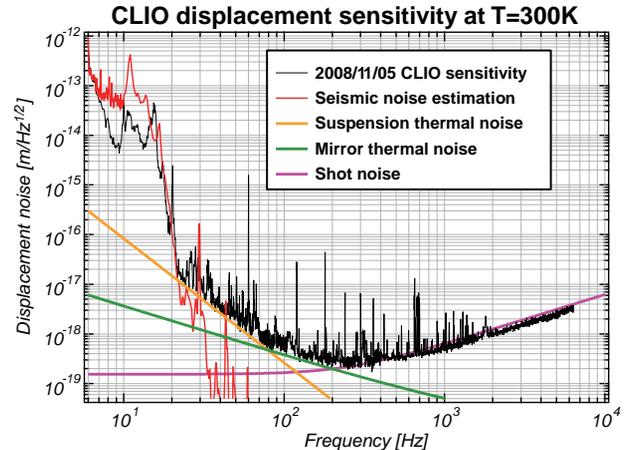
The project started to excavate the center room and two tunnels in 2002. The vacuum and the cryogenic systems were prepared by 2005. After the installation of the interferometer optics and the suspension systems, the operation of CLIO started in 2006. Test of the cryogenic systems and the interferometer operation with cryogenic mirrors were demonstrated in 2006.

After the demonstration of cryogenic operations, our activities have been focused on the noise investigation at room temperature. At the time of 2007, the strain sensitivity below 20 Hz is comparable to that of the LIGO detector, which is a couple of 4-km detectors in United State. Because the baseline length is forty times shorter than that of LIGO, it means that CLIO achieved forty times smaller noise level [1]. This is due to a big advantage of underground site. In February 2007, a 57-hours observation data was obtained to search for gravitational wave signal at 22.38 Hz from the pulsar J0835-4510 (Vela pulsar) [2].

To demonstrate the reduction of the noise level by mirror cooling, the detector noise should be limited by thermal noises at room temperature. In the summer of 2008, the sensitivity of CLIO at room temperature was improved in a wide frequency range above 20 Hz. This improvement was obtained by elimination of thermal noise from dissipation caused by eddy current induced to the coil support made of aluminum, as well as tuning of beam spot positions on the mirror surfaces, which indicates possible noise coupling from any kind of mirror rotation noise. Current best floor sensitivity is  $2.5 \times 10^{-19} \text{ m}/\sqrt{\text{Hz}}$  at 250 Hz as shown in Figure 1. We estimated the fundamental noises of the CLIO as follows. Seismic noise was estimated by the ground motion and transfer function of the vibration isolation systems. The

suspension thermal noise was estimated from the quality factor of the violin mode resonances. The mirror thermal noise at room temperature was supposed to be dominated by thermoelastic damping. Shot noise was estimated from incident optical power to the photodetector and its response. From these estimations, we concluded that the noise level is enough close to the fundamental limit at room temperature and the cooling of the mirrors is necessary for further noise investigations.

The next step is to realize noise reduction by mirror cooling. Now we are preparing noise hunting with the cryogenic mirrors including works to refine the cryogenic suspensions and the interferometer systems.



**Figure 1:** Current displacement sensitivity of CLIO (black). Noise estimations of seismic noise (red), suspension thermal noise (yellow), mirror thermal noise (green), and shot noise (purple) are also shown.

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# Detection of water ice grains on the surface of the circumstellar disk around HD 142527

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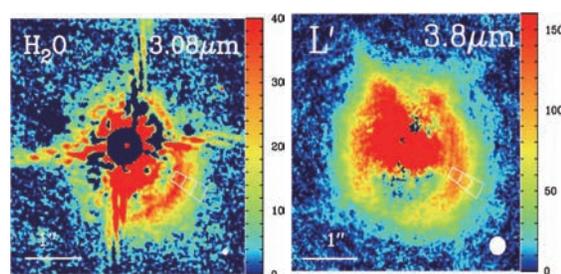
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Water ice grains theoretically play a number of important roles in protoplanetary and debris disks. Ice enhances the solid material in the cold outer part of a protoplanetary disk, which promotes the formation of cores of gaseous planets [1]. [2] show that ice absorption can also be found in the light scattered by icy grains and that multi-wavelength imaging in NIR wavebands, including in the H<sub>2</sub>O band at 3.1 μm, is a useful tool to constrain the ice distribution in the disk.

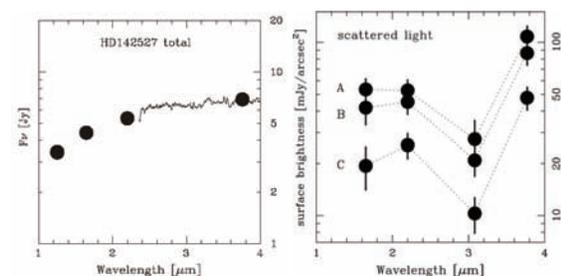
We have applied this method to the circumstellar disk around a Herbig Ae star, HD142527, which is an ideal target for imaging in the H<sub>2</sub>O band for a number of reasons. Coronagraphic imaging observations for the HD142527 using a L' band (3.77 μm) and a narrow-band H<sub>2</sub>O ice filter (3.08 μm) were performed using the Coronagraphic Imager with Adaptive Optics (CIAO) on the Subaru Telescope. Figure 1 shows the PSF-subtracted images of the circumstellar disk around HD 142527 in the H<sub>2</sub>O ice filter with the 0.6" mask and in the L' band without a mask, respectively. The eastern and western arcs discovered by [3] in the H and K bands were also detected in the H<sub>2</sub>O and L' band. The spectrum for the total flux densities is also presented for comparison in the left panel of Figure 2, which shows few or no features in the H<sub>2</sub>O band. The contribution from the scattered light to the total flux density in the H and K bands is estimated to be 3.1–3.2% [3], indicating that most of the flux density in the total spectrum is attributed to the emission from the star itself and thermal emission from hot grains in the vicinity of the star where icy grains cannot survive. Therefore, the absence of a 3 μm ice feature is a natural consequence. In contrast, the spectra of the scattered light show a clear reduction in the H<sub>2</sub>O ice band (3.08 μm), suggesting the presence of a 3 μm absorption feature. This result can be explained by the water ice grains in the disk surface where the light scattering occurs. Comparing the model predictions [2] with the observed color difference, the size of the scattering grains is ~1 μm or more.

Our results successfully demonstrate that multi-wavelength (K/H<sub>2</sub>O/L' imaging is useful to detect water ice grains in a (face-on) disk. The application of this method for other face-on disks will enhance our knowledge

on the presence and spatial distribution of water ice in circumstellar disks [4].



**Figure 1:** PSF-subtracted 3.08 μm (left) and L' band (right) image of the disk around HD142527. The unit of intensity is mJy arcsec<sup>-2</sup>. North is up; east is to the left. The three enclosed regions (regions A, B, and C from inner to outer regions) on the southwest structure are the positions from which the spectra shown in Figure 2 are extracted.



**Figure 2:** Left panel shows NIR photometry and a spectrum of the entire HD 142527 system, which is used as the representative of the illuminating source spectrum. Right panel represents the spectra of the scattered light from regions A, B, and C shown in Figure 1. The dip in 3 μm, probably due to H<sub>2</sub>O ice grains, is present in all the scattered light spectra, while it is not seen in the total spectrum, which is dominated by thermal emission from hot dust near the central star.

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# Small-JASMINE (Small Satellite with 30cm-class telescope)

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JASMINE Working Group

JASMINE is the acronym of the Japan Astrometry Satellite Mission for INfrared Exploration. The main objective of JASMINE is to study the fundamental structure and evolution of the Milky Way bulge components. In order to accomplish these objectives, JASMINE will measure trigonometric parallaxes, positions and proper motions of stars in the Galactic bulge with a precision of 10 microarcsec. Before the launch of JASMINE, we plan a launch of Small-JASMINE mission which is a downsized version of JASMINE.

The primary mirror for the telescope of Small-JASMINE has a diameter of 30cm with a focal length of about 5m. The wavelength to be observed is around  $2\mu\text{m}$  with the HgCdTe detectors. A target accuracy of the mission is the same with JASMINE,  $10\mu\text{as}$ . The size of the field of view is about  $0.87\text{deg} \times 0.87\text{deg}$  by using  $2 \times 2$  detectors on the focal plane. The telescope is designed to have only one field of view, which is different from the designs of other astrometric satellites (but the same with JASMINE). JASMINE will observe overlapping fields without gaps to survey a total area of about a few square degrees around the Galactic bulge. Accordingly we make a “large frame” of a square degrees by linking the small frames using stars in overlapping regions. Small-JASMINE will observe the region of Galactic bulge repeatedly during the mission life.

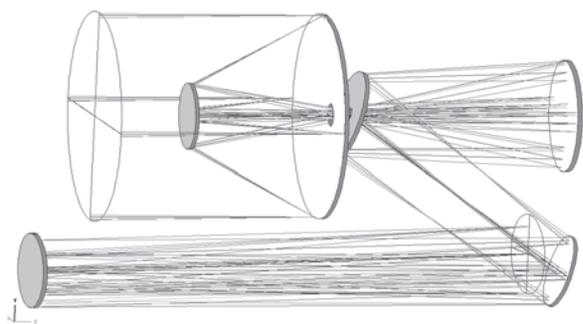


Figure 1: Overview of JASMINE Optics.

Optical Design	3 mirrors (Improved Korsch)
Wave length	$1.5\mu\text{m} \leq \lambda \leq 2.5\mu\text{m}$
Aperture	30 cm
Focal length	4.86 m
Field size	$0.87^\circ \times 0.87^\circ$
Size of Focal Plane	6 cm $\times$ 6 cm
Detector	HgCdTe
Number of pixels	2K $\times$ 2K
Number of detectors	4 (2 $\times$ 2)
Pixel size	15 $\mu\text{m}$
Pixel on sky	138 mas

Table 1: Summary of JASMINE Optics.

Number of Small Frames	4 $\times$ 4
Observing Magnitude	K=11 mag
Integration time	3sec
Time for making a Large Frame	0.25 h
Mission Life	1.5 year

Table 2: Parameters.

Pointing Stability	280 mas/3 sec
Thermal Stability	0.4 K/0.25 h
Orbit	600 Km Earth Orbit
Pointing Accuracy	0.1 deg
Data Transfer Rate	2 Mbps

Table 3: Requirements.

## Reference

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# Development of Very Small Astrometry Satellite: Nano-JASMINE

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JASMINE Project Office of National Astronomical Observatory of Japan, Kyoto University, University of Tokyo and Tokyo University of Marine Science and Technology develop a very small astrometry satellite: Nano- JASMINE. Nano-JASMINE is mounted a 5 cm effective diameter telescope and aims to measure positions and motions of stars with the accuracy of a few milliarcseconds in  $z$ -band ( $\lambda \sim 0.9 \mu\text{m}$ ). The satellite weighs about 25 kg and is developed aiming at the launch in 2010 [1]. JASMINE Project Office contributes to develop a small reflecting telescope and a driver circuit for a CCD detector. In 2008, we developed and evaluated these instruments and conducted the integration test and environmental test. This paper reports the results of development.

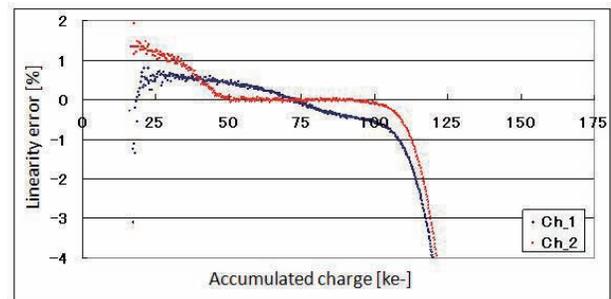
First, the small reflecting telescope was made of aluminum alloy because of the suppression of effects of strain induced by thermal changes. The reflecting surfaces were polished with a single-crystal diamond cutting tools and deposited with Cr and Au vapor. Developed telescope weighs 1.7 kg and occupies a volume of only  $17 \times 12 \times 12$  cm [2]. Wavefront errors of developed telescope were measured by an interferometer. As a result, the RSS wavefront error is  $\lambda/14$  ( $\lambda = 800$  nm). Then it is confirmed that developed telescope can achieve a diffractionlimited performance.

Second, new driver circuit for a CCD detector was developed. Developed circuit can be operated in TDI (Time Delay Integration) mode. In TDI mode, the rate of charge transfer is synchronized with the scan velocity of the spinning satellite. Developed circuit was evaluated through the measurement of system gain, dark current and CTE. Fig. 1 shows the relationship between accumulated charges and linearity error of CCD.

Next, Structure and Thermal Model (STM) of satellite was developed for the thermal test and the vibration test. Using STM, mechanical interfaces of satellite were confirmed and thermal vacuum test was conducted. Fig. 2 shows the overview of STM before the thermal vacuum test. Moreover, electric integration test was conducted and the electric interface, each function and power consumption were confirmed. These results were reflected to engineering model.

Over the next year, engineering model is to be developed and examined, and flight model is to be

developed. About the launch, NAOJ, University of Tokyo, Alcantara Cyclone Space and SDO Yuzhnoye agreed the memorandum of understanding (MOU) to launch Nano-JASMINE by Cyclone-4 rocket in the Federal Republic of Brazil.



**Figure 1:** The relationship between accumulated charges and linearity error of CCD. The linearity error is increased at 115ke- where accumulated charge is saturated.



**Figure 2:** Nano-JASMINE STM at the NAOJ space chamber just before the thermal vacuum test. Test heaters which control outer temperature are equipped around STM.

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# Neutrino-Nucleus Reaction Cross Sections for Light Element Synthesis in Supernova Explosions

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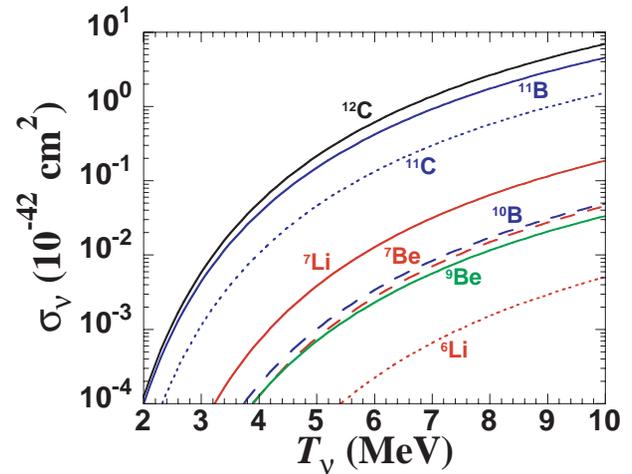
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HARTMANN, Dieter, H.  
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Neutrino nucleosynthesis (the  $\nu$ -process) in supernovae is one of main synthesis processes for rare light elements such as Li and B. In order to evaluate the light element yields, cross sections of neutrino-nucleus reactions should be precisely calculated. Recently, a new shell model Hamiltonian (the SFO Hamiltonian) reproduces well many characteristics of  $p$ -shell nuclei [1]. We evaluate cross sections of  $\nu$ - $^{12}\text{C}$  and  $\nu$ - $^4\text{He}$  reactions using the SFO and WBP Hamiltonians, respectively. Then, we discuss light element synthesis through the  $\nu$ -process in supernovae using the new cross sections [2].

Cross sections of  $\nu$ - $^{12}\text{C}$  reactions are calculated using the SFO Hamiltonian [1]. Both of neutral- and charged-current reactions are taken into account. Branching ratios for multiple particle emissions are evaluated using Hauser-Feshbach calculations. Cross sections of  $\nu$ - $^4\text{He}$  reactions are also calculated using the WBP Hamiltonian. Figure 1 shows the cross sections of neutral-current  $\nu$ - $^{12}\text{C}$  reactions as a function of neutrino temperature. Neutrino energy distribution is assumed to obey Fermi distributions with zero chemical potential. Only large branches are shown in this figure. Tables of the cross sections for  $\nu$ - $^{12}\text{C}$  and  $\nu$ - $^4\text{He}$  are listed in [2].

We evaluate the yields of light elements in a  $16.2 M_{\odot}$  supernova with the new cross sections. The yields of  $^7\text{Li}$  and  $^{11}\text{B}$  are  $2.67 \times 10.7 M_{\odot}$  and  $7.14 \times 10.7 M_{\odot}$ , respectively. These yields are larger than the corresponding ones with conventional cross sections by a factor of 1.1. The yields of  $^6\text{Li}$ ,  $^9\text{Be}$ , and a radioactive nucleus  $^{10}\text{Be}$  are a level of  $10^{-11} M_{\odot}$  and increase compared to the ones with the conventional rates. The temperature range of  $\nu_{\mu,\tau}$  and  $\bar{\nu}_{\mu,\tau}$  in supernovae appropriate for the Galactic chemical evolution of  $^{11}\text{B}$  is reevaluated to be 4.3–6.5 MeV. This range is slightly shifted to lower energy compared to the case with the conventional rates. The increase in the  $^7\text{Li}$  and  $^{11}\text{B}$  yields due to neutrino oscillations is also shown in [2]. This result is important in the determination of unknown neutrino oscillation parameters.



**Figure 1:** Calculation neutrino- $^{12}\text{C}$  reaction cross sections as a function of neutrino temperature  $T_{\nu}$  for the SFO Hamiltonian.

## References

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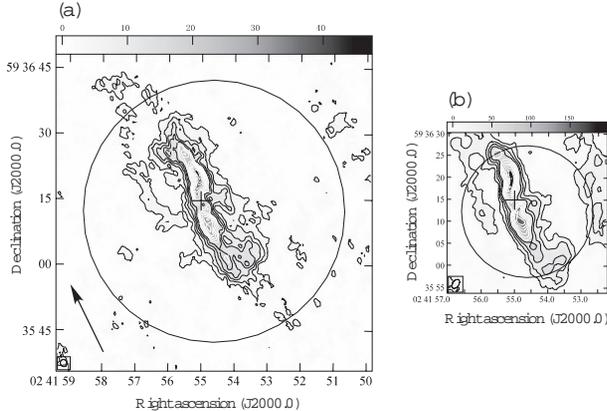
# Central Spiral Structure of Molecular Gas in Maffei 2

KUNO, Nario, NAKANISHI, Koichiro  
(NAOJ)

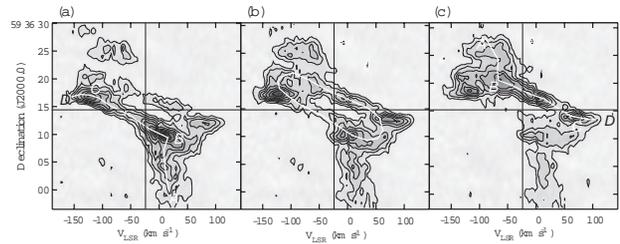
SORAI, Kazuo  
(Hokkaido University)

SHIBATSUKA, Toshihito  
(University of Tokyo)

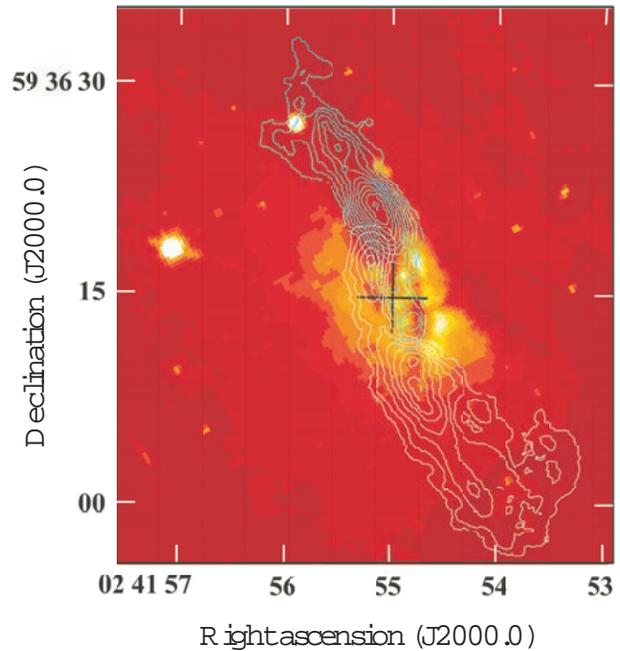
We made CO(1–0), CO(2–1), CS(2–1) lines and 103 GHz continuum observations of the nearby barred spiral galaxy Maffei 2 with Nobeyama Millimeter Array (NMA) [1]. The CO integrated intensity maps show two peaks in the central region (Figure 1). By investigating the velocity structure, however, we found the spiral structure of molecular gas in the central region (Figures 2, 3). The offset ridges of molecular gas along the leading side of the bar continue to the spiral structure embedded in the weak oval structure which is regarded as an  $x_2$  orbit. The size of these structure is less than  $R \sim 100$  pc. The mass within  $R \sim 35$  pc is estimated to be  $2 \times 10^8 M_\odot$  from the rotation velocity of molecular gas. The high mass concentration is consistent with theoretical predictions concerning the creation of such a nuclear spiral structure [2]. A comparison with tracers of dense gas and star-forming region suggests that the dense molecular gas traced by the CS(2–1) line is formed at the crossing points of the  $x_1$  and  $x_2$  orbits and that the star-forming region appears downstream of the peak of the dense gas.



**Figure 1:** (a) CO(1–0) and (b) CO(2–1) integrated intensity maps. The cross indicates the galactic center. The circles indicate the field of view of NMA. The arrow indicates the position angle of the major axis of Maffei 2.



**Figure 2:** P-V diagrams of CO(2–1) along major axis of Maffei 2 and parallel to the major axis. (a) 0.9'' eastern side of the major axis. (b) along the major axis. (c) 0.9'' western side of the major axis. Two parallel components that correspond to the spiral structure are seen.



**Figure 3:** Contour Integrated intensity map of CO(1–0) of the two components seen in Fig. 2. Color) The HST image (F814W).

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# Near-Infrared Imaging Polarimetry of M42: Aperture Polarimetry of Point-like Sources

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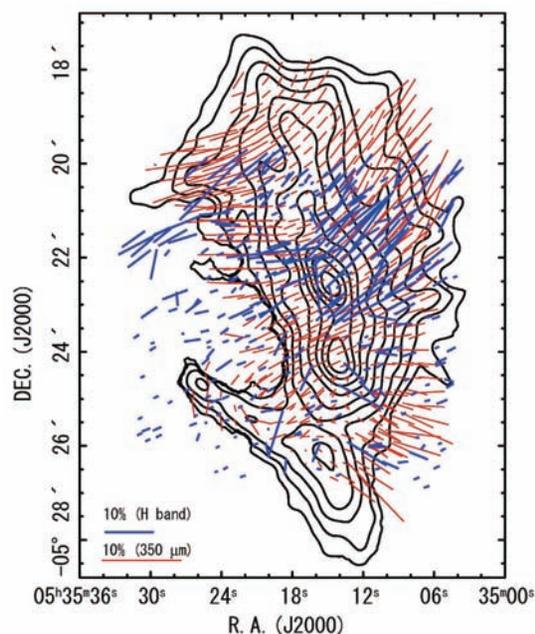
HOUGH, James H., LUCAS, Philip W.  
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We have conducted aperture polarimetry of  $\sim 500$  stars of the Orion Nebula Cluster (ONC) in M42 based on our wide-field ( $\sim 8' \times \sim 8'$ ) *JHKs* band polarimetry (Figure 1) [1]. Most of the near-infrared (NIR) polarizations are dichroic, with position angles of polarization agreeing, both globally and locally, with previous far-infrared (FIR) and submillimeter observations [2], having taken into account the  $90^\circ$  difference in angles between dichroic absorption and emission. This is consistent with the idea that both NIR dichroic polarizations and FIR/submillimeter thermal polarizations trace the magnetic fields in the OMC-1 region. The magnetic fields inferred from these observations show a pinch at scales less than  $0.5$  pc with a centroid near IRC2. The hourglass-shaped magnetic field pattern is explained by the models in which the magnetic field lines are dragged along with the contracting gas and then wound up by rotation in a disk. The highly polarized region to the northwest of IRC2 and the low-polarized region near the bright bar are also common among NIR and FIR/submillimeter data, although a few regions of discrepancy exist.

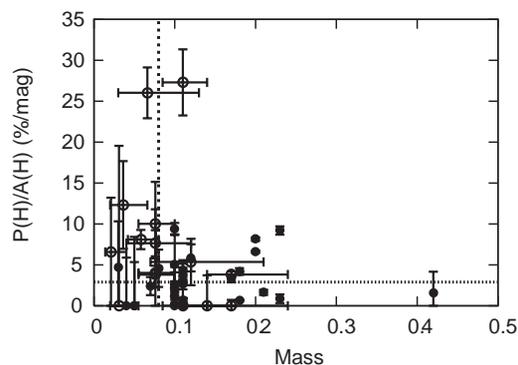
We have also discerned 51 possible highly polarized sources whose polarizations are more likely to be intrinsic rather than dichroic. Their polarization efficiencies ( $P(H)/A(H)$ ) are too large to be explained by the interstellar polarization. These include nine young brown dwarfs that suggest a higher polarization efficiency (see Figure 2), which may present geometrical evidence for (unresolved) circumstellar structures around young brown dwarfs.

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- [4] Riddick, F. C., et al.: 2007, *MNRAS*, **381**, 1077.



**Figure 1:** *H* band (shown in blue) polarization vectors and  $350\ \mu\text{m}$  (shown in red, rotated by  $90^\circ$ ) polarization vectors overlaid on the  $350\ \mu\text{m}$  contours. 10% polarization vectors for both wavelengths are shown at the bottom-left.



**Figure 2:**  $P(H)/A(H)$  for the 28 sources cross-identified with Slesnick et al. (2004) [3] (filled circle) and Riddick et al. (2007) [4] (open circle). The horizontal dotted line indicates  $P(H)/A(H) = 2.9$ , which corresponds to the assumed maximum value for interstellar polarization. The vertical dotted line indicates  $0.08 M_\odot$ , which corresponds to the brown dwarf mass limit.

# Infrared Imaging Polarimetry of S106 in Cygnus

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SIRPOL [1] is the polarization mode of the three-band simultaneous camera SIRIUS [2] on the 1.4-m telescope, IRSF, at the South African Astronomical Observatory (SAAO). Now, SIRPOL has two polarization modes, linear polarization and circular polarization modes.

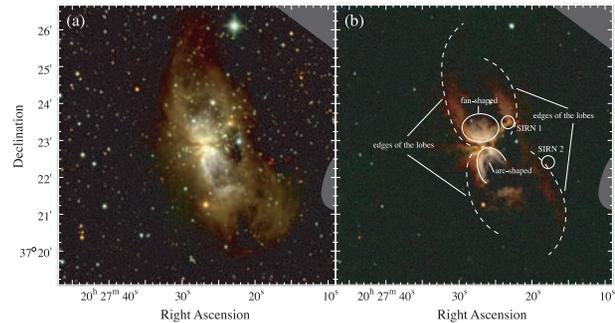
The polarization happens when dust reflects the light from the illuminating stars or absorb the light from the background stars. Thus, the imaging polarimetry can reveal the distribution of the reflection nebula and the characteristic of the dust in the region, and discover the illuminating stars. In addition, we can reveal the structure of the magnetic fields in the observing area using polarimetry of the background stars because dust grains are aligned by the magnetic fields. In particular, near infrared polarimetry makes it possible to reveal the structure of reflection nebulae and magnetic fields in star-forming regions because infrared observations can detect the objects in high dusty regions.

However, in spite of its usefulness for understanding the physical characteristics, “wide-field” infrared polarimetry has been scarcely carried out because of the lack of an appropriate polarimeter. Thus, the study with polarimetry toward even nearby star-forming regions was very limited.

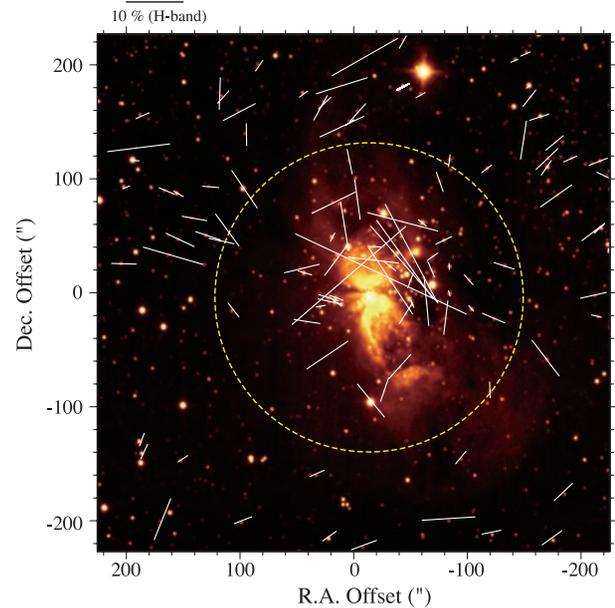
Figure 1 shows  $JHK_s$  composite images of the S106 cluster forming region. Fig. 1(a) is the image in intensity and Fig. 1(b) is the same image but in only polarized intensity. We could easily find the clear difference between these two images. The distribution of the polarized intensity has the characteristic structures on the eastern and western edges of the S106 nebula and around the exciting star, IRS 4. Compared with the distribution of the molecular gas, the ridges of the polarized intensity are corresponding to the boundary of the molecular cloud, and we suggest that the polarized components originate in the light reflected by the dust grain in the dense gas. In addition, we found two small infrared reflection nebulae (SIRNe). Two SIRNe in Fig. 1(b) have the polarized component with a small size, and in particular, SIRN 1 is associated with an early B protostar.

Figure 2 shows the aperture polarization vector map of point-like sources in the  $H$ -band. We found the difference between the polarization angle and the axis of the S106 nebula because the overall polarization angle distribution of the area outside the cluster is  $\sim 120^\circ$ . However, corresponding to the results of sub-millimeter observations, we revealed that the magnetic field was drawn into the center of the cluster. We suggested that

this feature was caused by the clear difference between the rotation axis and the position angle of the magnetic field.



**Figure 1:** NIR three-color composite images of the S106 nebula in intensity (a) and in polarized intensity (b).



**Figure 2:** Polarization vector of each point-like source on the  $H$ -band image.

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# Discovery of the Multiple-Core Systems toward Cluster-Forming Regions

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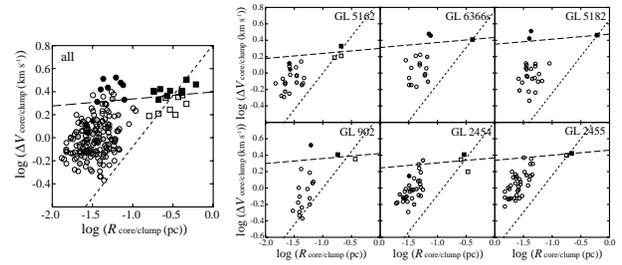
Most stars in the galactic disk are formed as a cluster in dense gas within giant molecular clouds as members of clusters [1]. The cluster consists of stars of various masses, and massive stars exist at the centers of clusters. Thus, clusters must play a critical role in the origins of some of the most fundamental properties of the galactic stellar population. Recently, observational results suggested that clusters are formed in the clumps with a cluster size scale [2] and that a few hot cores (0.03 pc) exist in cluster-forming regions. On the other hand, many cores must be formed in one clump to create a cluster with a large number of stars. Therefore, we carried out to reveal the inner structure of the clump with cluster formation.

We identified 171 cores in six clumps with cluster formation using Nobeyama Millimeter Array. The left panel of Figure 1 shows the relationship between the line width and the radius of the cores and 18 clumps identified by [2]. Although the line width of the clumps is in a narrow range, the line width of the cores is in a much wider range. Such features are seen even in the individual regions in the right panels of Figure 1. This indicates that the degree of the dissipation of the turbulent motion varies even within a single clump.

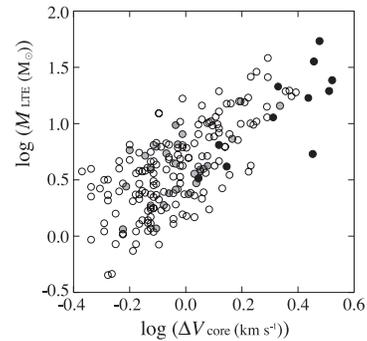
Figure 2 shows that the mass of the cores increases in proportion to the square of the line width, although the dispersion of the plot is large. Considering that the internal kinetic motion basically decreases with time, from this tendency, dense gas with a large kinetic motion in a clump seems to be necessary to form massive cores.

The stellar number density of the cluster is mentioned as an important characteristic of a cluster. The stellar number density of the cluster is generally much higher than that of typical low-mass star-forming regions. The recent results found a good relationship between the number density of YSOs with the clump and the  $H_2$  density of the clump [2]. Figure 3 shows the relationship between the number density of the cores in the clump and the  $H_2$  density of the clump. A good correlation exists between the core number density and the  $H_2$  density of the clump. This indicates that the stellar number density in the clump is controlled the core-formation mechanism.

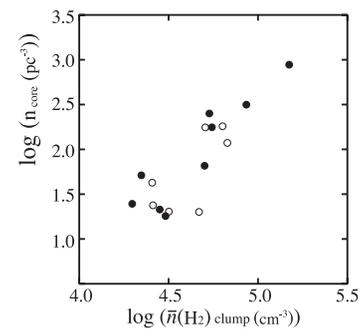
We determined that the mechanism of cluster formation would strongly depend on the mechanism of core formation and that the mechanism of core formation would be controlled by both the  $H_2$  density structure and the kinetic motion in a clump. Moreover, the characteristics of the cluster would be determined by the distribution of the physical parameters of cores in the clump [3].



**Figure 1:** The correlation between the radius and the line width of cores/clumps. The circles and squares indicate cores and clumps, respectively. The white, gray, and black symbols indicate the objects without, with YSOs, and with massive stars, respectively. The long- and short-dashed lines indicate the relationships of  $\Delta V \sim R^{0.06}$  and  $\sim R^{0.89}$ , respectively.



**Figure 2:** The LTE mass plotted against the line width of the cores. The symbols are the same as in Figure 1.



**Figure 3:** The number density of the cores in each clump plotted against the average  $H_2$  density of the clump.

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# Four-Quadrant Phase Mask Coronagraph with a Jacquinot-Lyot Stop

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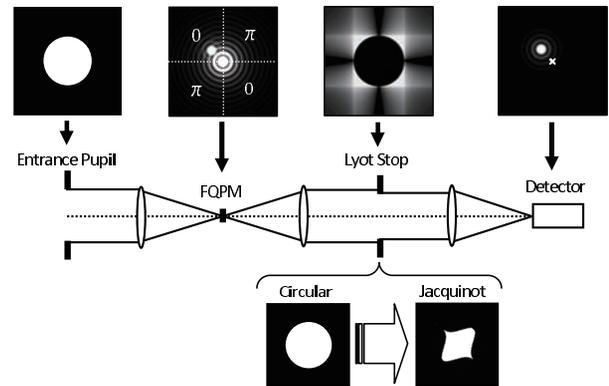
UEMURA, Ryosuke, BABA, Naoshi, SATO, Yoichiro  
(Hokkaido University)

NISHIKAWA, Jun, TAMURA, Motohide  
(NAOJ)

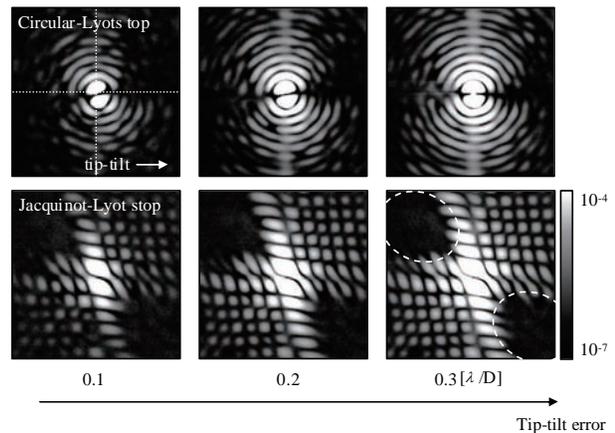
For direct detection of Earth-like extrasolar planets, extremely high dynamic range is required for coronagraphs. A four-quadrant phase-mask (FQPM) coronagraph is one of the most promising methods in that it has a relatively small inner working angle, large discovery space, and high optical throughput with technically simple instrumentation [1]. The FQPM is put on a focal plane to divide a stellar image into four-quadrant regions and provide  $\pi$ -phase difference between the adjacent quadrants (Fig. 1). This causes destructive interference inside the pupil area on a following reimaged pupil plane, where a Lyot-stop is put to block stellar light diffracted outside the pupil. Theoretically, the FQPM coronagraph can realize perfect stellar elimination for point-like stars. However, the FQPM coronagraph is very sensitive to tip-tilt errors and finite stellar angular sizes. Thus residual stellar noise would prevent us from detecting faint companions around partially resolved nearby stars.

To solve this problem, we propose to use a Jacquinot pupil as a Lyot stop of the FQPM coronagraph [2]. The Jacquinot pupil has a Gaussian-like shaped edge, as shown in Figure 1, to suppress a stellar diffraction halo in specified regions of the focal plane (so-called working region) where planetary signals can be detected. We manufactured the Jacquinot pupil by means of a chemical etching process on a nickel substrate (diameter of 1.125 mm and thickness of about  $20\ \mu\text{m}$ ), and conducted laboratory experiments of the FQPM coronagraph. As the FQPM, we used a four-quadrant polarization mask (FQPoM) utilizing polarization interferometry [3].

Figure 2 shows coronagraphic images, acquired with a He-Ne laser ( $\lambda = 633\ \text{nm}$ ), for the circular- and Jacquinot-Lyot stops for tip-tilt errors from 0.1 to 0.3  $\lambda/D$ . Dotted lines and an arrow in the image (top left) show boundaries of the FQPoM and a tip-tilt direction, respectively. As can be seen in the results, the residual halo for the circular-Lyot stop increases rapidly with the tip-tilt error, while that for the Jacquinot-Lyot stop can be well suppressed ( $1.4 \times 10^{-6}$  at  $7.5\ \lambda/D$ ) in the working regions (dashed circles) even for the large tip-tilt error of 0.3  $\lambda/D$ . Consequently, the laboratory experiments demonstrate that the Jacquinot-Lyot stop would enhance the FQPM-coronagraphic performance effectively.



**Figure 1:** A schematic optical setup of a FQPM coronagraph equipped with a Jacquinot-Lyot stop.



**Figure 2:** Results of laboratory experiments of a FQPM coronagraph with circular- and Jacquinot-Lyot stops for various tip-tilt errors.

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# An Eight-Octant Phase-Mask Coronagraph

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For direct detection and characterization of extrasolar planets, light from their parent stars must be strongly suppressed by coronagraphs. We propose an eight-octant phase-mask (EOPM) coronagraph [1], which is analogous to a previously proposed four-quadrant phase-mask (FQPM) one [2] (Fig. 1). It is expected that the EOPM coronagraph has fourth-order sensitivity to tip-tilt errors while the sensitivity of the FQPM one is second-order. This higher order behavior of the EOPM coronagraph will be advantageous for observations of partially resolved nearby stars.

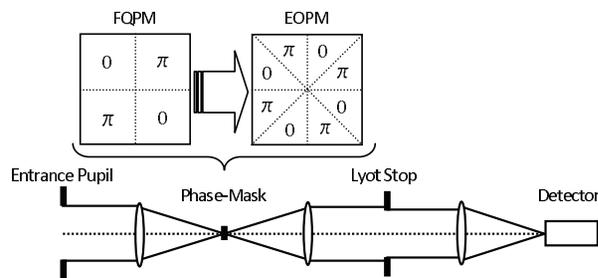


Figure 1: Principles of FQPM and EOPM coronagraphs.

We manufactured an eight-segmented phase-mask utilizing a nematic liquid-crystal (LC) device (Fig. 2). A phase retardation of each segment can be adjusted by an applied voltage to the device via flexible printed circuits (FPC). The LC phase-mask can be switched between FQPM- and EOPM-mode by applying voltages to appropriate segments (e.g., segments L1, L3, R2, and R4 for the EOPM-mode).

We carried out laboratory experiments of the phasemask coronagraphs by using a He-Ne laser ( $\lambda = 633$  nm). In Figure 2, acquired Lyot-stop images for the both coronagraphic modes are also shown together with those of numerical simulations. Figure 3 shows residual intensity of acquired coronagraphic images as a function of a tip-tilt error. The coronagraphic images for  $0.4 \lambda/D$  tip-tilt error are also shown. Note that we subtract a zero-tip-tilt coronagraphic image from each one to eliminate speckle noise in order to evaluate residual intensity due to only the tip-tilt error. As can be seen in the results, the residual intensity can be well suppressed by the EOPM coronagraph because of its fourth-order behavior.

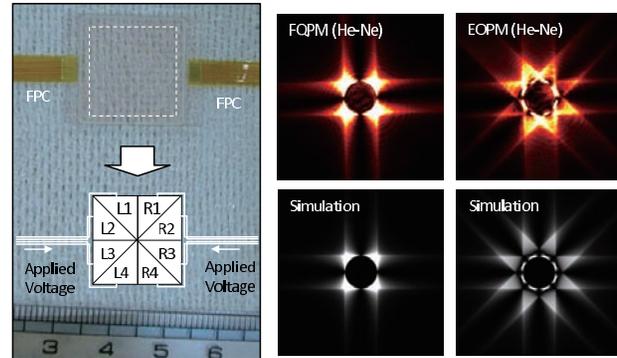


Figure 2: A manufactured liquid-crystal phase-mask, and acquired and numerically simulated Lyot-stop images of FQPM- and EOPM-mode.

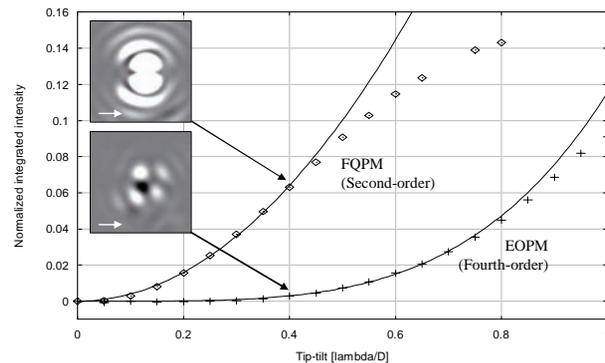


Figure 3: Residual intensity of FQPM and EOPM coronagraphs as a function of a tip-tilt error. Coronagraphic images for  $0.4 \lambda/D$  tip-tilt error are also shown.

## References

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- [2] Rouan, D., et al.: 2000, *PASP*, **112**, 1479.

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# Variation of the Subhalo Abundance in Dark Matter Halos

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MAKINO, Junichiro  
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Moore et al. (1999) and Klypin et al. (1999) calculated the evolution of galaxy-scale dark matter halos using high resolution cosmological  $N$ -body simulations. They found that dark matter halos of the mass comparable to the Local Group contained far too many subhalos compared to the number of known dwarf galaxies in the Local Group. In the case of dark matter halos of the size of typical clusters of galaxies, the theoretical prediction and observation agreed pretty well. The theoretical prediction for galaxy-sized and cluster-sized halos are similar, and observations of clusters of galaxies and that of Local Group are very different [1, 2]. This discrepancy is now called “the missing dwarf problem”.

We analyzed the statistics of subhalo abundance of galaxy-sized and giant-galaxy-sized halos formed in a high-resolution cosmological simulation of a 46.5 Mpc cube with the uniform mass resolution of  $10^6 M_\odot$ . We analyzed all halos with mass more than  $1.5 \times 10^{12} M_\odot$  formed in this simulation box. The total number of halos was 125. We found that the subhalo abundance shows large halo-to-halo variations. The results of recent ultra-high-resolution runs fall within the variation of our samples. We found that the concentration parameter and the radius at the moment of the maximum expansion show fairly tight correlations with the subhalo abundance. This correlation suggests that the variation of the subhalo abundance is at least partly due to the difference in the formation history. Halos formed earlier have smaller number of subhalos at present [3, 4].

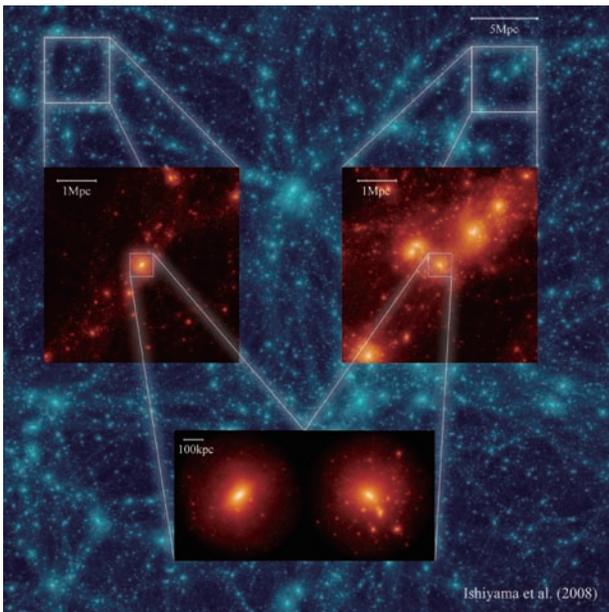


Figure 1: Snapshot ( $z = 0$ ).

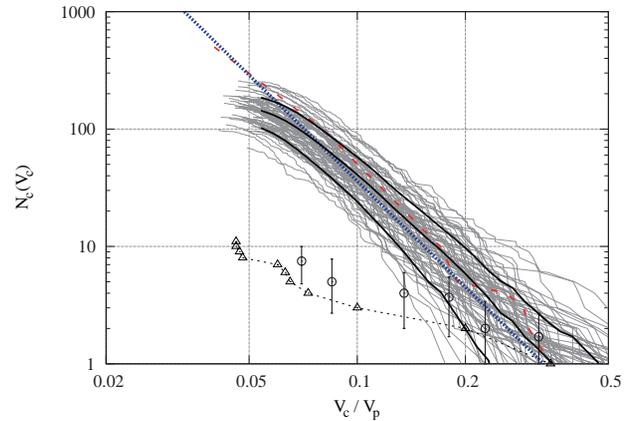


Figure 2: Cumulative numbers of subhalos as a function of their maximum rotation velocities  $V_c$  normalized by those of the parent halos  $V_p$  for 68 galaxy-sized halos with  $1.5 \times 10^{12} M_\odot \leq M < 3 \times 10^{12} M_\odot$ . Three thick solid curves show the average (middle) and  $\pm 1\sigma$  values (top and bottom). Red line are the result of [2] for a galaxy-sized halo. The thin dashed curve with open triangles denotes the number of dwarf galaxies in our galaxy. The open circles with error bars show the number of dwarf galaxies in the Local Group.

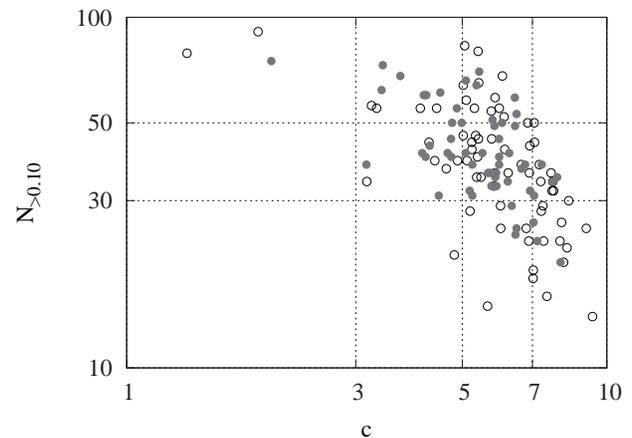


Figure 3: Dependence on the concentration parameter  $c$  of the subhalo abundance. More centrally concentrated halos are generally formed earlier.

## References

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- [2] Moore, B., et al.: 1999, *ApJL*, **524**, L19.
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- [4] Ishiyama, T., et al.: 2009, *ApJ*, **696**, 2115.

# Tycho's Supernova is of Type Ia ~Super explosion in 16th century caught by Subaru in 21st century~

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(NAOJ)

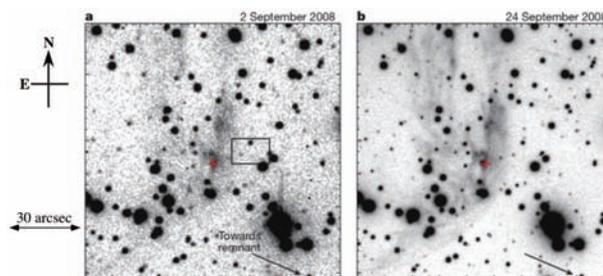
KRAUSE, Oliver, BIRKMANN, Stephan M., GOTO, Miwa  
(Max-Planck-Institut for Astronomy)

TANAKA, Masaomi, NOMOTO, Ken'ichi  
(University of Tokyo)

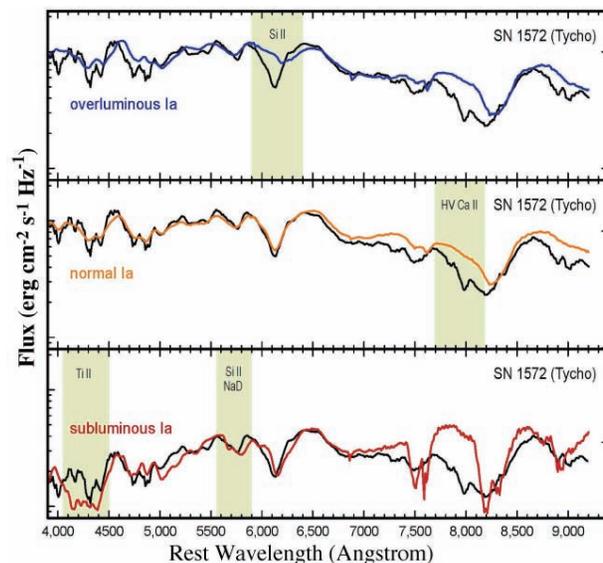
On the evening of 11 November 1572, Tycho Brahe, a Danish astronomer observed a bright star in the constellation Cassiopeia, outshining even Venus. He studied the brightness and color of the “new star” until March 1574 when it faded from view. On the basis of historical records of the light curve and colour evolution, it has been interpreted as a type Ia supernova of either a normal or somewhat overluminous type or of a subluminous type. Core-collapse supernovae of type Ib or II-L have also been suggested to be compatible with the light curve, however the determination of the exact supernova type has not been possible without spectroscopic information.

We started to observe the identified light-echo fields of SN 1572 using the 2.2-m and 3.5-m telescopes at the Calar Alto Astronomical Observatory, on 23 August 2008 and 2 September 2008, and found a bright and extended nebulosity with a peak surface brightness of  $R = 23.6 \pm 0.2$  mag arcsec<sup>-2</sup> (Figure 1a).

The region was re-observed using the Faint Object Camera And Spectrograph (FOCAS) at the Subaru 8.2-m telescope on Mauna Kea, Hawaii, on 24 September 2008 (Figure 1b). The peak of the emission, with a surface brightness of  $R = 23.5 \pm 0.2$  mag arcsec<sup>-2</sup> has again shifted away from SN 1572. A long-slit spectrum of the brightness peak of the echo structure was obtained with FOCAS on the same night, covering the wavelength range from 3,800 to 9,200 Å with a spectral resolution of 24 Å. The acquired echo spectrum unambiguously shows light of a supernova origin (Figure 2), a number of broad absorption and emission features from neutral and singly ionized intermediate mass elements. Type I supernovae are distinguished from those of type II by the absence of hydrogen, and type Ia supernovae are further distinguished from types Ib and Ic by a prominent Silicon 6,355 Å absorption feature at maximum light. We have compared the spectrum of SN 1572 with thermonuclear supernovae of different luminosity. Both sub- and overluminous type Ia supernovae, such as SN 1991bg and SN 1991T, respectively, showed peculiarities in their spectra near maximum light. We concluded that SN 1572 belongs to the majority class of normal type Ia supernovae [1].



**Figure 1:** (a) Optical R-band images of the Tycho's supernova light echo taken by Calar Alto 3.5 m telescope (black means bright). The rectangle shown in a indicates the location of a previous light-echo detection in 2006. The vector towards Tycho's supernova remnant is indicated (arrow). (b) R-band images taken by FOCAS on Subaru Telescope.



**Figure 2:** Spectrum of Tycho's supernova obtained by FOCAS. Black solid lines show the spectrum of Tycho's SN. Comparison with templates of subluminous (red), normal (orange) and overluminous (blue) type Ia supernova.

## Reference

[1] Krause, O., et al.: 2008, *Nature*, **320**, 1195.

# Supernova Remnant Cassiopeia A: Decoding Light to Uncover the Mystery of its Birth

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(NAOJ)

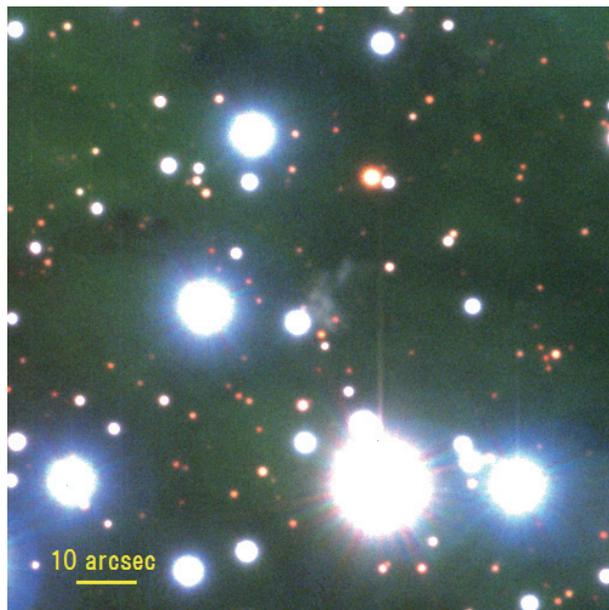
KRAUSE, Oliver, BIRKMANN, Stephan M., GOTO, Miwa  
(Max-Planck-Institut for Astronomy)

RIEKE, George H., MISSELT, Karl A.  
(Steward Observatory)

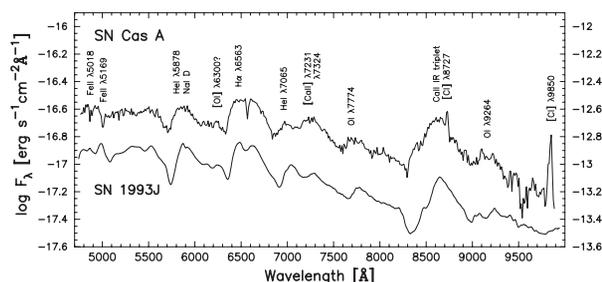
The supernova remnant Cassiopeia A (Cas A) is one of the most-studied objects in the sky, with observations from the longest radio waves to gamma rays. The lack of a definitive sighting means that there is almost no direct information about the type of the explosion, and the true nature of its progenitor star has been a puzzle since the discovery of the remnant. There is no chance to see the precise date when it happened, or how it looked like at the time of the outburst. The outburst supposedly happened in approximately 1680, judging from the current expansion rate. It is a mystery, however, why there is no widely recorded sighting of such an extravagant outburst in the 17th century, happened just 11,000 light years away from the earth. The light echo is a phenomenon that the wave of light is reflected or re-emitted by the material located away from its origin, and arrives to the observer with a delay. The light echo of Cas A we see today is therefore the light of the outburst that arrived on the earth 300 yrs after the direct light sweeps out the planet. No need to say, this light echo hides the secret of the birth of Cas A within. We started regular patrols of the region by Spitzer Space Telescope in the infrared wavelength, and wide field monitoring observation by MPIA Calar Alto 2.2 m telescope in Spain. The candidates were found in September 2007.

On the night of 9 October 2007, the light echo ( $R = 23.5$  mag) was confirmed with FOCAS (Fig. 1), and we started the spectroscopy. After 5.5 hours exposure, we got a spectrum showing the clear spectral lines, the finger prints of a very young SN (Fig. 2). Comparing with the data in the spectral database of various supernovae, we found that the spectra of the light echo is almost identical with that of SN 1993J (Fig. 2). The progenitor of Cas A must be the same with that of SN1993J: a red super giant more massive than 10 times of the solar mass. And they went through the same type of outburst: Type IIb. A supernova outburst of the type IIb is relatively faint ( $M_V = -17.5$ ) and fades away fast. Unfortunately weather of mere several days, that somehow overlapped with the maximum of the light curve, would be enough to let this young supernova escaped away from the history. At a relatively faint brightness, the light of the Cas A supernova event would have been extinguished relatively quickly by the dust grains in the foreground of the expanding supernova remnant along with the characteristically fast decay of light from a Type IIb

outburst. With these findings, the answer to our historical puzzle seems to be finally obtained.



**Figure 1:** Optical image of the echo region. It is a color composite of I-band (red), R-band (green), and V-band (blue) by FOCAS at Subaru Telescope. The faint white features in the middle of the image are the light echoes.



**Figure 2:** Spectrum of the Cas A supernova (upper) taken by FOCAS on 9 October 2007 and SN 1993J for a reference (bottom).

Reference

[1] Krause, O., et al.: 2008, *Science*, **320**, 1195.

# Detections of Strong Ionizing Radiation from High-redshift Galaxies

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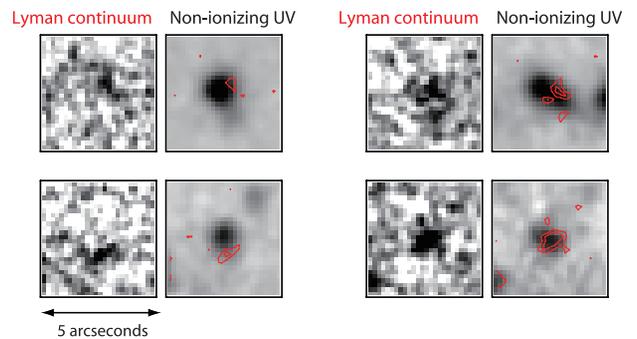
Ionizing radiation from star-forming galaxies is a likely primary source of cosmic reionization. Although the emissivity of Lyman continuum (LyC) from galaxies is a key to evaluate the contribution of galaxies to cosmic reionization, it has been poorly constrained due to the fact that LyC photons are easily absorbed by the intergalactic medium (IGM). Since direct observation of LyC at  $z > 4$  is virtually impossible because of a rapid increase in the number density of Lyman limit systems toward high redshifts, we must focus on  $z \sim 3$  where the IGM optical depth is still about unity on average. So far the number of galaxies with direct observation of LyC was too small to reveal the typical LyC emissivity. It is because the spectroscopy for LyC search requires much observing time for each target galaxy, and the number of galaxies observed was limited.

In order to confront this issue, we adopted a novel technique: imaging observation with a narrow-band filter to trace the LyC from galaxies at a specific redshift enables us to examine LyC from a large number of galaxies simultaneously. The target field of the present study is the SSA22 field where the prominent proto-cluster of galaxies at  $z = 3.09$  has been discovered, and there are 198 galaxies with spectroscopic redshifts in one Subaru / Suprime-Cam field. We fabricated a special narrow-band filter NB359 (central  $\lambda = 359$  nm with FWHM = 15 nm) to examine LyC from these  $z \approx 3.1$  galaxies, and executed observation in Sep. 2007.

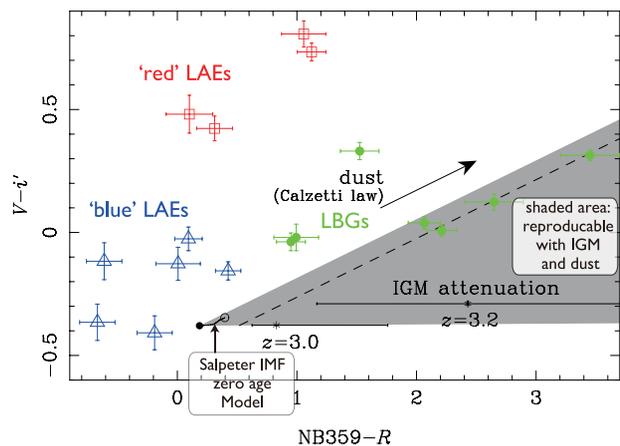
We detected 7 Lyman break galaxies (LBGs) and 10 Lyman  $\alpha$  emitters (LAEs) in NB359 among 198 galaxies with spectroscopic redshifts ([1]; Fig. 1). There are some objects with significant spatial offsets between LyC and non-ionizing UV radiation.

In Figure 2 colors of the detected galaxies are plotted. NB359-R color stands for the ratio of LyC and nonionizing UV radiation, while  $V-i'$  color represents UV slope. Most LAEs have NB359-R color bluer than LBGs. We also show predicted colors of model spectral energy distribution generated by the population synthesis code with Salpeter initial mass function, low metallicity and zero age. The colors which can be reproduced by considering dust and IGM attenuation are indicated with the shaded area. Because increasing age or metallicity make NB359-R color redder, the colors of detected galaxies (especially LAEs) cannot be explained with such

standard stellar populations. This result would indicate the existence of strong sources of ionizing radiations among high redshift star-forming galaxies, which might be a counterpart of galaxies in the cosmic reionization epoch. Further investigation including deep spectroscopy is required to clarify the nature of such strong LyC emitters.



**Figure 1:** Examples of galaxies with LyC detection. In UV ( $R$ -band) images contours of NB359 images are overlaid.



**Figure 2:** NB359-R and  $V-i'$  colors of galaxies with LyC detection.

Reference

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# Okayama Astrophysical Observatory has discovered ten exoplanets

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ITOH, Yoichi, MURATA, Daisuke  
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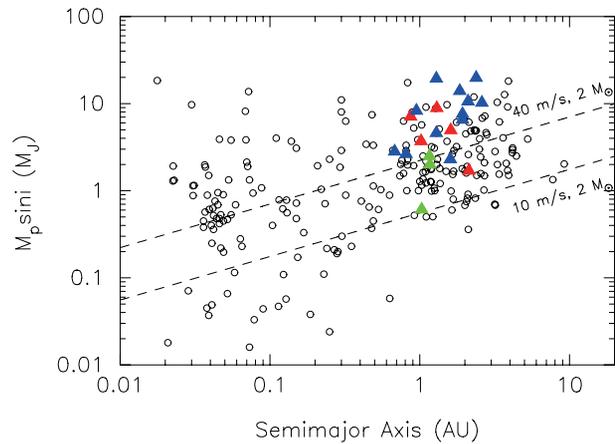
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Since the first discovery of the extrasolar planet in 1995, more than 300 extrasolar planets have been discovered so far. While most of them orbits around FGK dwarfs ( $0.7\text{--}1.5 M_{\odot}$ ), so-called solar-type stars, about 20 planets orbit around GK-type giants ( $1.5\text{--}4 M_{\odot}$ ). Such intermediate-mass stars are unsuitable for Doppler planet searches when they are on the main sequence (early-type stars) because they have few absorption lines in their spectra, which are often rotationally broadened, and thus it is difficult to attain high precision in radial velocity measurements. On the other hand, late-G to early-K type giants, which are evolved from early-type dwarfs, have many sharp absorption lines in their spectra appropriate for precise radial velocity measurements due to low surface temperature and small rotational velocity. Thus, giant stars are now regarded as ideal targets for planet searches around intermediate-mass stars.

Since 2001 we have been carrying out a Doppler planet search program targeting about 300 GK giants using the 1.88 m telescope and High Dispersion Echelle Spectrograph at Okayama Astrophysical Observatory. We discovered seven new planets in 2008 from the program [1, 2]. The planets have minimum masses of  $1.6\text{--}10 M_{\text{JUP}}$ , semimajor-axes of  $0.7\text{--}2.6 \text{ AU}$ , and eccentricities of  $0\text{--}0.2$ . Their host stars have masses of  $1.6\text{--}2.7 M_{\odot}$ . The number of our discoveries reached at 10, which is about a half of the currently known planets around giant stars.

One of the interesting properties of planets around giants is paucity of inner planets. All of the known planetary companions to giant stars orbit beyond  $0.7 \text{ AU}$  and no planets have been found within the limit. Two scenarios can account for the orbital distribution; one is that inner planets are primordially deficient around intermediate-mass stars and the other is that they have been engulfed by the expanding central stars due to stellar evolution. When we assume that many of clump giants are post-RGB stars (core-helium-burning stars), planets in orbits with  $\leq 0.5 \text{ AU}$  could have been engulfed by the central stars at the tip of RGB ( $R_{\text{star}} \sim 25\text{--}40 R_{\odot}$  for a  $2\text{--}3 M_{\odot}$  star) due to tidal force from the central stars. It means that we can not reject the possibility that planets had originally existed in short orbital distances around progenitors of clump giants. Other interesting properties

of planets to giants, such as higher frequency of super-massive planets with  $\geq 5 M_{\text{JUP}}$  compared to the case of solar-type stars and lack of metal-rich tendency in planet-host giants, have begun to emerge [3]. We try to increase the number of targets and detected planets further in order to confirm their statistical properties.



**Figure 1:** Planet mass plotted against semimajor axis. Planets around low-mass ( $< 1.6 M_{\odot}$ ) giants, intermediate-mass subgiants ( $1.6\text{--}1.9 M_{\odot}$ ), and clump giants ( $1.7\text{--}3.9 M_{\odot}$ ), are plotted by red, green, and blue triangles, respectively. Planets around solar-type dwarfs are plotted by open circles. Dashed lines express detection limits with velocity semi-amplitude of  $10 \text{ m s}^{-1}$  and  $40 \text{ m s}^{-1}$  around a  $2 M_{\odot}$  star. Since giants have intrinsic variability in radial velocity of  $10\text{--}20 \text{ m s}^{-1}$  due to stellar activity, it is difficult to detect planets which impart velocity variations less than  $40 \text{ m s}^{-1}$  to central stars.

## References

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- [3] Takeda, Y., Sato, B., Murata, D.: 2008, *PASJ*, **60**, 781.

# Fe XIII Density Diagnostics in the EIS Observing Wavelengths

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The EUV Imaging Spectrometer (EIS) on board the *Hinode* satellite observes solar emission lines at EUV wavelengths with high spatial and spectral resolution [1]. The two wavelength bands observed with EIS are 170–210 Å and 250–290 Å, and include emission lines of HeII and a wide range of iron ions from FeVIII through to FeXXIV, formed in the temperature range  $4.7 \leq \log T_e \leq 7.2$ .

Among the emission lines of the various iron species, the Fe XIII lines found in the EIS short wavelength band (170–210 Å) are particularly important, as they provide several line pairs sensitive to the density of the coronal plasma. As an independent parameter to the emission measure, the density is crucial to determining the plasma filling factor, i.e., the fraction of the observed plasma volume that is emitting radiation.

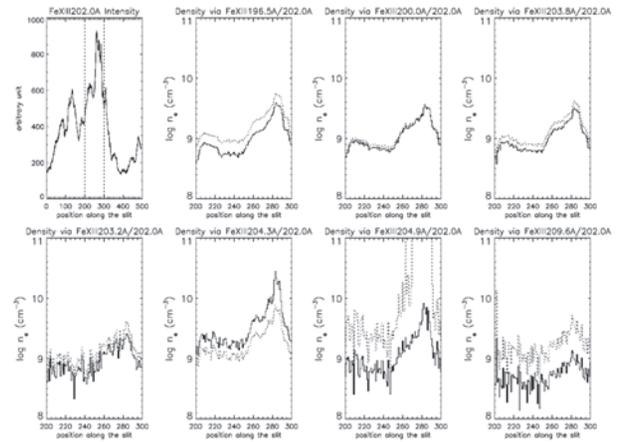
The diagnostic capability of Fe XIII line features seen in the EIS observing wavelengths has been studied. Full CCD slit spectra of the EIS first-light active region AR 10921 observed in the short wavelength band (170 Å–210 Å) are used to provide a dataset of Fe XIII emission lines, and their intensities are compared with two theoretical models for density diagnostics: The first is a modification of the CHIANTI v5.2 model [2] that is modified to use the Gupta-Tayal (1998) [3] electron collision file, and the second is the model of Yamamoto et al. (2008) [4] (hereafter Y08) that uses the electron collision data of Aggarwal and Keenan (2005) [5].

Derived electron densities are plotted against the position along the slit in Figure 1. In the density range from  $10^{8.5}$  to  $10^{9.5} \text{ cm}^{-3}$  the four density diagnostic line pairs;  $\lambda 96.5 \text{ Å}$ ,  $\lambda 200.0 \text{ Å}$ ,  $\lambda 203.2 \text{ Å}$ , and  $\lambda 203.8 \text{ Å}$  vs  $\lambda 202.0 \text{ Å}$  give consistent densities. These are all strong lines appearing in the EIS waveband except  $\lambda 203.2 \text{ Å}$ . The Y08 model gives slightly better consistency between the ratios than the modified CHIANTI model for these strong line pairs.

For other Fe XIII lines we find problems with both of the theoretical models: from the Y08 model,  $\lambda 204.3 \text{ Å}$  and  $\lambda 209.6 \text{ Å}$  are not well predicted; while for the modified CHIANTI model  $\lambda 204.3 \text{ Å}$  and  $\lambda 204.9 \text{ Å}$  are inconsistent with observations. They are generally weak lines seen in the EIS waveband, and the problems with the  $\lambda 209.6 \text{ Å}$  and  $\lambda 209.9 \text{ Å}$  may be partly due to the instrument calibration, as they lie very close to end of the EIS waveband.

The accuracy of the high density limit of the important  $\lambda 203.8 \text{ Å}/\lambda 202.0 \text{ Å}$  density sensitive ratio has also been investigated, using an EIS observation of a C-class flare from on 2007 January 16. The measured ratio is consistent with the theoretical predictions both from the Y08 and modified CHIANTI models.

Watanabe et al. (2009) [6] describe the detail of the analysis.



**Figure 1:** Densities derived from observed line ratios at the positions (pixel numbers: 200–300) along the slit. Total intensity of Fe XIII  $\lambda 202.0 \text{ Å}$  is plotted in the top-left panel. Solid lines are for the model of Y08 and dotted lines are for the CHIANTI.

## References

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- [3] Gupta, G. P., Tayal, S. S.: 1998, *ApJ*, **506**, 464.
- [4] Yamamoto, N., et al.: 2008, *ApJ*, **689**, 646.
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- [6] Watanabe, T., et al.: 2009, *ApJ*, **692**, 1294.

# Laser interferometric high-precision geometry monitor for JASMINE

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A heterodyne metrology interferometer was stabilized down to a noise level of 20 picometers (pm) as a rootmean-square (RMS) value integrated between 0.3 mHz and 1 Hz. This interferometer can be used in precise metrology applications, such as characterization of deformation for satellite optical components against thermal exposure.

JASMINE (Japan Astrometry Satellite Mission for Infrared Exploration) [1] is one of next-generation astrometry satellite missions and aims to measure parallaxes, positions and proper motion with the precision of 10 microarcsec for ten million stars in the bulge of the Milky Way. Now, we are developing small-JASMINE, which is a downsized version of JASMINE and planned to be launched around 2015. In order to realize high-precision astrometric observation, geometry of a small-JASMINE telescope with the structure size of approximately 1 m should be stabilized with an accuracy of 100 pm in RMS value for a frequency range from 1 mHz to 1 Hz. In order to ensure such requirement by ground experiments, thermal deformation properties of telescope materials should be measured with the same accuracy under the similar environment to space. A heterodyne metrology interferometer can be used as a high-precision displacement sensor for such a laboratory test. Then, we had begun to develop a heterodyne metrology interferometer [2] and stabilized it so as to verify that a noise level satisfies the requirement for small-JASMINE.

A conceptual design of our heterodyne metrology interferometer is shown in Fig. 1. The heterodyne interferometer was built on a 50 mm square ultra-low expansion glass plate (OHARA, CLEARCERAM® -Z) in order to reduce an optical path length change caused by temperature variation. Optical components are bonded to the glass plate by hydroxide-catalysis bonding so as to make the whole structure virtually monolithic. An optical configuration of the interferometer is a Mach-Zehnder interferometer with a design as symmetric as possible so that a detection signal can be insensitive to homogeneous thermal expansion of the glass plate. The heterodyne frequency is actively controlled in order to suppress residual noises that are caused by optical path length changes outside of the glass plate as well as phase fluctuations of the heterodyne frequency source. This configuration allowed us to concentrate our effort for stabilization only at a 50 mm square optical base plate.

Figure 2 shows the power spectrum density and the

RMS of the displacement measured by the stabilized interferometer. The displacement RMS noise level is 20 pm for the frequency range from 1 Hz to 1 mHz. This meets the requirement for small-JASMINE.

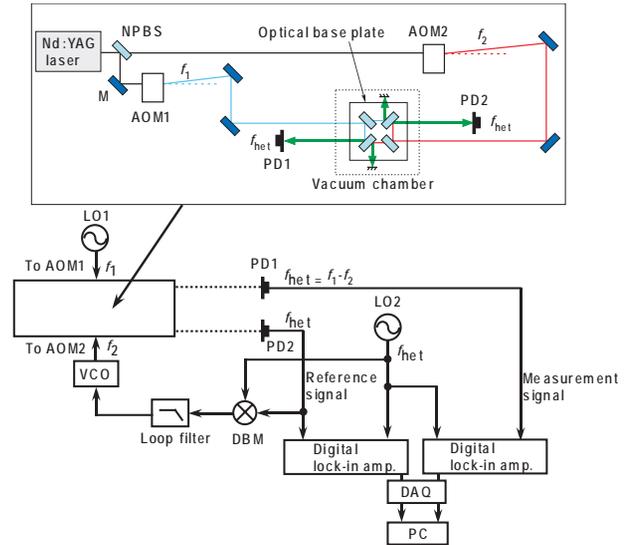


Figure 1: Schematic drawing of a heterodyne interferometer.

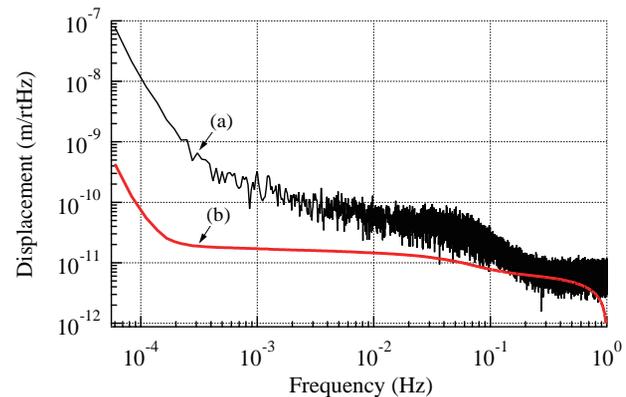


Figure 2: The displacement noise level of the stabilized interferometer: (a) the power spectrum density, (b) the RMS.

## References

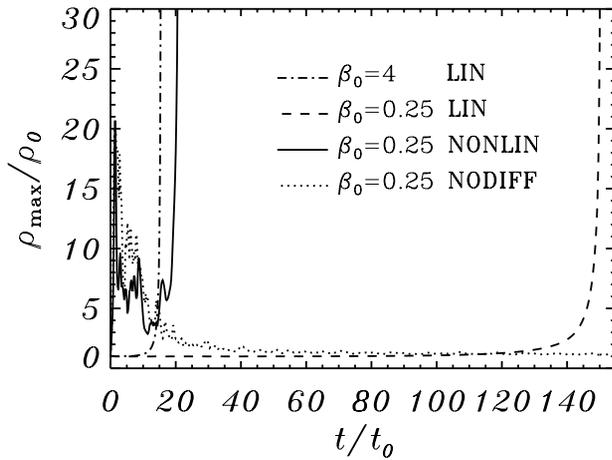
- [1] Gouda, N., et al.: 2008, *Proc. of IAU Symp.*, **248**, 248–251.
- [2] Niwa, Y., et al.: 2008, *SPIE*, **7018**, 70183D-1–70183D-12.

# Three-dimensional Simulation of Magnetized Cloud Fragmentation Induced by Nonlinear Flows and Ambipolar Diffusion

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(NAOJ)

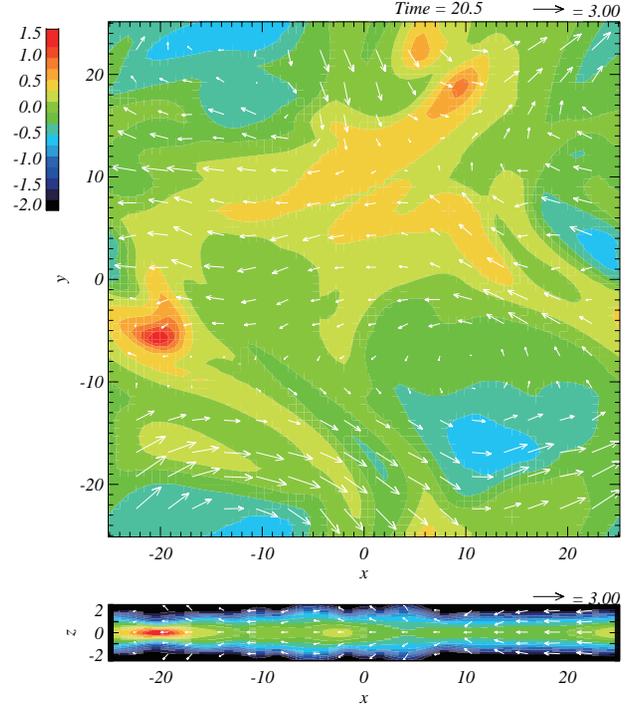
BASU, Shantanu  
(University of Western Ontario)

We demonstrate that the formation of collapsing cores in subcritical clouds is accelerated by nonlinear flows, by performing three-dimensional non-ideal MHD simulations [1]. (In the subcritical clouds, magnetic force suppress the gravitational collapse unless the magnetic diffusion works.) This is the first 3D simulation that would alter the previous estimation of the starformation time scale.



**Figure 1:** The time evolution of maximum densities at  $z = 0$ . The solid line shows the evolution for an initially nonlinear supersonic perturbation and  $\beta_0 = 0.25$  (see Fig. 2). The dashed line ( $\beta_0 = 0.25$ ) and the dash-dotted line ( $\beta_0 = 4$ ) show the evolution for models with a linear initial perturbation as calculated by Kudoh et al. (2007) [2]. The dotted line shows the evolution for an initially nonlinear supersonic perturbation and  $\beta_0 = 0.25$ , but without ambipolar diffusion. The time unit  $t_0$  is about  $2.5 \times 10^5$  years in this model.

An initial random supersonic (and trans-Alfvénic) turbulent-like flow is input into a self-gravitating gas layer that is threaded by a uniform magnetic field (perpendicular to the layer) such that the initial mass-to-flux ratio is subcritical. (In this model, the subcritical means  $\beta_0 < 1$ , where  $\beta_0$  is the initial ratio of gas to magnetic pressure at  $z = 0$ .) Magnetic ambipolar diffusion occurs very rapidly initially due to the sharp gradients introduced by the turbulent flow. It subsequently occurs more slowly in the traditional near-quasistatic manner, but in regions of greater mean density than present in the initial state. The overall timescale for runaway growth of the first core(s) is several  $\times 10^6$  yr, even though previous studies have found a timescale of several  $\times 10^7$  yr when starting with linear perturbations and similar physical parameters (Fig. 1, 2).



**Figure 2:** Logarithmic density image at  $t = 20.5 t_0$  for the nonlinear perturbation case of the subcritical cloud ( $\beta_0 = 0.25$ ). The top panel shows the cross section at  $z = 0$ , and the bottom panel shows the  $x$ - $z$  cross section at  $y = -5.9 H_0$ , where  $H_0$  is the scale height of the gas layer ( $H_0 \sim 0.05$  pc). A collapsing core is formed in the vicinity of  $x = -20 H_0$  and  $y = -6 H_0$ .

Large-scale supersonic flows exist in the cloud and provide an observationally testable distinguishing characteristic from core formation due to linear initial perturbations. However, the nonlinear flows have decayed sufficiently that the relative infall motions onto the first core are subsonic, as in the case of starting from linear initial perturbations. The ion infall motions are very similar to those of neutrals; however, they lag the neutral infall in directions perpendicular to the mean magnetic field direction and lead the neutral infall in the direction parallel to the mean magnetic field.

## References

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- [2] Kudoh, T., et al.: 2007, *MNRAS*, **380**, 499.

# How was the mushroom-shaped GW 123.4–1.5 formed in the Galactic disk?

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(NAOJ)

GW 123.4–1.5 is a mushroom-shaped HI structure composed of a stem and a cap and is dissimilar to any common shells or chimneys in our Galactic disk. To understand the origin and velocity structure of GW 123.4–1.5, we extend the previous two-dimensional hydrodynamical simulation [1] of the collision of a high-velocity cloud (HVC) with the Galactic disk to three dimensions [2].

Figure 1 shows the evolution of column-density on the  $x$ - $z$  plane (integrating along the  $y$ -axis) for model O3 (the angle from the vertical axis of the Galactic disk is  $3^\circ$ , the number density is  $0.1 \text{ cm}^{-3}$ , and the incident speed is  $100 \text{ km s}^{-1}$ ). We can find the mushroom-shaped structures between  $t = 56 \text{ Myr}$  (or  $t/t_0 = 4.0$ ) and  $t = 67 \text{ Myr}$  (or  $t/t_0 = 4.8$ ). The height of the mushroom at the age of  $t = 67 \text{ Myr}$  is about  $350 \text{ pc}$ , which is the same size as that of GW 123.4–1.5. The lifetime of the mushroom structure is  $\sim 10^7 \text{ yr}$  and the width ratio of the cap to the stem is about 3:1. The mass of the cap is about twice that of the stem, and the total mass is estimated to be  $\sim 10^5 M_\odot$ . The mean column-density contrast between the mushroom and ambient medium is about 2:1. Although the mass ratio is a little smaller than the observed one 4:1, most physical parameters of the mushroom-shaped structure in model O3 coincide with those of GW 123.4–1.5.

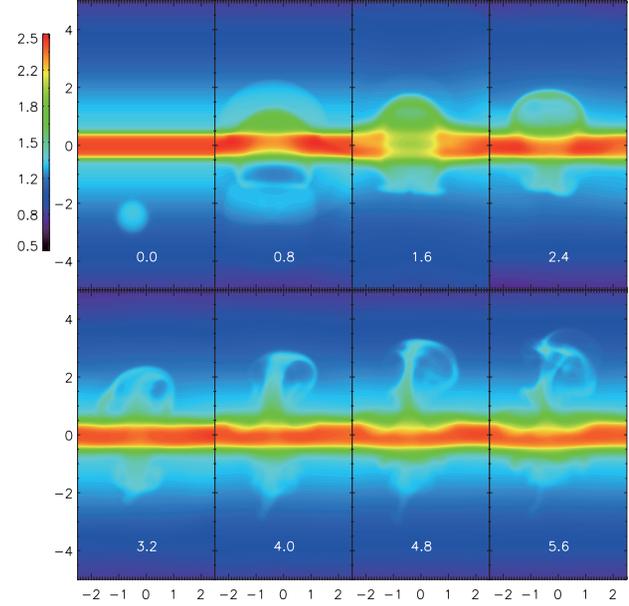
Figure 2 shows the column-density and the position-velocity maps, which are equivalent to the observational longitude-velocity ( $l$ - $v$ ) maps, for model O3. The velocity gradient in Figure 2b is very similar to the observed value of GW 123.4–1.5.

GW 123.4–1.5 is expected to be formed by the almost head-on collision of a HVC with velocity  $\sim 100 \text{ km s}^{-1}$  and mass  $\sim 10^5 M_\odot$  about  $5 \times 10^7 \text{ yr}$  ago. A mushroom-shaped structure like GW 123.4–1.5 must be infrequent on the Galactic plane, because the head-on collision which explains the mushroom structure seems rare for observed HVCs. HVC-disk collision explains not only the origin of the mushroom-shaped structure but also the formation of a variety of structures like shells, loops, and vertical structures in our Galaxy.

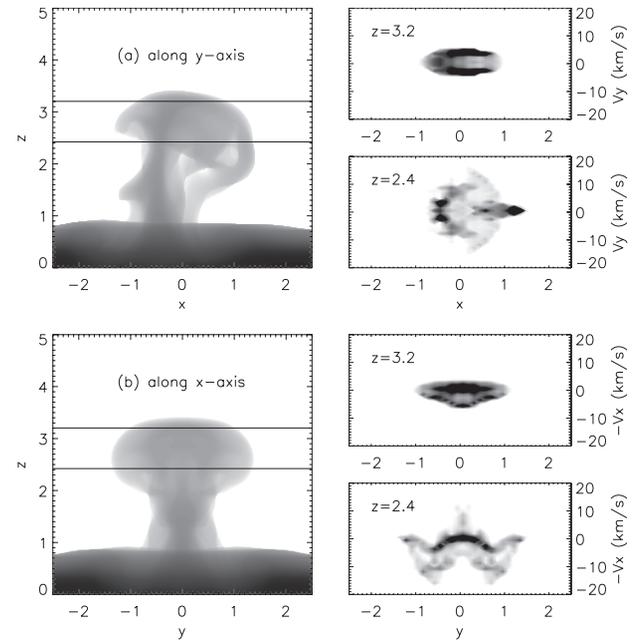
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\* Chang Hyun Baek was invited to NAOJ as a Post-Doctoral Fellow by Korea Research Foundation from October 2006 to September 2007.



**Figure 1:** Time evolution of column-density of cool gas on the  $x$ - $z$  plane (integrating along the  $y$ -axis). The snapshots are made at  $t/t_0 = 0, 0.8, 1.6, 2.4$  (top: left to right),  $3.2, 4.0, 4.8, 5.6$  (bottom: left to right) for model O3.



**Figure 2:** Column-density (left panels) and position-velocity (right panels) at  $t/t_0 = 4.8$  in model O3. Position-velocity is measured for the lower ( $z = 2.4 H_0$ ) and upper ( $z = 3.2 H_0$ ) regions of the caps on (a) the  $x$ - $z$  plane (along the  $y$ -axis) and on (b) the  $y$ - $z$  plane (along the  $x$ -axis).

# Development of a hybrid N-body simulation code for planet formation process that can handle number of large particles

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In the standard model of planet formation scenario proposed by Hayashi et al., planetesimals evolve to protoplanets through mutual collisions. N-body simulations have been used for the study of this process, and the formation mechanisms such as the runaway growth and the oligarchic growth have been found. In previous studies, the number of particles have been limited to around  $10^4$ , and therefore minimum particle mass used was around  $10^{23}$  g. However, we need to examine the behavior of small planetesimal, since it is believed that planetesimals of the mass  $10^{16}$ – $10^{19}$  g are first formed from dust. Moreover, because there is a limitation in the number of particles that can be handled, in almost all N-body simulations the assumption of perfect accretion has been used. This assumption is not appropriate for the process of the planet growth in the late stage. Thus a simulation code that can handle a large number of particles is necessary to study these problems. It can also handle fragmentation that increases the number of planetesimals.

In N-body simulations, the calculation cost increases as the second power of the number of particles. Moreover, in planet formation simulations, since the formation time ( $10^6$ – $10^7$  years) is far longer than the orbital timescale of planetesimals, very long calculations are necessary. We developed a high speed simulation code that can handle a large number of particles. Since gravity decrease as the inverse square of the distance, distant particles are organized to calculate gravity with keeping calculation cost lower. In addition, it can integrate with long time steps, because gravity of distant particles have smaller time variation than that of near particles. Hence, we can calculate with high accuracy while keeping low calculation cost and time.

Specifically, we split the mutual gravity of particles as the distance [1]. The gravity can be divided into gravity of near particles with rapid variation and that of distant particles with slow variation. Moreover, the distant particle gravity is calculated by the tree method [2], because it is relatively small compared to the solar gravity and forces from nearby particles gravity. They are calculated by the fourth Hermite scheme with small and individually variable time steps. It can integrate at high speed while keeping high accuracy by combining these methods [3]. Figure 1 shows the relative energy error as a function of time. The energy error changes randomly, so this code can integrate the system with small error for long calculations. The calculation cost of this code is about factor 60 smaller than conventional

direct calculating code for  $10^4$  particles. Therefore, for  $10^6$  particles our code is 6000 times faster than the direct calculation. Then we can calculate over  $10^6$  particles.

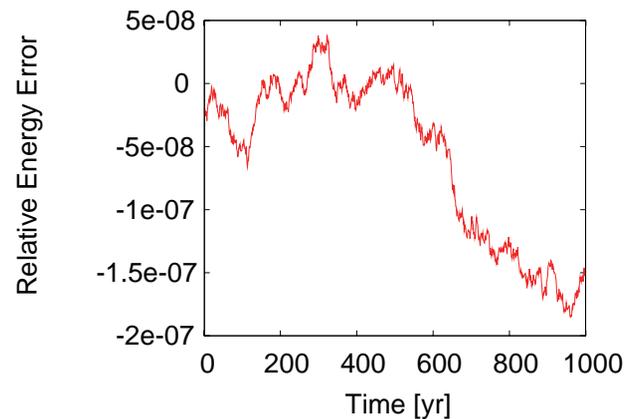


Figure 1: Relative energy error as a function of time.

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# Wide-Field Infrared Imaging Polarimetry of the NGC 6334 Region: A Nest of Infrared Reflection Nebulae

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(NAOJ)

NAGATA, Tetsuya, NAGAYAMA, Takahiro  
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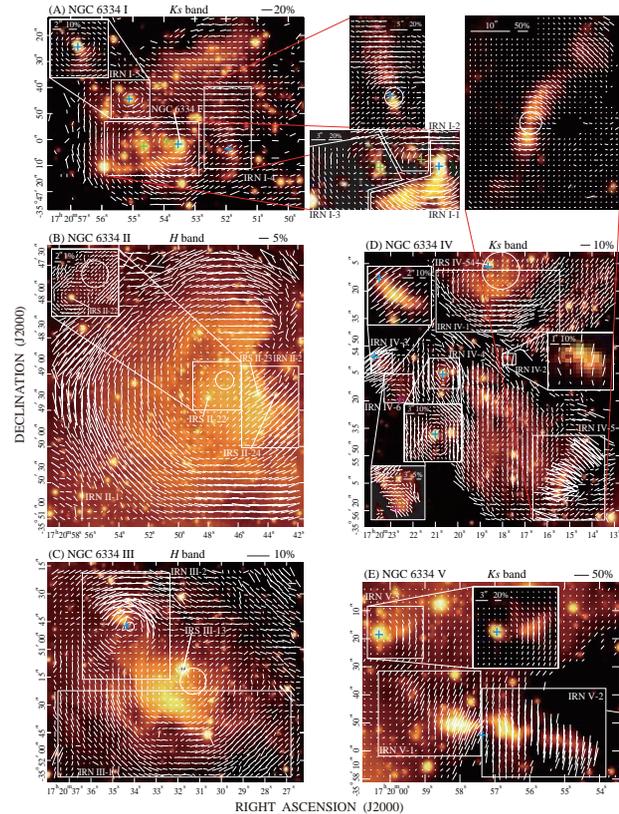
KURITA, Mikio  
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HOUGH, Jim, CHRYSOSTOMOU, Antonio  
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We report the detection of eighteen infrared reflection nebulae (IRNe) in the *J*, *H*, & *Ks* linear polarimetric observations of the NGC 6334 massive star-formation complex [1], a total bolometric luminosity  $\sim 1.9 \times 10^6 L_{\odot}$  [2], of which 16 IRNe are new discoveries. Our images cover  $\sim 180$  square arcminutes, one of the widest near-infrared polarization data in star-formation regions so far. These IRNe are most likely associated with embedded young OB stars at different evolutionary phases, showing a variety of sizes, morphologies, and polarization properties, which can be divided into four categories of Type A to D. We argue the above classification can be placed along the approximate evolutionary sequence, i.e., Type A is the youngest and evolve into D through B and C.

We suggest that the difference of the polarized patterns reflects the evolutionary sequence of IRNe; the youngest IRNe with well aligned polarization (*Type A*) evolve into monopolar/bipolar IRNe with roughly centrosymmetric polarization (*Type B*). When IRNe have concentric polarizations (*Type C*), the illuminating sources create HII regions, and then, IRNe become less distinct near the central sources with a centrosymmetric circular pattern at the region of the ionization boundary (*Type D*). We note, however, that care must be taken to infer above conclusions from the polarization pattern alone. Further high resolution radio observations for kinematical information should be undertaken to verify our IRN evolutionary hypothesis.

This research was partly supported by Grants-in-Aid for Scientific Research on Priority Areas from the Ministry of Education, Culture, Sports, Science and Technology (MEXT) and No.16340061, 16077204, and 16340061.



**Figure 1:** Polarization vectors are superposed on the intensity images in the *H* or *Ks* band. In the magnifications, except for IRN IV-6, the vectors are superposed on the polarized intensity images in the same band. The white square indicates each IRN, and the circle shows each position of the illuminating sources with  $1\sigma$  error. Blue crosses represent the positions of *Spitzer*  $8.0\mu\text{m}$  sources, green crosses represent the MIR sources for IRN I-2 and 3, and purple cross represents the radio source for IRN IV-6. For IRN I-1, I-5, IV-4, and V-3, since illuminating sources are identified with NIR sources, the error circles are not marked.

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# New Constraint on Radiative Decay of Long-lived Particles with New ${}^4\text{He}$ Photodisintegration Data and Primordial ${}^6\text{Li}$ Problem

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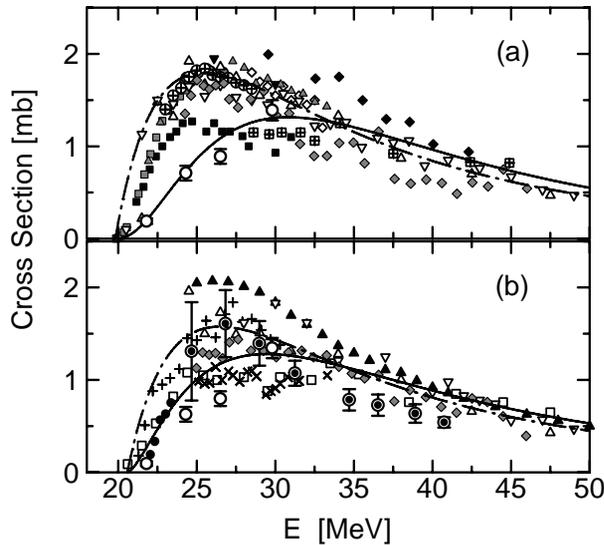
YOSHIDA, Takashi  
(NAOJ)

SHIMA, Tatsushi, NAGAI, Yasuki  
(Osaka University)

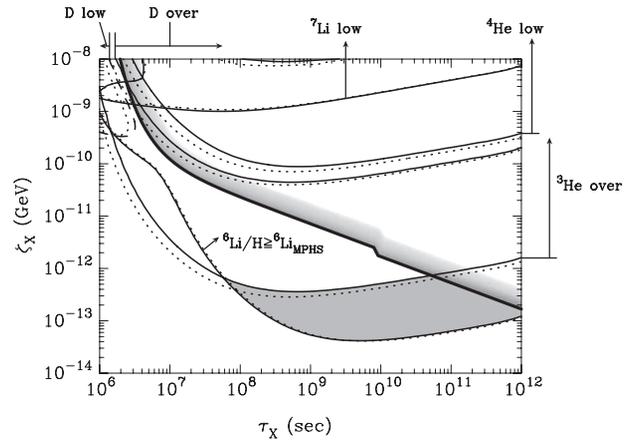
KII, Toshiteru  
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Spectroscopic observations of metal-poor halo stars (MPHSs) indicate that abundances of  ${}^6\text{Li}$  and  ${}^7\text{Li}$  on the stellar surface are independent of the metallicity. The observed value of  ${}^7\text{Li}$  is a factor of 2–4 smaller and that of  ${}^6\text{Li}$  a factor of  $\sim 10^3$  higher than predictions by standard big bang nucleosynthesis (BBN) model. The radiative decay of long-lived particles in the early universe could trigger through nonthermal processes production of  ${}^6\text{Li}$  up to about ten times more than observed in MPHSs [1]. The photodisintegration of  ${}^4\text{He}$ , which is involved in the nonthermal nucleosynthesis triggered by the radiative decay, was newly measured with precise quasimonochromatic laser-Compton photon beam [2]. They show smaller cross sections than deduced from previous experiments. We then calculate the decay-induced nonthermal nucleosynthesis with new cross sections fitted to the new experiment (Fig. 1) and derived a new constraint on the radiative decay (Fig. 2) [3].

long-lived particles  $X$ , a lifetime  $\tau_X$  and  $\zeta_X = (n_X^0/n_\gamma^0)E_{\gamma 0}$ . ( $n_X^0/n_\gamma^0$ ) which is the number ratio of  $X$  to photon before  $X$ -decay and  $E_{\gamma 0}$  is the photon energy emitted in the radiative decay. The change in cross sections of  ${}^4\text{He}$  photodisintegration affects the yields of D,  ${}^3\text{He}$  and  ${}^4\text{He}$ , which are related to the  ${}^4\text{He}$  photodisintegration at low energy ( $\sim 30$  MeV). The upper limit of allowed region of  $\zeta_X$  for these nuclides shifts upward by  $\sim 300\text{--}30\%$  for  $\tau_X = 10^6\text{--}10^{10}$  s accordingly. The nonthermal yield of  ${}^6\text{Li}$  is not sensitive to the change in cross sections at low energy, since the  ${}^6\text{Li}$  production is mainly through secondary processes which need energetic photons of  $E_\gamma \geq 50$  MeV. We study a possible uncertainty in the cross sections at high energies associated with sum rules, and found that it causes a factor of  $\leq 3$  change in the constraint on the parameter  $\zeta_X$ .



**Figure 1:** The cross sections of the  ${}^4\text{He}(\gamma,p){}^3\text{H}$  (a) and  ${}^4\text{He}(\gamma,n){}^3\text{He}$  (b) reactions as a function of the photon excitation energy  $E_\gamma$ . The open circles correspond to the new data from experiments with quasi-monochromatic laser-Compton photon beams [2]. The solid curves are the most probable excitation functions [3] while the dash-dotted curves the previous fitting functions [4]. This figure is reprinted from [3].



**Figure 2:** Contours corresponding to constraints on the primordial light element abundances for calculations with new cross sections of  ${}^4\text{He}$  photodisintegration [3] (solid lines) and with the previous ones [4] (dotted lines). Overabundance (over) and underabundance (low) regions of D,  ${}^3\text{He}$  and  ${}^7\text{Li}$  are excluded from a comparison with observations. The region above the thick solid line is excluded by the consistency requirement of the background radiation with a black body. The grey region is the allowed region where  ${}^6\text{Li}$  is produced more than observed in MPHSs. This figure is reprinted from [3].

We calculate the non-thermal nucleosynthesis, and constrain the allowed parameter region for

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# Production mechanism of circular polarization in the ORION BN/KL region: Near-infrared circular polarimetry by SIRPOL

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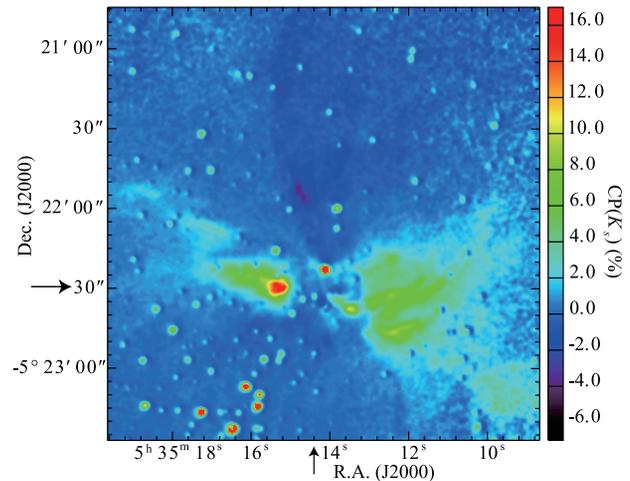
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We reported a deep  $K_s$  band circular polarization (CP) image of the Orion BN/KL region, a massive star-forming region [1, 2]. Further observations of CP from star-forming regions were necessary to reveal the production mechanism of CP.

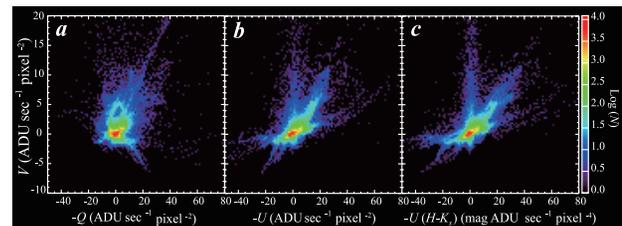
This study is the first result of circular polarimetry by a polarimeter SIRPOL mounted on the IRSF 1.4 m telescope in South Africa. SIRPOL in CP mode saw first light in the end of 2006, after SIRPOL in linear polarization (LP) mode saw first light in the end of 2005.

Figure 1 shows the distribution of CP around the most massive star in this region, IRC2. This distribution has a quadrupolar structure of right- and left-handed CP. Figure 2 shows, for the first time, correlations of CP, linear polarization (LP), and  $H-K_s$  color representing extinction. These quality correlations were yielded from CP and LP [3] data which were obtained by the same instrument and telescope.

We theoretically derived a simple relation between dichroic extinction, color excess, CP, and LP. This derivation for a young stellar object was yielded by analysis of the basic equation of vector radiative transfer. This derived relation agreed with the observed correlation between the Stokes parameters and the color excess. This result suggests a major contribution of dichroic extinction to the production of CP in this region.



**Figure 1:** Image of the degree of CP (%) in the  $K_s$  band ( $2.14 \mu\text{m}$ ) of the BN/KL region. The field of view is  $2'.7$  square. A positive sign for CP indicates that the electric vector is rotated counterclockwise in a fixed plane relative to the observer. The embedded massive star IRC2, which has about 25 solar masses, is indicated by arrows.



**Figure 2:** Pixel-by-pixel correlation using logarithmic color: (a) between Stokes  $V$  and  $Q$ ; (b) between Stokes  $V$  and  $U$ ; (c) between Stokes  $V$  and the product of Stokes  $U$  and  $H-K_s$  color. Stokes  $Q$  and  $U$  express linear polarization, and  $V$  expresses circular polarization.

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# New nuclear cosmochronometer for supernova neutrino process

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Long-lived radioactivities are celestial nuclear cosmochronometers to date when stellar nucleosynthesis occurred, which have been widely used for estimation of ages of early generations of stars and the Milky way. Since the number of the long-lived radioactivities is limited, there exist only six cosmochronometers,  $^{40}\text{K}$ ,  $^{87}\text{Rb}$ ,  $^{176}\text{Lu}$ ,  $^{187}\text{Re}$ ,  $^{232}\text{Th}$ , and  $^{238}\text{U}$ , whose half-lives are of the cosmological time scale 1–100 Gyr. However, there is no chronometer for neutrino reaction nucleosynthesis in supernovae ( $\nu$ -process).

We have proposed a  $^{138}\text{La}$  ( $T_{1/2}=102$  Gyr) -  $^{138}\text{Ce}$  -  $^{136}\text{Ce}$  system as a chronometer to measure the time elapsed from a single  $\nu$ -process [1]. These three nuclei are classified to the p-nuclei and  $^{138}\text{La}$  is considered to be synthesized dominantly by the  $\nu$ -process [2]. This chronometer is applied to samples affected by a single nucleosynthesis episode, for example, presolar grains in primitive meteorites. Recent studies of the presolar grains suggest that such samples are available. If the abundances of  $^{136,138}\text{Ce}$  and  $^{138}\text{La}$  in a sample affected strongly by a single supernova are known, an age  $T$  elapsed from the supernova episode is calculated by

$$T = -\frac{T_{1/2}(^{138}\text{La})}{\ln 2} \times \ln\left(\frac{\frac{N(^{138}\text{La})}{N(^{138}\text{Ce})}}{\frac{N(^{138}\text{La})}{N(^{138}\text{Ce})} + \frac{1}{b}(1 - R_{pp}(\text{Ce})) \times \frac{N(^{136}\text{Ce})}{N(^{138}\text{Ce})}}\right), \quad (1)$$

where  $N$  is the observed abundance in the sample,  $R_{pp}$  is the initial  $N(^{138}\text{Ce})/N(^{136}\text{Ce})$  ratio at the freezeout and  $b$  is the branching ratio of the  $\beta^-$  decay of  $^{138}\text{La}$ . The  $R_{pp}(\text{Ce})$  ratio can be calculated by an empirical scaling law between two p-nuclei with the same atomic number, which was found in the solar system abundances. The detailed mechanism of the empirical scaling law was presented in our previous paper [3]. The  $R_{pp}(\text{Ce})$  ratio is estimated to be 1.2.

In addition, we have pointed out that the nuclear structure of  $^{138}\text{La}$  is crucial for the performance of the chronometer and the  $\nu$  process origin of  $^{138}\text{La}$ . The lowest  $1^+$  state has not been established experimentally. If this  $1^+$  excited state exists at energy lower than 72 keV that is the energy of the first excited state, the  $1^+$  state may be a  $\beta$  unstable isomer. If so,  $^{138}\text{La}$  synthesized by the  $\nu$ -process is destroyed via the isomer (see Fig. 1).

Neumann-Cosel et al., commented to us that a recent nuclear experiment on  $^{138}\text{La}$  [4] indicates that the

existence of a low-energy  $1^+$  state in  $^{138}\text{La}$  is extremely unlikely [5]. This is consistent with our shell model calculation in a relatively large model space [1]. We conclude that there is no evidence that suggests the existence of the  $1^+$  isomer in  $^{138}\text{La}$  and our proposed cosmochronometer is robust from this viewpoint [6].

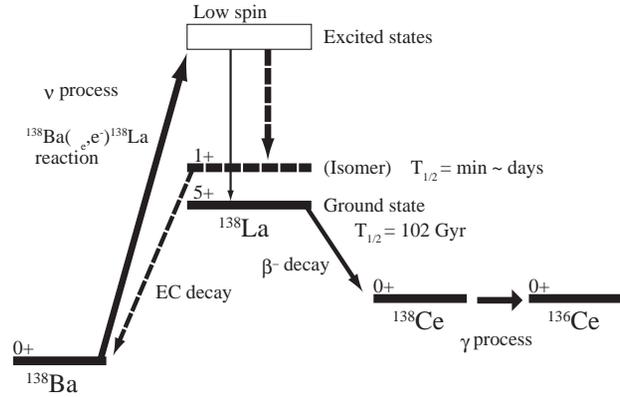


Figure 1: Schematic view of nucleosynthesis of  $^{138}\text{La}$ .

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# Necessary conditions for super-integrability of Hamiltonian systems

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A super-integrable system is defined to be a system having maximum number of first integrals among integrable Hamiltonian systems. The Kepler problem is a super-integrable system. In a super-integrable system, all bounded orbits are closed (periodic) like the elliptic orbits in the Kepler problem (Fig. 1).

A particle motion in a centrifugal potential  $V = V(r)$  is always integrable. It is known that among them, potentials in which bounded orbits are always periodic are limited to only two cases; the Kepler problem with  $V(r) = -1/r$ , and the isotropic harmonic oscillator with  $V(r) = r^2$ . This is known as Bertrand's theorem (1873). Since periodicity of all bounded orbits is a consequence of super-integrability of the system, the classical Bertrand's theorem is considered to have listed up all super-integrable centrifugal potentials. In the present work, a necessary conditions for super-integrability of Hamiltonian systems with homogeneous potential is given, which is based on a theorem of Morales-Ramis (2001) on necessary conditions for integrability.

For a potential system  $\ddot{\mathbf{q}} = -\nabla V(\mathbf{q})$  with a 2-dim. homogeneous potential  $V(\mathbf{q})$  of degree  $k$ , fix a solution  $\mathbf{q} = \mathbf{c}$  of the simultaneous algebraic equation  $\nabla V(\mathbf{q}) = \mathbf{q}$ . Then compute the quantity

$$\lambda = \nabla^2 V(\mathbf{c}) - (k - 1)$$

The Morales-Ramis theorem says that if the system under consideration is integrable, then the  $\lambda$  must be one of the discrete values (two spectrums)

$$\lambda \in \{0, 1, k - 1, k + 2, 3k - 2, \dots\}$$

$$\lambda \in \left\{ \frac{k-1}{2k}, \frac{k-1}{2k} + k, \frac{k-1}{2k} + 3k, \dots \right\}$$

Furthermore,  $\lambda$  can be arbitrary when  $k = \pm 2$ , and additional spectrums are permitted when  $k = \pm 3, \pm 4, \pm 5$ .

Since super-integrable system is integrable,  $\lambda$  must satisfy at least the necessary condition for integrability in order to be super-integrable. Only some subsets of  $\lambda$  values are compatible with super-integrability. Below is the main result [1].

- when  $k \geq 3$  or  $k \leq -3$  :

$$\lambda \in \left\{ \frac{k-1}{2k}, \frac{k-1}{2k} + k, \frac{k-1}{2k} + 3k, \dots \right\}$$

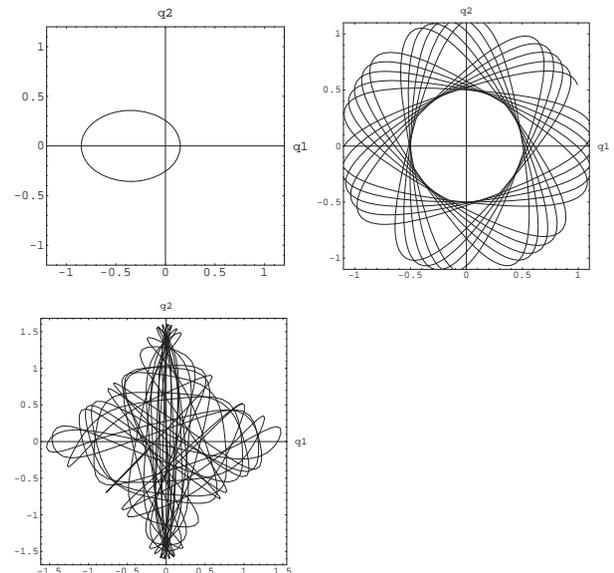
and additional spectrums when  $k = \pm 3, \pm 4, \pm 5$ .

- $k = 2$  :  $\lambda = (n/m)^2$ , a square of non-zero rationals
- $k = 1$  :  $\lambda = 0$
- $k = -1$  :  $\lambda = 1$
- $k = -2$  :  $\lambda = 1 - (n/m)^2$

Application of this theorem to homogeneous cent-

rifugal potentials  $V = \pm r^k$  always gives the value  $\lambda = 1$ . This value of  $\lambda$  is compatible with super-integrability only when  $k = -1$  and  $k = 2$ . This means that system can be super-integrable only in the case of Kepler problem ( $k = -1$ ) and isotropic harmonic oscillator ( $k = 2$ ). Thus, it turns out that the classical Bertrand's theorem is just a simple example of a wide applicable theorem.

When this new condition is applied to known sequences of integrable homogeneous polynomial potentials of arbitrary degree  $k$ , it is shown that the system cannot be super-integrable whenever  $k \geq 3$ . This suggests that there are no more super-integrable systems besides known ones; harmonic oscillators with rational frequency ratios.



**Figure 1:** Typical orbits in a super-integrable system (upper-left), in an integrable system (upper-right), and in a nonintegrable system (lower-left).

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# Trojan Stars near the Galactic Center

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Young and massive stars have been found within one parsec from the Galactic center (GC) [1]. How these massive stars were brought to the vicinity of the supermassive black hole (SMBH) has been a mystery. One possible scenario is the in-situ formation in an accretion disk. Giant molecular clouds fall into the GC and form massive gaseous disks around the central BH. Stars form in the disk if it becomes gravitationally unstable and results in fragmentation [2]. Another possible scenario is the following. A star cluster formed at a few tens of pc from the GC, and then spiraled in due to the dynamical friction [3].

We performed a fully self-consistent  $N$ -body simulation in which the internal dynamics of the cluster, that of the parent galaxy, and interactions between cluster stars and galaxy stars are correctly handled. We used the Bridge code [4] to handle the interaction between the parent galaxy and the star cluster fully self-consistently. The Bridge scheme is a tree-direct hybrid scheme. In our simulation, we adopted collisions of stars in a star cluster and formation of an intermediate-mass black hole (IMBH) in the cluster.

Figure 1 shows the projected distribution of stars at the end of the run. This figure looks very similar to K-band images of the GC [5].

We found that an IMBH is formed in the star cluster and stars escaped from the cluster are captured into a 1:1 mean motion resonance with the IMBH. These “Trojan” stars are brought close to the SMBH by the IMBH, which spirals into the GC due to the dynamical friction. Furthermore, star clusters selectively carry massive stars close to the GC. Massive stars sink to the center of their parent star clusters due to the mass segregation. As a result, massive stars tend to remain in star clusters and are carried to a few parsec from the GC.

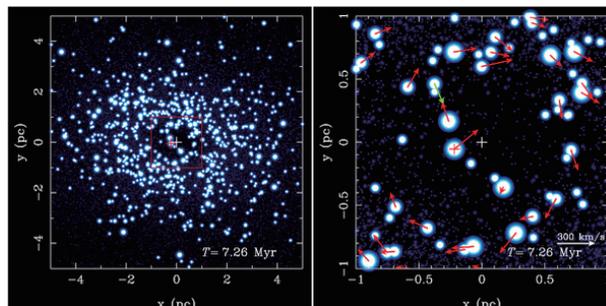


Figure 1: Projected distribution of stars at  $T = 7.26$  Myr.

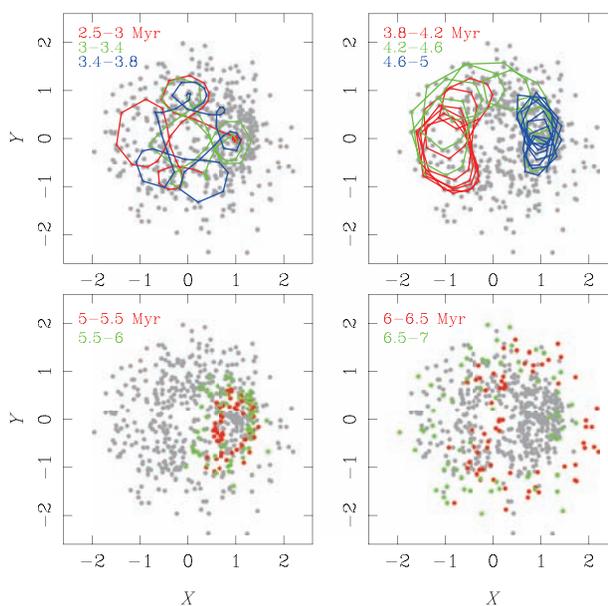


Figure 2: Orbit of a star in a rotational frame, where the IMBH is fixed at (1.0, 0.0).

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# Detection of Fluorine in the Halo Planetary Nebula BoBn1: Evidence For a Binary Progenitor Star

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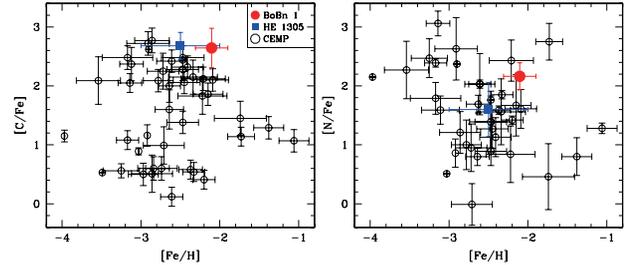
Currently, over 1,000 objects are regarded as planetary nebulae (PNe) in the Galaxy, while 14 of them have been identified as halo members from their location and kinematics. Halo PNe are interesting objects as they provide direct insight into the final evolution of old, low-mass halo stars, and they are able to convey important information for the study of low-mass star evolution and the early chemical conditions of the Galaxy.

Five objects are known as C- and N-rich ( $[C, N/O] \geq 0$ ) halo PNe with  $[Ar/H] \leq -1.7$ . The progenitors of halo PNe are thought to be  $\sim 0.8 M_{\odot}$  stars, which is the typical mass of a halo star. These C- and N-rich halo PNe, however, show signatures that they have evolved from massive progenitors. For example, these halo PNe can become N-rich, but not C-rich, if they have evolved from  $\sim 0.8 M_{\odot}$  single stars, according to the current stellar evolution models. To become C-rich PNe, the third dredgeup (TDU) must take place in the late AGB phase. However, the TDU takes place in stars with initial masses  $\geq 1.2\text{--}1.5 M_{\odot}$ . Also, current stellar evolutionary models predict that the post-AGB evolution of a star with an initial mass  $\sim 0.8 M_{\odot}$  proceeds too slowly for a visible PN to be formed. Resolving these issues would offer important results to build realistic stellar evolution and Galactic chemical evolution models. Chemical abundance analysis should help solve these issues, by revealing the signatures of internal nucleosynthesis and mixing taking place in progenitor stars. In this paper, we report the results of chemical abundance analysis of BoBn1 using VLT/UVES archive spectra [1]. BoBn1 is an extremely metal-poor ( $[Ar/H] \sim -2.1$ ) and C- and N-rich halo PN ( $[C/O] \sim +1.5$ ,  $[N/O] \sim +1.0$ ).

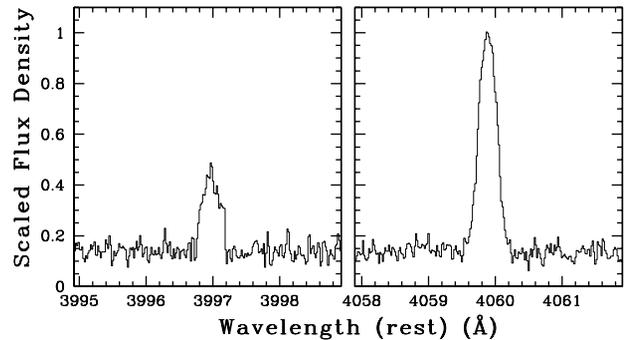
Our analysis reveals that the  $[C/Fe]$  and  $[N/Fe]$  abundances of BoBn1 are compatible with those of carbon-enhanced metal-poor (CEMP) stars (Fig. 1). C and N overabundances of CEMP can be explained by theoretical binary interaction models. We have detected two fluorine (F) lines (Fig. 2). The  $[F/H]$  of this object is  $+1.1 \pm 0.1$ , which makes BoBn1 the most F-enhanced and metal-poor PN among F-detected PNe and provides new evidence that F is enhanced by nucleosynthesis in low mass metal-poor stars (Fig. 2 in [1]). We have found that the C, N, and F overabundances of BoBn1 are comparable to those of the CEMP HE1305+0132 (HE1305, hereafter).

The chemical similarities between BoBn1 and CEMP stars suggest that this PN shares a similar origin and evolutionary history. The theoretical model by [2] demonstrated that HE1305 consists of  $\sim 2 M_{\odot}$  (primary)

and  $\sim 0.8 M_{\odot}$  (secondary) stars, and the enhanced C and F can be explained by binary mass transfer from the primary star. The C and F abundances of BoBn 1 can be explained by the  $\sim 2 M_{\odot}$  and  $\sim 0.8 M_{\odot}$  binary model of [2]. Also, the issue of evolutionary time scale can be resolved if BoBn 1 has evolved from a binary and experienced mass transfer in the mid-course of its evolution.



**Figure 1:** Location of BoBn1 on the diagrams of  $[C/Fe]$  (left) and  $[N/Fe]$  (right) vs.  $[Fe/H]$ . For BoBn1, we used Ar instead of Fe.



**Figure 2:** Emission line-profiles of  $[F\text{ IV}] \lambda 3996.92 \text{ \AA}$  (left) and  $[F\text{ IV}] \lambda 4059.90 \text{ \AA}$  (right) of BoBn 1.

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# Project Milkyway: Shock-Induced Starburst at a Collision Interface during the First Encounter of Interacting Galaxies

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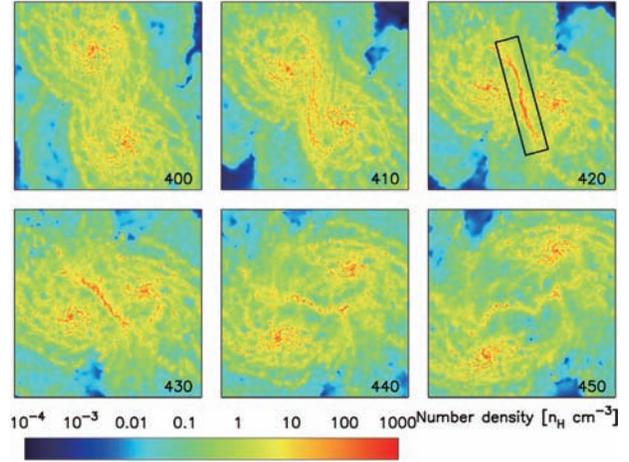
DAISAKA, Hiroshi  
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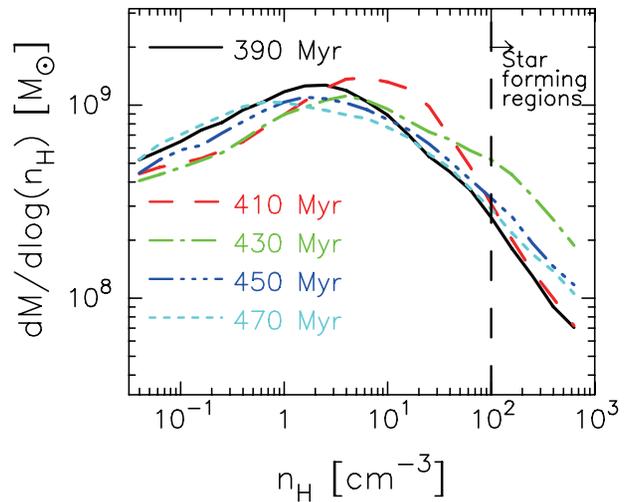
The goal of “Project Milkyway” is to reveal the formation process of galaxies, especially that of the Milkyway galaxy, by the state-of-the-art simulations. For this purpose, we have developed (i) PC-clusters whose computational nodes consist of ordinary CPUs (AMD Opteron) and special purpose computer, GRAPEs, which accelerates the calculation of gravity, and (ii) a special code for parallel simulations with GRAPEs, ASURA. The latest version of ASURA employs a Phantom-GRAPE library, which is a software emulator of GRAPE tuned with SIMD operations.

We performed high resolution simulations of galaxy-galaxy collisions with ASURA on Cray XT4 at CfCA [1]. In this simulation, we resolved a multiphase nature of the interstellar medium and adopted a realistic conditions for star formation [2]. Previous simulations of galaxy-galaxy merging, on the contrary, did not employ a lower part of a cooling function ( $< 10^4\text{K}$ ) and considered relatively diffuse regions ( $n_H > 0.1\text{cm}^{-3}$ ) as star forming regions. These simulations could not resolve neither wide spread starbursts nor star cluster formation during the first encounter.

Our simulations successfully resolved not only a wide spread starburst during the first encounter but also the formation of star clusters. The starburst took place along a filamentary region, which was induced by hydrodynamical shocks. This starburst rapidly quenched because of the effect of supernovae. After  $3 \times 10^7$  yr later, we found about ten self-bound star clusters along the filament. Interestingly, these star clusters have no dark matter. Our results are consistent with the observational properties of colliding galaxies.



**Figure 1:** The evolution of gas density during the first encounter. A large and dense filament appears at the collision interface.



**Figure 2:** The evolution of gas mass as a function of density during the first encounter. Gas, which was compressed by the large scale shock, satisfies star formation conditions and a burst-like star formation took place.

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# Limit of primordial magnetic field from the matter power spectra

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Recent observations show that magnetic fields in clusters of galaxies are  $0.1\text{--}1.0\mu\text{G}$ . Many authors also suggest the existence of a primordial magnetic field (PMF) from this fact, and study effects of PMF in the early universe, such as the cosmic microwave background (CMB) [1, 2, 3], or the matter density field [4, 5]. In this article we introduce effects of PMF on a matter power spectra of density fluctuations in early universe, and show that the observational constraints on cosmological density fluctuations, as parameterized by  $\sigma_8$ , lead to strong constraints on the amplitude and spectral index of the PMF [4, 5].

Pressure and tension of the PMF affect the evolution of density fluctuations of ionized baryons. Since baryon and cold dark matter (CDM) interact with each other through gravitational force, the PMF also affects the evolution of density fluctuations of the CDM. Eventually, the total matter density fluctuations are affected by the PMF [4, 5].

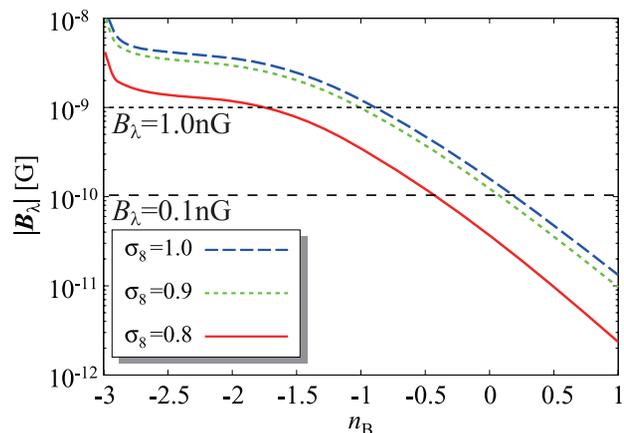
The alternative normalization parameter  $\sigma_8$  is the root-mean-square matter density fluctuation in a comoving sphere of radius  $8h^{-1}\text{Mpc}$ . It is given by a weighted integral of the matter power spectrum. We can study the physical processes of density field fluctuations on cosmological scales within the linear regime to determine  $\sigma_8$ . From this we can obtain strong constraint on the PMF parameters by numerically calculating  $\sigma_8$  under the influence of PMF effects.

Since  $\sigma_8$  is affected by other cosmological parameters, we should consider the degeneracy between the PMF and other cosmological parameters as mentioned above. Fortunately, these cosmological parameters are constrained by recent CMB observations on larger scales ( $l < 1000$ ), while it was shown in our previous work [1] that the effect of the PMF mainly affects the CMB anisotropies on smaller scales ( $l > 1000$ ). Hence, we expect that the degeneracy between the PMF parameters and the other cosmological parameters is small. For this reason in the present analysis we are justified to fix the other cosmological parameters at their best fit values.

We adopt the constraint that PMF parameters giving  $\sigma_8 > 1$  are excluded by observations. Figure 1 shows that a PMF amplitude of  $B_\lambda \gtrsim 1\text{ nG}$  is excluded when  $n_B > -0.9$ . Furthermore, PMF amplitudes of  $B_\lambda \gtrsim 0.1\text{ nG}$  are excluded when  $n_B > 0.2$ . The magnetic field strength in galaxy clusters is  $\sim 1\mu\text{G}$ . Therefore, if isotropic collapse is the only process which amplifies the magnetic field strength, the lower limit to the PMF is  $\sim 1\text{ nG}$  at  $z \sim 0$ . Hence, we

can obtain a strong constraint on this PMF evolution model for a PMF spectral index in the range  $n_B < -0.9$ .

It is very important to constrain the PMF as precisely as possible since a magnetic field affects the formation of large scale structure, for example, magnetic pressure delays the gravitational collapse. Recently, several observations of the CMB anisotropies and polarizations for higher multipoles  $l$ , e.g. via the *Planck*, *QUIET*, and *PolarBear*, are progressing as planned. If we combine these future plans and our works, we will be able to limit parameters of the PMF more accurately and develop studies of evolution and formation of large scale structure with the PMF.



**Figure 1:** Curves of constant values for  $\sigma_8$  in the parameter plane of PMF amplitude  $B_\lambda$  vs. spectral index  $n_B$ . All curves are as indicated in the legend box in the figure.

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# Coronal Plasma Motions near Footpoints of Solar Active Region Loops

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(NAOJ)

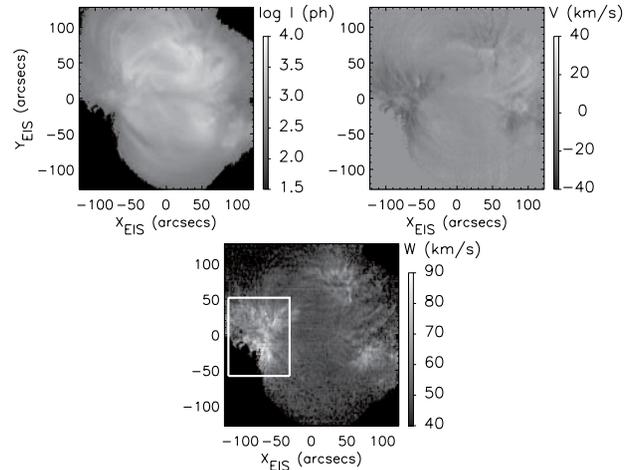
YOUNG, P. R.  
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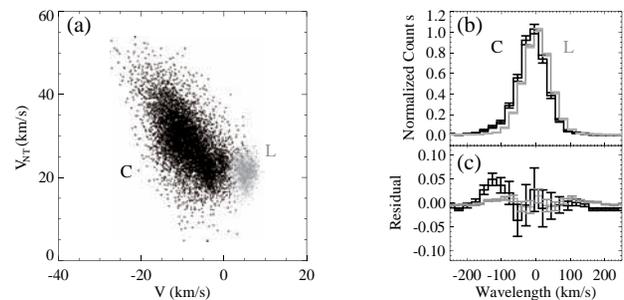
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The Sun has a tenuous  $10^6$  K corona located above the photosphere of 6000 K. We try to fully understand how it is heated to such a hot temperature. We are investigating the heating mechanism of the solar activeregion corona from observations with the *Hinode* EUV imaging spectrometer (EIS) [1]. Initial results of this study [2] are presented here.

Figure 1 shows an observation of a solar active region, located at the disk center, with EIS in the  $\text{Fe XIV}$  emission line at 27.4 nm. The formation temperature of the emission line is about  $1.8 \times 10^6$  K. In this figure a single Gaussian function is applied for fitting the emission line profile to evaluate the spatial distribution of the line intensity  $I$ , Doppler velocity  $V$ , and line width  $W$ . The enhanced line width is wider than the thermal Doppler width and the enhanced component is called nonthermal velocity  $V_{\text{NT}}$ . Subsonic upflow motions of  $10\text{--}20\text{ km s}^{-1}$  and enhanced nonthermal velocities have been found near the footpoints of the active region loops. When the active region is observed near the west limb, both Doppler velocity and nonthermal velocity at the same footpoint region essentially decrease. There is a strong correlation between the upflow velocity and the nonthermal velocity (Fig. 2a). Significant deviations from a single Gaussian profile with a weak emission enhancement are found in the blue wing of the line profiles for the upflows (Fig. 2b and 2c). These suggest that there are unresolved high-speed upflow components. Bluewing enhanced profiles show fast upflow components of heated plasmas whose velocities along the loop structure are nearly the coronal sound speed. The *Hinode* EIS has been revealing the characteristic motion of heated coronal plasmas.



**Figure 1:** Solar active region located at the disk center of the Sun observed with *Hinode* EUV imaging spectrometer in the emission line  $\text{Fe XIV}$  at 27.4 nm. Line intensity  $I$ , Doppler velocity  $V$ , and the line width in full width at half maximum  $W$  are shown. Red (blue) shift is indicated in plus (minus) sign for the Dopplergram.



**Figure 2:** (a) Correlation between Doppler velocity  $V$  and nonthermal velocity  $V_{\text{NT}}$  in  $\text{Fe XIV}$  at 274 Å. (b) Blue-side enhanced line profiles at a footpoint of a coronal loop and (c) residual from a single Gaussian fit. C (L) in black (gray) shows the line profile in the disk-center (limb) observation.

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# Subaru Witnesses the Assembly of Distant Clusters 8 Billion Years Ago

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Galaxies are dynamic star systems that have enormous numbers (millions to billions) of stars in wide varieties of colors and shapes. Some galaxies have blue beautifully knotted spiral arms, while others are red diffuse light without arms. As most celestial objects, galaxies are NOT uniformly distributed throughout the Universe. They tend to flock together and form a “cosmic web”, which represents the large-scale structure of the Universe. At intersections within this webwork, astronomers observe massive concentrations of galaxies called Galaxy Clusters.

Astronomers have observed that the colors and shapes of galaxies change depending on their location. In clusters, where so many galaxies live together, galaxies tend to be red and elliptical. On the other hand, galaxies living alone are predisposed to be blue and spiral shaped. Astronomers do not yet have a clear idea why galaxies change their colors and shapes contingent on their location and their surrounding environments. Nonetheless, one way to explain this phenomenon is to observe galaxies at great distances that represent periods long ago in the past. Simply stated, the Universe is like a time machine - as we look deeper into the Universe, we go further back in time. Using the Subaru Telescope, astronomers explored answering the question, “How do galaxies change their colors and shapes over time?”

Recently, an international team of astronomers used Subaru to observe two huge groups of galaxies, galaxy clusters, in the distant Universe. One cluster (RDCS0910) was 8 billion light years away while another (RDCS1252) was 8.5 billion light years distance. These two clusters are among the most distant clusters known to date. Galaxies at such a long distance appear very faint, but Subaru is sensitive enough to collect the dim light from such great distances.

Galaxy shape, color, and distribution in these structures are keys to solving the riddles of galactic evolution in the distant reaches of Universe. Using Suprime-Cam, the astronomers discovered a lot of galaxies within the clusters and found that the clusters are surrounded by small groups of galaxies. Not all these small groups were confirmed to lie at exactly the same distance as the clusters, but this discovery of prominent large-scale structures in the distant Universe is a seed to understanding overall galactic evolution.

Small galaxy groups at this distance ( $z = 1.2$ ) have previously not been studied in detail, and, as such, are very interesting objects for astronomers. The small galaxy groups may get closer to the nearby larger cluster through gravity and eventually will collide a few hundred million to a few billion years from now. In fact, collisions and mergers are considered primary factors in galactic cluster growth. That is, clusters grow by accreting small galaxy groups around them. It should not be surprising if colors and shapes of galaxies in small groups change during these large-scale collisions.

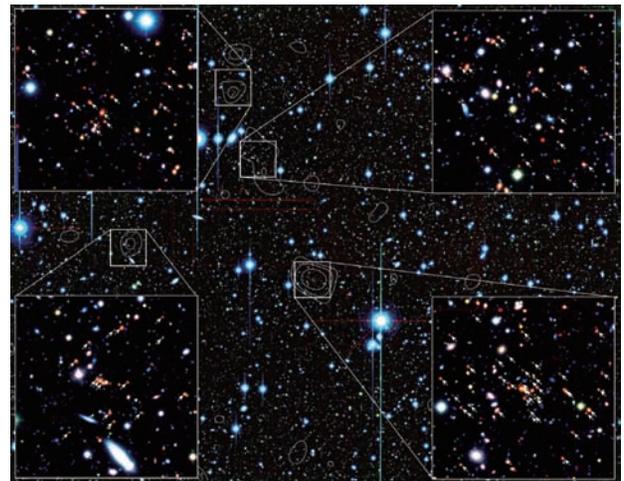
Looking at the small groups of galaxies in Figures 1, the first thing noted is the fantastic scale. The large zoom-in boxes are 3 million light years across, indicating that the galaxy groups are huge. Looking closely at the boxes, many reddish (well, orange-ish) galaxies are observed. These galaxies are so far away from us that it is impossible to spatially resolve their shapes. However, it is clearly visible that the small groups of galaxies are already dominated by red galaxies at least 8 billion years ago.

The fundamental question for this study was “What do these observations mean?” As mentioned earlier, galaxies living alone tend to have blue colors, while galaxies in clusters have red colors. Galaxies likely change their colors from blue to red during the course of relocation and structural growth. Nearby galaxies pulled together by gravity, form a group of galaxies, and eventually collide with other groups and form large clusters. At some point in this structural growth, blue galaxies become red. The team found in their study that small groups are already dominated by red

galaxies before they collide with big clusters. Astronomers have long believed that large clusters are the key environment changing colors and shapes of galaxies, but the findings at Subaru now point astronomers at small groups. Galaxies become red in groups before they get together in large clusters - small groups are the key environment for these changes in shape and color.

Astronomers know several physical processes that can affect colors and shapes of galaxies. Among them, collisions and mergers between galaxies (not between groups) are the ones that work very efficiently. Because groups have many galaxies in them, galaxies can attract each other by gravity and eventually collide and merge. In fact, a computer simulation has shown that a collision between two blue galaxies leaves a single red galaxy. The Universe 8 billion years ago shows us that collisions with other galaxies might be the reason why galaxies change their colors and shapes depending on environment around them. While the astronomers observed just a tiny fraction of the Universe, galaxy collisions alone cannot explain everything and some mysteries remain. Subaru will continue to observe at great distances and at various wavelengths to get a better understanding of galactic transformation.

This work is published in two papers [1] and [2].



**Figure 1:** This image shows a cluster of galaxies 8.5 billion light years away in the constellation Hydra. The pseudo-color picture is made with images taken with Suprime-Cam on the Subaru and WFCAM on the United Kingdom Infrared Telescope (UKIRT). The picture covers  $35' \times 28'$  of the sky. This area corresponds approximately to 60 million  $\times$  45 million light years. The contours show how densely galaxies flock together 8.5 billion years ago. The zoom-in views on the corners show small groups and a big cluster of galaxies. The one on the bottom-right is a known large cluster. The newly discovered small groups shown in the other corners surround this massive cluster. Each panel illustrates 3 million light years across. The arrows point at cluster member candidates.

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# Evidence for a Companion to BM Gem, a Silicate Carbon Star

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A small number of carbon stars among those identified in the optical were found to show silicate dust emission features in the IRAS Low Resolution Spectra in 1986 [1, 2]. Silicate dust forms in oxygen-rich environment where the oxygen abundance exceeds the carbon abundance, while it does not in carbon-rich environment where the opposite is the case.

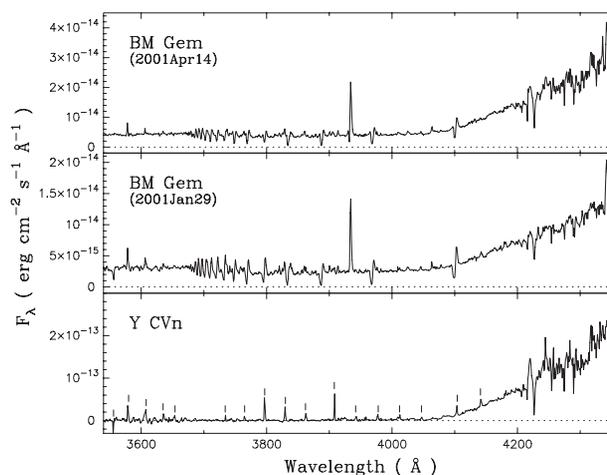
Contemporary stellar evolution theories tell us the occurrence of repetitive runaways of the He shell burning (“thermal pulses”) in AGB stars, which are descendants of low- or intermediate-mass stars. In every thermal pulse synthesized carbon is carried to the outer layer, which eventually causes an originally oxygen-rich M-type star evolve into a carbon star.

Two explanations were proposed for the origin of silicate carbon stars. One was that silicate carbon stars were those that had newly evolved from mass-losing, M-type AGB stars so that we could see remnant silicate dust around them in the infrared emission features [2]. The other was that silicate carbon stars were binary systems, each with a mass-losing, M-type AGB star companion that was enshrouded in a thick silicate dust shell that shone in the infrared features [3]. These two propositions were, however, both rejected by follow-up studies.

In more than ten years two new pictures were proposed for the origin of silicate carbon stars. Both postulate a low mass, low luminosity companion and invoke an unknown mechanism that is driven by the binary interaction and gives rise to either a circum-binary reservoir [4] or a circum-companion disk [5] that retains silicate grains captured when the primary was an M-type mass-losing star. However, there has been no observational evidence for the presence of a companion in silicate carbon stars.

We obtained high resolution, high sensitivity ultraviolet to visible spectra of the brightest silicate carbon star BM Gem with the high dispersion spectrograph HDS on the Subaru telescope. In the spectra we discovered ultraviolet continuum emission that is unusual for carbon stars (Fig. 1). The emission has spectral features typical for an ionized gas. Further, we detected the Balmer series lines, each of which has a P Cygni-type profile that indicates the presence of a high velocity outflow as fast as  $400 \text{ km s}^{-1}$  and is time variable [6]. No carbon star has been known that shows such spectral features. It is difficult to imagine a new radiative mechanism that produces the observed spectral features in a cool carbon star. On the other hand, we have found that the continuum luminosity is plausibly explained if we interpret the

spectral features as emission from an ionized gas that is caused by the heat from the gravitational potential energy release during the accretion of a part of the primary’s mass-loss flow by the companion. Furthermore, we have noticed that almost identical spectral features have been observed in Mira B, the companion to the M-type AGB star Mira A (*o* Cet). With these observational results we have concluded the presence of a low-luminosity, low-mass companion to the silicate carbon star BM Gem.



**Figure 1:** Ultraviolet to visible spectra of BM Gem observed with HDS on Subaru telescope (top and middle, different dates). A spectrum of Y CVn, which looks similar to BM Gem in the visible region, taken with the same setting for comparison (bottom). Small vertical ticks in the bottom show the positions of artifacts that mimic emission lines.

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# Discovery of strange filamentary structures (“fireballs”) around a merger galaxy in the Coma cluster of galaxies

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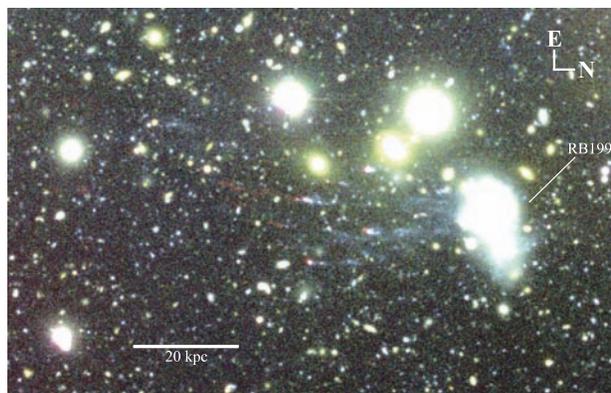
Clusters of galaxies are ideal experimental laboratories for studying environmental effects on evolution of galaxies. Several environmental effects which explain the evolution of cluster galaxies have been proposed so far and many observational studies have been made to identify the dominant environmental processes in clusters. However, it is still not clear what kinds of physical processes are dominant for rapid galaxy evolution in clusters.

Recently, two peculiar galaxies were found near the central regions of two rich clusters at  $z \approx 0.2$  [1]. These galaxies have disturbed morphology and are associated with many small blue blobs extending toward the onesides of galaxies. It is suggested that these phenomena may be snapshots of transformation of spiral galaxies to S0 galaxies and may give insights to the origin of faint end population of clusters [1].

Here we report a discovery of a strange complex of narrow blue filaments and knots (Figure 1; its morphology resembles a complex of “fireballs (large meteors)”) which extends up to 80 kpc to the south from a disturbed E+A galaxy, RB199, in the Coma cluster using Subaru Telescope [2]. The morphology of the structure is very similar to that of the phenomena found in the two  $z \approx 0.2$  clusters. This structure may be a nearby counterpart of the transition galaxies found in the clusters at  $z \approx 0.2$ .

The narrow blue filaments extend in straight shapes toward the south from the galaxy and several bright blue knots are located at the south ends of the blue filaments (Figure 1). The  $R_C$  band absolute magnitudes, half light radii and estimated masses of the bright knots are  $\sim -12$  –  $-13$  mag,  $\sim 100$  –  $200$  pc and  $\sim 10^{6-7} M_\odot$ , respectively. Strong  $H\alpha$  emission is associated with the bright knots. Long narrow  $H\alpha$  emitting filaments are connected at the south edge of the knots. The average color of the “fireballs” is  $B - R_C \approx 0.5$ , which is bluer than RB199 ( $B - R = 0.99$ ), suggesting that most of the stars of which the “fireballs” is composed were formed within several times  $10^8$  yr. There is no  $H\alpha$  emission in the narrow blue filaments, whereas strong  $H\alpha$  and UV emission are emitted at the bright knots. These characteristics indicate that star formation has been recently ceased in the blue filaments and is now in progress in the bright knots. The gas stripped by some mechanism from the disk of RB199

may be traveling in the intergalactic space, forming stars and leaving the formed stars along its trajectories. Most plausible mechanism of forming the “fireballs” is ram pressure stripping by high speed collision between the galaxy and the hot intra-cluster medium. The “fireballs” may be a snapshot of a formation process of diffuse intra-cluster population or halo star population of a cluster galaxy.



**Figure 1:** False color ( $B$  band: blue,  $R_C$  band: green,  $H\alpha$  NB: red) image of area around RB199. North is right and east is top. RB199 is located at the right side of the image. Several blue filaments extend from RB199 to the center of this image and bright blue knots are located at the tips of the filaments. Red  $H\alpha$  emission is associated with the bright knots.

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# Lunar Global Topographic Map by the KAGUYA Laser Altimetry

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The main orbiter of Japanese Lunar probe KAGUYA (SELENE) which was launched on 14 September, 2007 (UT; the same hereafter) has successfully closed its observation by the controlled impact onto the Lunar surface on 10 June, 2009. We have constructed a global and precise topographic map of the Moon for the first time using the KAGUYA laser altimeter (LALT) topographic data (Fig. 1) [1]. LALT has obtained about 10 million range data during the nominal observation period from 30 December, 2007 to 27 October, 2008. The along and cross track resolution of the geo-located points in the equatorial region is 1.5 km and less than 10 km, respectively. The mean number density of the geo-located points is about  $1/(200\text{ m})^2$  in both polar regions.

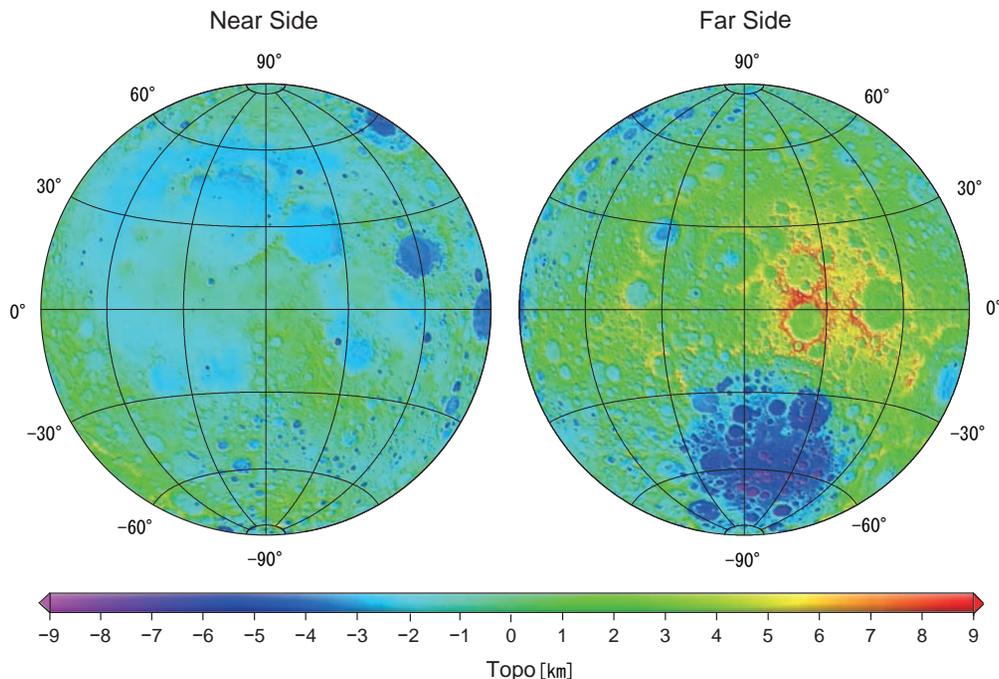
It has been derived that Lunar mean radius is 1737.15 km and the offset between the center of mass (COM) and the center of figure (COF) is 1.93 km by the spherical harmonic analysis of 6.67 million lunar topographic data obtained by 31 March, 2008. The maximum topographic difference is found to be 19.81 km, where height reference is a sphere whose radius is 1737.4 km centered on the COM. The spherical

harmonic analysis has also revealed that the amplitude of the topography for the horizontal scale less than 180 km shows 2 or 3 times larger spectrum than the previous estimation. This may suggest that the small scale topography is not compensated and more rigid than other terrestrial planets [1].

In the Lunar polar regions, volatile materials such as water ice are expected to be buried within the regolith in the deep crater floor where the sun lit rate is zero due to the low solar elevation; while the high place where the sun lit rate is high such as on the crater rim is promising candidate site for the Lunar base. We have calculated and mapped the solar illumination condition on the both polar regions for more than 85 degrees based on the LALT data. The results are as follows; (i) the sun lit rate is 89% and 86% at most for the north and south region respectively, (ii) the permanent shadow area for more than 87.5 degrees is 844 km<sup>2</sup> and 2751 km<sup>2</sup> for the Lunar north and south polar region, respectively [2].

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**Figure 1:** Lunar global topographic map derived from LALT altimetry data from December 30, 2007 to June 30, 2008 shown in Lambert equal area projection. Lunar coordinates are based on the mean Earth/polar axis system. Reference of the height is a sphere whose radius is 1737.4 km and origin is set to the center of mass.

# Precise wavefront correction with an unbalanced nulling interferometer for exo-planet imaging coronagraphs

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In 1995, an exoplanet was discovered by a Doppler method for the first time. Since then, about 350 exoplanets were detected almost by indirect methods while a few very bright exoplanets were recently detected by direct coronagraphic imaging in 2009. Spectroscopic analysis of exoplanet light is one of the most important theme in modern astronomy.

The direct imaging and the spectroscopy of Earth-like exoplanets is disturbed by bright diffracted light and speckle noise of a central star, and the required dynamical range is  $10^9 \sim 10^{10}$  on small angular separation of a few  $\lambda/D$ . Here an input wavefront quality of approximately ten thousandths of a wavelength rms is required to reduce the speckle noise. We studied pre-optics with nulling interferometer [1], and complex compensation after interference [2], and, as a combined extension of them, developed a method presented in the following.

We propose a novel method based on a 4-stage configuration of a coronagraph system (Figure 1) [3, 4, 5], in which a very important role is achieved by a middle-stage optics setup (UNI-PAC) that behaves partly as a low-efficiency coronagraph, and partly as a high-sensitivity wavefront aberration compensator (phase and amplitude).

A first adaptive optics (AO) system compensates for the wavefront aberrations of mainly the primary mirror of a telescope with a quality of  $\lambda/1000$  rms. Second, an (intensity-) unbalanced nulling interferometer (UNI) performs a rejection of part of the wavefront electric field. Then, the input aberrations of the recombined output wavefront are magnified to  $\lambda/100$  rms. Because of the unbalanced recombination scheme, aberrations can be free of phase singular (zero intensity) points and can therefore be compensated by a downstream phase and amplitude correction (PAC) AO system, using two deformable mirrors with a quality of  $\lambda/1000$  rms. In the image plane, the central star's peak intensity and the noise level of its speckled halo are reduced by the UNI-PAC combination: the output-corrected wavefront aberrations from the UNI-PAC can be interpreted as an improved compensation of the initial incident wavefront aberrations, implying that the wavefront quality was "virtually" improved to  $\lambda/10000$  rms beyond the AO system capabilities.

Finally the residual star light is reduced by a coronagraph and the speckle level is kept as low as the

planet intensity even with a wavefront quality of  $\lambda/1000$  rms. A significant advantage of using the UNI-PAC is the relaxation of the specification for the entire optics of the coronagraph system to  $\lambda/1000$  rms surface error quality, which is an important step towards developing a terrestrial-planet detection system in a cost-effective way.

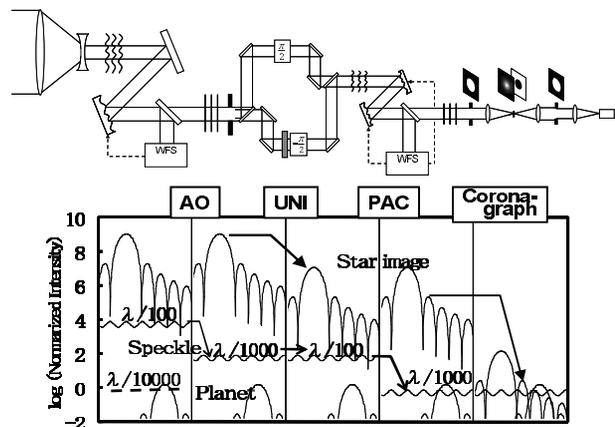


Figure 1: Schematic of the 4-stage coronagraph optics including the UNI-PAC method [3, 4, 5]. Wavefront corrections are made twice at the first AO and the PAC AO system, and star light suppressions are also made twice at the UNI and the coronagraph.

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# TCP-AIC: Tandem common-path achromatic interfero-coronagraph

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New scheme for an achromatic interfero-coronagraph, incorporating two common-path interferometers in tandem (Figure 1), is proposed to attain a deeper nulling for an extended incoherent stellar disk. The predicted performance improvement, by using TCP-AIC has the possibility to reach a  $10^{10}$  achromatic coronagraphic contrast for an almost  $10^{-2} \lambda/D$  effective source size (stellar disk size).

Direct imaging of extra-solar Earth-like planets, using stellar coronagraphs and nulling interferometers, requires the realization of star-to-planet intensity contrasts, as high as  $\sim 10^{10}$  in the visible wavelength range and  $\sim 10^6$  in the mid-IR range. For the detection of a faint planet signal, such optical systems require an extremely accurate wavefront control and a highly stable system configuration. Currently an achromatic  $\pi$  shifter with sufficient inner working angle remains a key device for nulling interferometers and an AIC coronagraph concepts.

The AIC realizes coronagraphic nulling by shifting the geometric phase in a common-path interferometer. This common-path configuration makes the detection system highly immune to environmental disturbances. The coronagraphic image is a centro-symmetrical pair of companion images.

For a monochromatic laser point source of 633 nm, single CPAIC has demonstrated light suppression with a stable coronagraphic contrast of  $10^6$  at a  $1 \lambda/D$  off-axis distance [1].

For a star-disk diameter of  $10^{-2} \lambda/D$  and an off-axis planet separation distance of  $1 \lambda/D$ , nulling contrast of only  $10^4$  can be achieved with an AIC, which condition corresponds to a Sun-Earth system at 5 pc distance using a 1-meter-class diameter telescope operated at a  $1 \mu\text{m}$  center wavelength. The resulting non-null residual starlight remains at 2 orders greater than the target planet contrast of  $10^6$ , which is required for the mid infrared range. Residual visible light remains 6 orders greater to the  $10^{10}$  contrast required.

To cope with the problem of insufficient spatial coherence caused by an extended source, we proposed the TCP-AIC which use two interferometers cascaded in tandem, effectively making four-beam interference (Figure 2). The TCP-AIC was started experimentally to realize a  $10^{10}$  coronagraphic contrast [2].

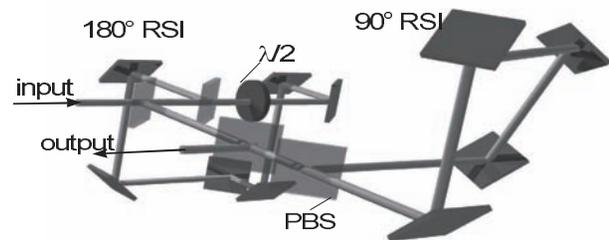


Figure 1: Schematic of the tandem common-path achromatic interferometer coronagraph.

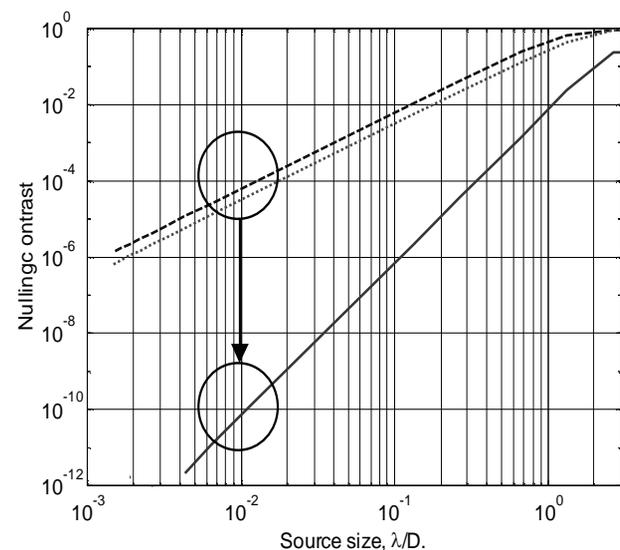


Figure 2: Nulling contrast of the single and tandem CPAIC.

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