

Chromospheric Structure Derived from Flash Spectra of the Total Solar Eclipse

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(Received September 30, 2002)

Abstract

A chromosphere model for the analysis of emission lines in a flash spectrum is constructed. Emission gradients of metallic and Balmer lines in flash spectra give height distributions of the total hydrogen and the product of electron and proton density in the high chromosphere, respectively. The derived distributions imply the presence of “spicule” structure which has a filling factor of 0.05 at 4,000 km above the base of the chromosphere. They explain the averaged eclipse curves of Ca II H and K, and H ϵ line profiles observed in a 1958 flash spectrum and the Balmer and Sr II emissions observed in a 1962 flash spectrum. Their excitation and ionization seem to match the radiation field of the chromosphere. They are applied to 24 Ca II H and K spicules in the higher chromosphere observed during the 1958 eclipse. The analysis shows that they have a turbulence of 22 km/s on the average and 19 of them are thinner than 2,000 km. The Ca II H and K, and H ϵ emissions of the active region observed during the 1958 eclipse are enhanced mainly by the increase of their source functions due to an increase in their excitation temperatures.

Key words: Chromosphere, Flash spectra, Spicules, Total solar eclipse.

1. Introduction

Physical parameters of the solar chromosphere are determined in two ways. One is the way that they can reproduce the solar UV spectrum. The other uses flash spectra of the total solar eclipse. Height distributions of the total hydrogen density and of the electron times proton density are shown in figure 1b and d. The thin curve distributions are the latest result from fitting the UV spectrum (Fontenla et al. 1993). In this determination the hydrostatic assumption confines the vertical extension and the chromosphere extends to around a height of 1800 km by the usual lifting forces of pressure and turbulence against gravity. In the second case the hydrostatic assumption is not necessary because the flash spectrum is a “side view” of the chromosphere. The thick curve distributions extending to the higher levels in figure 1b and d are the results of this study which will explain the emission lines visible in the flash spectrum of these levels.

The analysis made here is rather coarse and uses two simplifications: 1) The emission gradients of the eclipse curves reflect the density gradient and 2) Each analyzed emission line has its own constant excitation temperature, i.e., constant source function and ionization temperature throughout the chromosphere. In the low density chromosphere the relative populations among the atomic levels are mainly governed by radiative processes which have far longer scale heights than the density. This makes the excitation and ionization temperatures of each emission line change less against height and the chromospheric emission gradient equal to the density gradient. Zirker (1958) shows this in the case of metallic lines. In the higher density atmosphere collisions become more effective and change the populations of the atomic levels, which may decrease the temperature scale heights. The emission gradi-

ents will thus deviate from the density gradient.

Section 1 gives the derivation of the above distributions from emission gradients of metallic and Balmer lines in the high chromosphere, which extend the low chromosphere model (Hiei 1963) obtained with the continuous spectrum. The derived model reveals the spicule structure in the high chromosphere. Section 2 is an application of the model to Ca II H and K, and H ϵ line profiles obtained during the 1958 total solar eclipse (Suemoto and Hiei 1959, 1962). The same analysis is made for Balmer and Sr II lines using 1962 eclipse data (Dunn et al. 1968) in section 3. Section 4 studies 24 Ca II H and K spicules observed during the 1958 eclipse. Attached appendices give line profile data of Ca II H and K, and H ϵ observed during the 1958 eclipse and other information necessary for the text.

The reference coordinate system for the flash spectrum data used in this paper is depicted in figure 2 (ref. Thomas and Athay 1961). The whole area outside the moon’s limb contributes to the emission in the slitless flash spectrum.

Numerical constants in the text are described in cgs-units.

2. Model of the Chromosphere

An empirical model of the low chromosphere has been obtained by Hiei (1963) with the use of a continuous spectrum of the 1958 total solar eclipse. The higher chromosphere, where many emission lines are visible during the eclipse, may be studied using their emission gradients. These are measured where the emission lines are weak and are closely related to the gradient of the emitting atom density, $n(h)$. Let $\beta'(h)$ be the density gradient of the emitting atom, then

$$\beta'(h) = -\frac{d \ln n(h)}{dh} \quad (1)$$

By integration

$$n(h) = n(0) \exp \left[-\int_0^h \beta'(h) dh \right] \quad (2)$$

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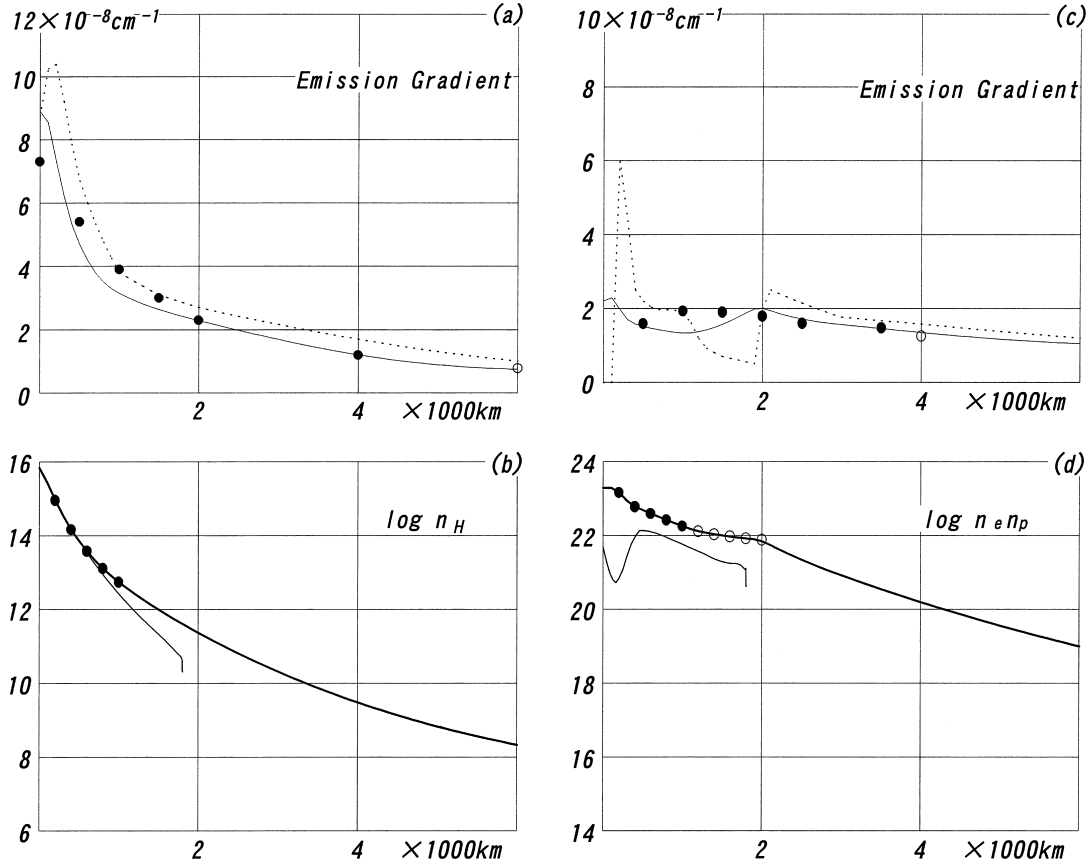


Fig. 1. Emission gradients and density distributions of the chromosphere. (a) Filled circles are emission gradients of the metallic lines compiled by Unsöld (1955). The open circle is obtained from the Ca II H and K lines observed during the 1958 eclipse. The solid curve is the emission gradient calculated from the dotted curve (the density gradient of hydrogen). (b) Filled circles are the densities of hydrogen from Hiei (1963). The thick curve is the hydrogen density distribution determined in this work. The thin curve is from the FAL-C model (Fontenla et al. 1993). (c) Filled circles are the emission gradients of the Balmer lines from the 1962 eclipse. The open circle is obtained from H ϵ during the 1958 eclipse. The solid curve is the emission gradient calculated from the dotted curve (the gradient of $n_e n_p$ distribution). (d) Filled and open circles show the $n_e n_p$ distribution from Hiei's work (Hiei 1963). The open circles are the values from the assumption, $n_e = n_p$. The thick curve is the $n_e n_p$ distribution determined in this work. The thin curve is from the FAL-C model (Fontenla et al. 1993). The height of the FAL-C model is shifted down by 400 km to reduce to the one used here, the origin of which is at the base of the chromosphere.

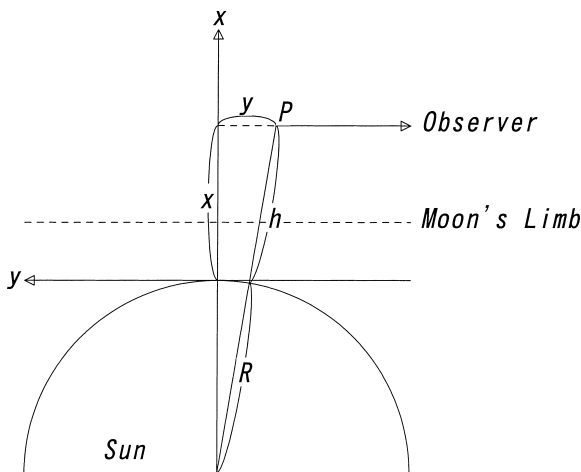


Fig. 2. The coordinate system (x, y) for the flash spectrum data. R : Sun's radius, h : height of an emitting atom at P from the base of the chromosphere. All the emitting atoms outside the moon's limb contribute to the emission of the slitless flash spectrum.

where $n(0)$ is the density at the base of the chromosphere, which is defined by the visible sun's limb. The number of the emitting atoms in the line of sight, $N(x)$, at the projected distance from the sun's limb, x , is

$$N(x) = \int_{-\infty}^{\infty} n(h) dy \quad (3)$$

where

$$h = x + \frac{y^2}{2R} \quad h \ll R \quad (4)$$

The total number, $N'(x)$, contributing to the intensity is

$$N'(x) = \int_x^{\infty} N(x) dx \quad (5)$$

Therefore, the emission gradient, $\beta(x)$, determined from weak emissions, should be

$$\beta(x) = -\frac{d \ln N'(x)}{dx} \quad (6)$$

If an approximation, $\beta'(h = x + \text{emission scale height}) \sim \beta(x)$ (e.g., van de Hulst 1953), is applied, equation (6) can be calculated with the use of equations (1)–(5). The resultant $\beta(x)$ is compared with the observed one, and, by iteration, the

final $n(h)$ is obtained from equation (2).

According to the study by Zirker (1958), the emission gradient of the metallic lines is equal to the total hydrogen density gradient. If in equation (1) n is replaced with the total hydrogen density, n_H , equation (6) provides the emission gradient of the metallic lines. The compiled values by Unsöld (1955) are adopted as the observed emission gradients. The gradient at $x = 6,000$ km is added from the Ca II H and K eclipse curves derived from 1958 eclipse data (see figures 4a and b, ref. appendix 1). Zirker's density gradients are the emission gradients assigned to special higher heights where the emission has the intensity of the "plate limit" (Zirker 1958) and were not used. The final iteration is shown in figure 1a. The discrepancy between the observed and calculated distribution is seen below 2,000 km, where the emission gradient of the metallic lines may not be equal to the density gradient. Hiei's hydrogen density is adopted below 1,000 km.

Balmer line intensities are proportional to the number of emitting atoms when the atmosphere is thin. Let n_j be the number density of the j -th level of hydrogen. Then by the Saha-Boltzmann equation it is

$$n_j = j^2(n_e n_p) T^{-3/2} \exp \left[\frac{X_{ion} - X_j}{kT} \right] / C_1 \quad (7)$$

where

$$C_1 = (2\pi mk)^{3/2} / h^3 = 2.4147 \times 10^{15},$$

k : Boltzmann constant,

n_e : electron density,

n_p : proton density,

T : temperature,

X_j : excitation potential of the j -th level,

X_{ion} : ionization potential of hydrogen.

This equation shows that the gradient of the Balmer line intensity is equal to that of the cross product, $n_e n_p$, as long as the temperature gradient is negligible. This will occur in the high chromosphere where the density is low and the radiation is dominant, and equations (1)–(6) are applicable when $n(h)$ is replaced with $n_e n_p(h)$. The $\beta(x)$ are derived from Balmer line intensities from 1962 eclipse data (Dunn et al. 1968, see appendix 3). The gradient at $x = 4,000$ km is added from the $H\epsilon$ eclipse curves derived from 1958 eclipse data (see figure 4c, ref. appendix 1). The result is shown in figure 1c and d. The discrepancy between observed and calculated distributions is seen in figure 1c below around 2,000 km and is probably due to an improper assumption. Hiei's $n_e n_p$ distribution is adopted below 2,000 km.

The n_H and $n_e n_p$ distributions in figure 1 give the ratio

$$n_e n_p / n_H = n_e x_H \quad (8)$$

where x_H is the ionization degree of hydrogen. Since $x_H \leq 1$, the ratio (8) gives the minimum of n_e . At $h = 5,000$ km, $\log n_H = 8.83$ and $\log(n_e n_p / n_H) = 10.73$. This means the minimum n_e exceeds n_H by a hundred times even although it is the main source of the electrons. A spicule structure can solve this contradiction. If the filling factor of the spicule has spherical symmetry, the derived values in figure 1 are taken as the ones of the spicule reduced by the filling factor. This solution is probable because a metallic line of Ca II H and a hydrogen line of $H\epsilon$ in the 1958 eclipse data are well correlated (see appendix 2) and they may be emitted from the same volume.

By assuming a spicule structure, equation (8) is rewritten as

$$n_e n_p / n_H = n_{es} x_H \quad (9)$$

where n_{es} is the electron density of the spicule and n_H , $n_e n_p$ are taken as reduced or smoothed densities. In order to obtain the parameters of the spicule structure the following two assumptions are made:

1) $x_H = 1$ above $h = 4,200$ km. At this height $\log n_{es}$ becomes ~ 10.73 and equal to the hydrogen density, n_{HS} , of the spicule, since the main electron source is hydrogen. n_{HS} thus obtained is close to a recent theoretical value, $\log n_{HS} = 10.6$ at $h = 4,000$ km, obtained by Kudoh and Shibata (1999). An higher n_{HS} with a lower x_H leads to the denser atmosphere.

2) $n_{es} = n_e$ below $h = 2,000$ km. This is equivalent to $f = 1$. However, the separate determination of n_H from that of $n_e n_p$ causes a slight inconsistency around $h = 2,000$ km as shown in table 1 and figure 3. This assumption is reasonable since the hydrostatic chromosphere can extend to around $h = 2,000$ km.

With the above two assumptions the spicule electron density is drawn by hand between $h = 2,000$ km and 4,200 km. Once n_{es} is fixed, x_H is obtained from equation (9). On the other hand n_{HS} is calculated for $h > 2,000$ km by the following formula,

$$\log n_{HS} = \log n_{es} - \log x_H \quad (10)$$

under the assumption that the main source of the electrons is hydrogen. Comparison of n_{HS} with n_H gives the filling factor,

$$\log f = \log n_H - \log n_{HS} \quad (11)$$

The chromospheric model with the spicules is given in figure 3 and table 1.

3. Ca II H and K, and $H\epsilon$ Line Profiles during the 1958 Total Solar Eclipse

Unlike other slitless spectrum, the flash spectrum obtained during the 1958 eclipse (Suemoto and Hiei 1959, 1962)

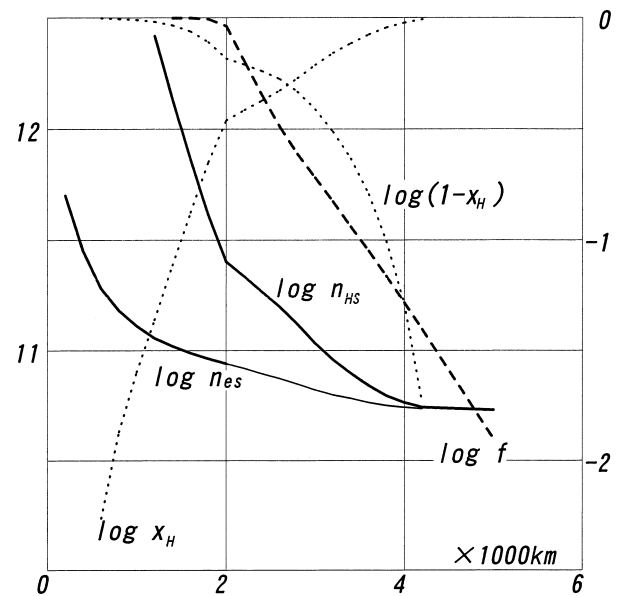


Fig. 3. Chromospheric model with spicules. n_{HS} : hydrogen density and n_{es} : electron density of the spicule, respectively, read by the left scale. x_H : ionization degree of hydrogen and f : filling factor of the spicule, read by the right scale. Thin part of the n_{es} curve is drawn by hand (see the text).

Table 1. Chromosphere Model.

height x1000km	Log \mathcal{N}_H	Log $\mathcal{N}_e \mathcal{N}_p$	Log($\mathcal{N}_e \mathcal{N}_p$ / \mathcal{N}_H)	Log \mathcal{N}_{es}	Log χ_H	Log(1- χ_H)	Log \mathcal{N}_{Hs}	Log f
5.0	8.829	19.555	10.726				10.73	-1.901
4.8	8.946	19.676	10.730					
4.6	9.069	19.800	10.731					
4.4	9.199	19.929	10730					
4.2	9.335	20.060	10.725	10.733	-0.008	-1.739	10.741	-1.406
4.0	9.479	20.196	10.717	10.740	-0.023	-1.288	10.763	-1.284
3.8	9.630	20.334	10.704	10.749	-0.045	-1.007	10.794	-1.164
3.6	9.789	20.476	10.687	10.763	-0.076	-0.794	10.839	-1.050
3.4	9.955	20.621	10.666	10.781	-0.115	-0.633	10.896	-0.941
3.2	10.130	20.769	10.639	10.798	-0.159	-0.513	10.957	-0.827
3.0	10.313	20.920	10.607	10.820	-0.213	-0.412	11.033	-0.720
2.8	10.505	21.079	10.574	10.846	-0.272	-0.332	11.118	-0.613
2.6	10.706	21.251	10.545	10.871	-0.326	-0.277	11.197	-0.491
2.4	10.916	21.437	10.521	10.893	-0.372	-0.240	11.265	-0.349
2.2	11.135	21.638	10.503	10.918	-0.415	-0.211	11.333	-0.198
2.0	11.364	21.842	10.478	10.940	-0.462	-0.184	11.402	-0.038
1.8	11.605	21.919	10.314	10.963	-0.649	-0.110	11.612	-0.007
1.6	11.860	21.971	10.111	10.987	-0.876	-0.062	11.863	-0.003
1.4	12.130	22.033	9.903	11.018	-1.115	-0.035	12.133	-0.003
1.2	12.423	22.116	9.693	11.055	-1.362	-0.019	12.417	0.006
1.0	12.740	22.252	9.512	11.110	-1.598	-0.011		
0.8	13.110	22.420	9.310	11.180	-1.870	-0.006		
0.6	13.575	22.592	9.017	11.280	-2.263	-0.002		
0.4	14.165	22.785	8.620	11.450	-2.830			
0.2	14.960	23.160	8.200	11.700	-3.500			
0.0	15.836	(23.290)	(7.454)					

is one of a few which can provide line profiles. The peak intensity, E_0 , total intensity, E , and $1/e$ -line width, $\Delta\lambda_{1/e}$, of the Ca II H and K, and $H\epsilon$ profiles have been measured (see appendix 1) in active and less active regions. They will be related to the chromospheric model in the previous section by the equations below.

Under the assumption of constant source function, S , the intensity at the projected height x and the wavelength distance from the line center $\Delta\lambda$ is

$$I(x, \Delta\lambda) = S[1 - \exp(-\tau(\Delta\lambda))] \quad (12)$$

τ is the optical thickness in the line of sight and given by

$$\tau(\Delta\lambda) = C_2 \lambda f_{abs} \int_{-\infty}^{\infty} \frac{n(h)}{V_D(h)} H(a, v) dy \quad (13)$$

where

$$C_2 = \sqrt{\pi} e^2 / (mc) = 1.4977 \times 10^{-2},$$

λ : wavelength of the line,

f_{abs} : absorption transition probability of the line,

$n(h)$: number density of the absorbing atom,

V_D : Doppler width in velocity scale,

$H(a, v)$: Voigt function with damping parameter a and $v = \Delta\lambda / V_D$.

The Doppler width taken from Suemoto's result (see figure 4, Suemoto 1963) is

$$V_D(\text{km/s}) = \begin{cases} 1.7 + 0.0051 h & h < 3,580 \text{ km} \\ 19.959 & h \geq 3,580 \text{ km} \end{cases} \quad (14)$$

The measured quantities are

$$E_0(x) = \int_x^{\infty} I(x, 0) dx \quad (15)$$

$$E(x) = \int_{-\infty}^{\infty} d(\Delta\lambda) \int_x^{\infty} I(x, \Delta\lambda) dx \quad (16)$$

and $\Delta\lambda_{1/e}$ is obtained from

$$E_0(x) = e \int_x^{\infty} I(x, \Delta\lambda_{1/e}) dx \quad (17)$$

The number density in the ground level of Ca II is

$$n(h) = A_{Ca} n_H(h) (1 - x_{CaII}) \quad (18)$$

where the calcium abundance relative to hydrogen, A_{Ca} , is -5.65 in logarithms, $n_H(h)$ is from table 1 and the ionization degree of the Ca II atom, x_{CaII} , is

$$x_{CaII} = \left[1 + n_{es} T_{CaII}^{-3/2} \exp\left(\frac{\chi_{CaII}}{k T_{CaII}}\right) / C_1 \right]^{-1} \quad (19)$$

from the Saha equation. The electron density in equation (19) is the spicule's value in table 1, T_{CaII} is the ionization temperature of Ca II and $\chi_{CaII} = 11.871$ eV is the ionization potential of Ca II. The fitting parameters to the observed Ca II H and K lines are the source function, S , the ionization temperature, T_{CaII} , and the damping parameter, a , in equations (12)–(19). The other constant parameters are $f_{abs}(\text{Ca II K}) = 0.682$, $f_{abs}(\text{Ca II H}) = 0.331$, $\lambda(\text{Ca II K}) = 3,934$ Å, and $\lambda(\text{Ca II H}) = 3,968$ Å. The result is shown in figures 4a and b. S is equivalent to the excitation temperature, T_{ex} , of 4,300 K and $T_{CaII} = 6,000$ K in the less active region. In the active region only an increase of the excitation temperature to 4,690 K seems to be enough. Here, ionization of Ca II is necessary, otherwise it leads to a stronger intensity at greater heights and a broader width at lower heights. $a = 0.001$ can explain the intensity and width increases near $x = 0$.

The number density in the second level of hydrogen is

$$n(h) = 4[n_e n_p](h) T_{BaC}^{-3/2} \exp\left[\frac{\chi_{ion} - \chi_2}{k T_{BaC}}\right] / C_1 \quad (20)$$

(see equation (7)), where $[n_e n_p](h)$ is given by table 1 and

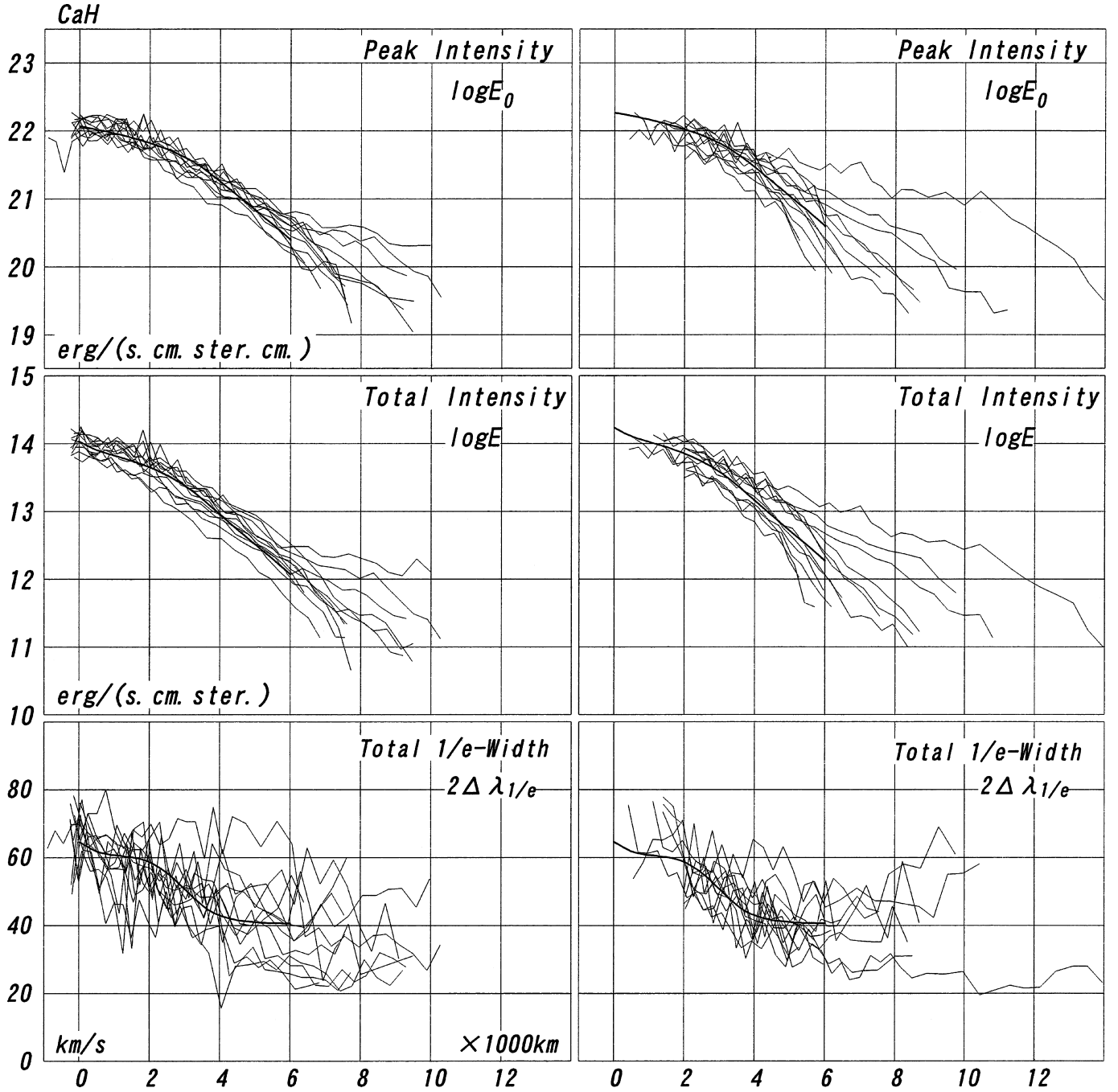


Fig. 4a. Fitting of the calculated curves (thick) to the observed Ca II H eclipse curves (thin).

T_{BaC} is the ionization temperature from the second level. The Doppler width should take into account the thermal broadening of hydrogen and equation (14) thus changes to

$$V_D = \sqrt{\frac{2kT_{\text{BaC}}}{m_H} + [\text{equation}(14)]^2} \quad (21)$$

where m_H is the hydrogen mass and the kinetic temperature is assumed to be T_{BaC} . The other constant parameters are $f_{\text{abs}} = 0.0127$, $\lambda(H\epsilon) = 3,970 \text{ \AA}$, $\chi_{\text{ion}} = 13.595 \text{ eV}$, $\chi_2 = 10.20 \text{ eV}$. Fitting to the observation of $H\epsilon$ is made with the use of equations (12)–(17), (20) and (21). Figure 4c is the result with $T_{\text{ex}} = 4,520 \text{ K}$, $T_{\text{BaC}} = 5,200 \text{ K}$ and $a = 0$ for the less active region. For the active region only an increase of the excitation temperature to $5,110 \text{ K}$ is sufficient.

4. Balmer and Sr II Emissions during the 1962 Eclipse

The chromospheric model in section 2 is applied to the total intensity of Balmer lines and Ca II-like Sr II resonance lines observed during the 1962 eclipse (Dunn et al. 1968).

$H\alpha$, $H\beta$, $H\gamma$, $H\epsilon$, H10, H15, and H20 are calculated with the use of the scheme for $H\epsilon$ in section 3 but changing the wavelengths and transition probabilities (ref. Allen 1973). The result is shown in figure 5, with predictions of the peak intensities and 1/e widths. The fitting parameters are $a = 0$, $T_{\text{BaC}} = 5,200 \text{ K}$, the same as for the 1958 eclipse, and excitation temperatures of $5,020 \text{ K}$, $4,860 \text{ K}$, $4,740 \text{ K}$, $4,720 \text{ K}$, $4,800 \text{ K}$, $4,840 \text{ K}$, $4,790 \text{ K}$ for the above Balmer lines, respectively. Let b_j be NLTE factors and T_e the electron temperature, the relations

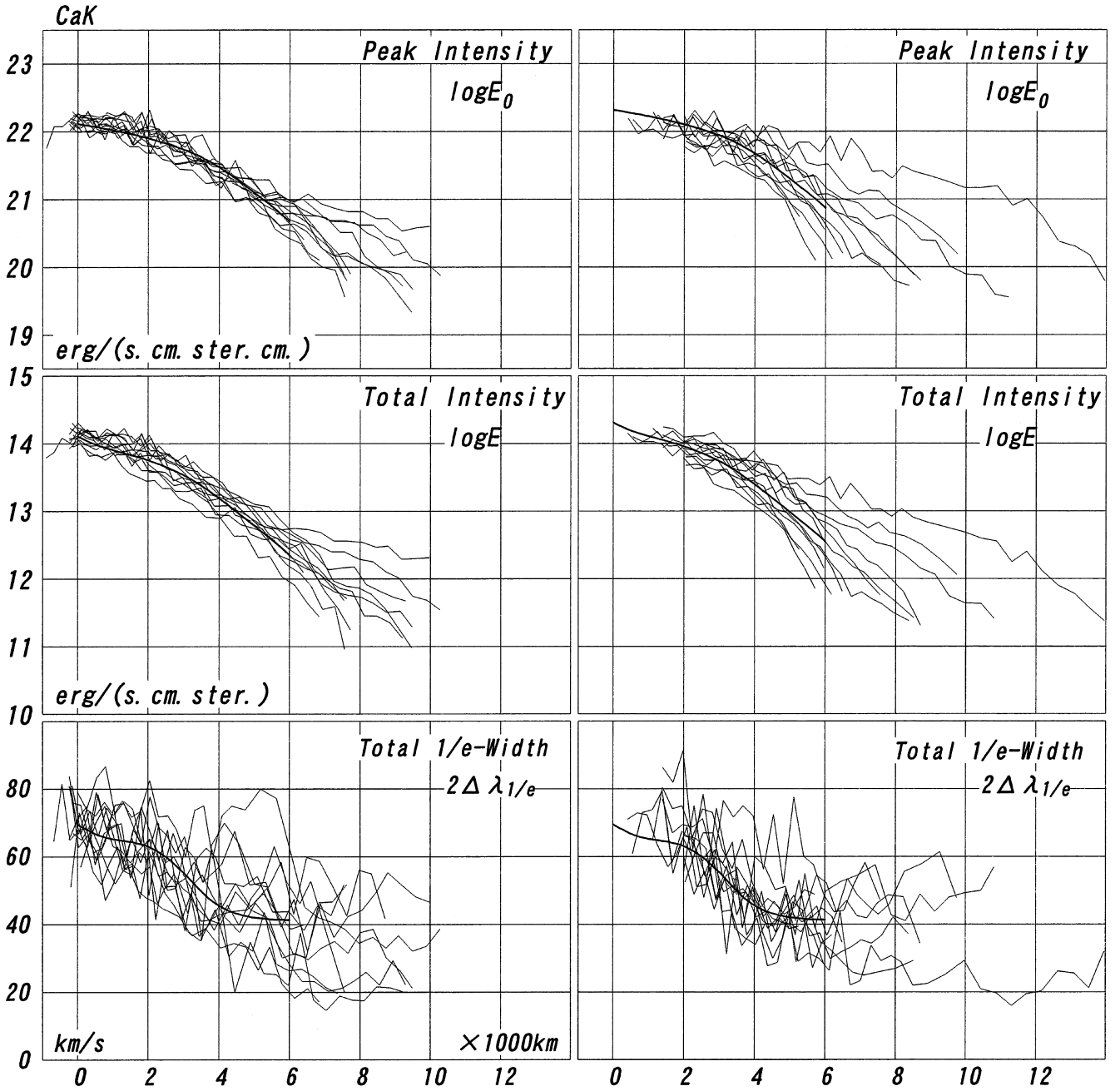


Fig. 4b. Fitting of the calculated curves (thick) to the observed Ca II K eclipse curves (thin).

$$T_{\text{BaC}}^{-3/2} \exp\left(\frac{\chi_{\text{ion}} - \chi_2}{kT_{\text{BaC}}}\right) = b_2 T_e^{-3/2} \exp\left(\frac{\chi_{\text{ion}} - \chi_2}{kT_e}\right) \quad (22)$$

from the Saha equation and

$$\exp\left(\frac{\chi_2 - \chi_j}{kT_{\text{ex},j}}\right) = \frac{b_j}{b_2} \exp\left(\frac{\chi_2 - \chi_j}{kT_e}\right) \quad (23)$$

from the Boltzmann formula are obtained (ref. Thomas and Athay 1961). $T_{\text{ex},j}$ is the excitation temperature of the j -th Balmer line. If $b_{20} = 1$, the combination of the above relations gives $T_e = 7,900$ K, $b_2 = 24.8$, $b_3 = 5.1$, $b_4 = 2.4$, $b_5 = 1.5$, $b_7 = b_{10} = b_{15} = 1.1$. These values can be compared with those of the VAL-C model (see tables 12 and 17, Vernazza et al. 1981). The electron temperature is comparable with Matsuno and Hirayama's result derived from Balmer and metallic line widths obtained from 1966 eclipse data (Matsuno and Hirayama 1988).

Sr II 4078 and 4216 intensities are calculated with the use of the scheme for Ca II in section 3 but changing the wavelengths, transition probabilities, abundance ratio, and ionization potential. The constant parameters used are f_{abs} (Sr II 4078)=0.708, f_{abs} (Sr II 4216)=0.339, $\chi_{\text{Sr II}} = 11.03$ eV, $\log A_{\text{Sr}} = -9.1$. The result is shown in figure 6, with predictions of the peak intensities and 1/e widths. The fitting parameters are the ionization temperature of 5,140 K, the excitation temperature of 4,500 K and $a = 0$.

The Ca II H and K lines observed during the 1962 eclipse can also be explained by the chromospheric model with a slightly higher excitation temperature (see appendix 1).

5. Ca II H and K Spicules observed during the 1958 Eclipse

The total and peak intensities of the Ca II H and K

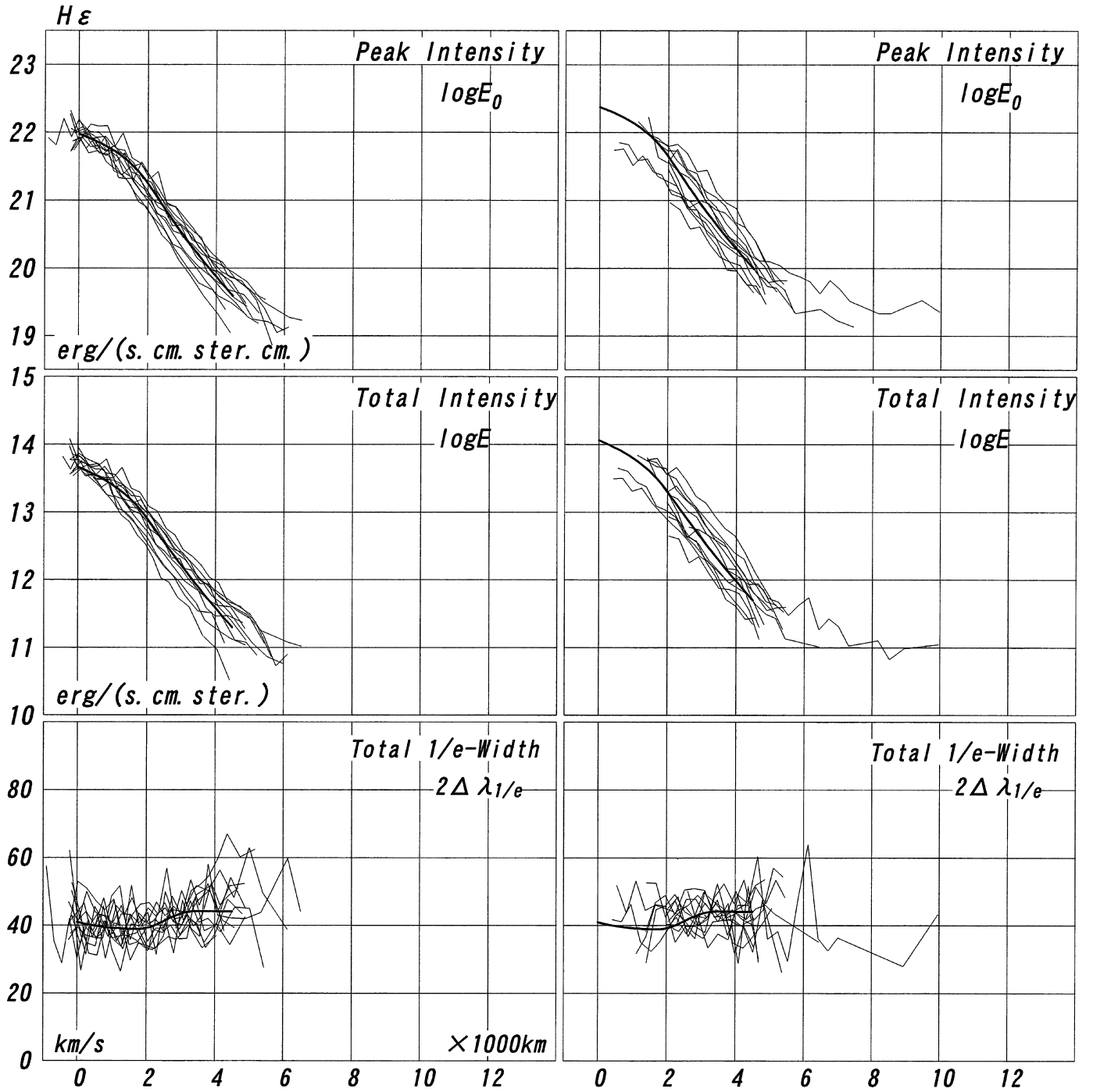


Fig. 4c. Fitting of the calculated curves (thick) to the observed $H\epsilon$ eclips curves (thin).

spicules have nearly the same emission gradients in the high chromosphere. Figure 7 shows the intensities overlapped by a shift on the Ca II K total intensity. Special attention is paid to the top height regions where the curves are linear. The spicules there are assumed not to be overlapped but isolated.

According to equation (13), in the thin atmosphere, the shifts give the intensity ratio of (Ca II K)/(Ca II H) as equal to 2 and the total to peak intensity ratio as equal to $\sqrt{\pi}V_D$, respectively. The emission gradients, intensity ratios and Doppler widths of the 24 spicules are listed in table 2. The averaged intensity ratio, (Ca II K)/(Ca II H), is less than 2 and might suggest thick spicules. However, Hirayama's examination (private communication, ref. Hirayama 1964) found that the color correction increased the Ca II K intensity relative to the Ca II H

intensity by a small amount of 0.05 in logarithms and therefore the spicules are concluded to be thin.

Correlations between the top height, the total intensity at the top height, the emission gradient, and the Doppler width are shown in figure 8. Taller spicules have smaller emission gradients (a), spicules with brighter intensity at the top might be shorter (b) and have stronger turbulence (e). The other correlations are not remarkable. A difference between the active spicules and the less active ones is not obvious. However, their average total intensities at the top height (see table 2) are different by 0.2 in logarithms, which is equivalent to the excitation temperature difference obtained in section 3. The derived Doppler widths are compared with the measured 1/e-widths in figure 9. The measured ones have a tendency to be narrower

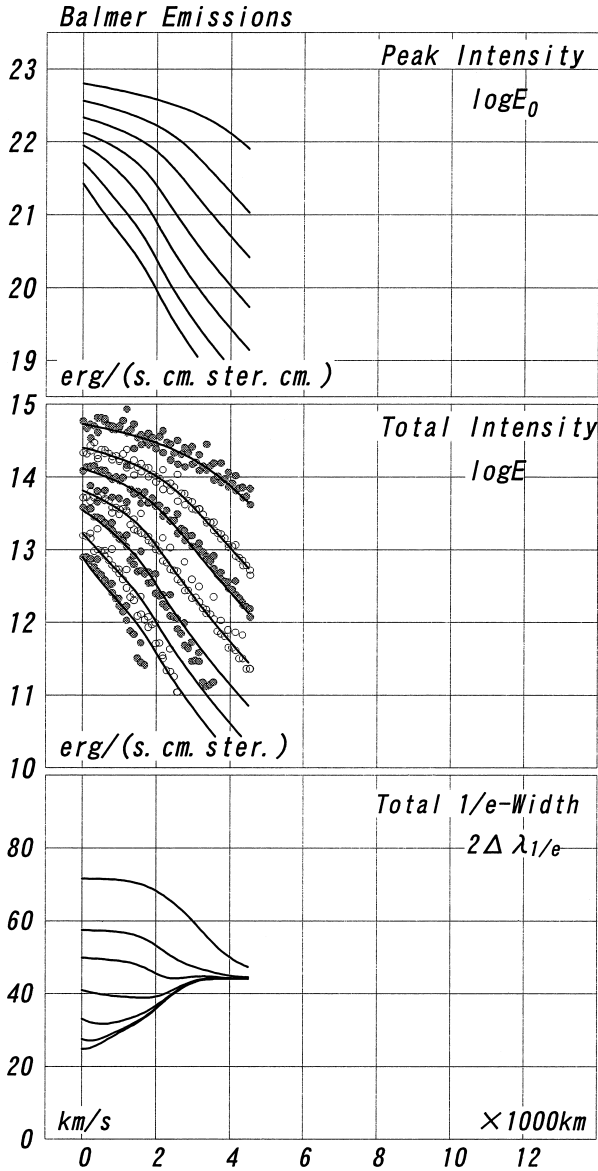


Fig. 5. Fitting of the calculated curves (solid) to the observed total intensities of the Balmer lines (circles). Each frame shows, from top to bottom, curves or circles for $H\alpha$, $H\beta$, $H\gamma$, $H\epsilon$, $H10$, $H15$, and $H20$.

than the calculated ones. This may suggest that the line profile has a narrower core and broader wing than the Doppler profile does.

The total intensity of the spicules can give the intensity at the projection height x by the following formula (ref. equations (3) and (5)),

$$I(x) = \beta(x)E(x) \quad (24)$$

Since the spicules are thin,

$$I(x) = C_3 N_2 A_{21} / \lambda \quad (25)$$

where

$C_3 = hc/(4\pi) = 1.581 \times 10^{-17}$, $A_{21} = 1.5 \times 10^8$ is the spontaneous emission coefficient from the second level, and N_2 is the number in the line of sight of the second level of the Ca II atom. The number in the line of sight of the ground state is

$$N_1 = \frac{N_2}{2} \exp\left(\frac{\chi_2}{kT_{ex}}\right) \quad (26)$$

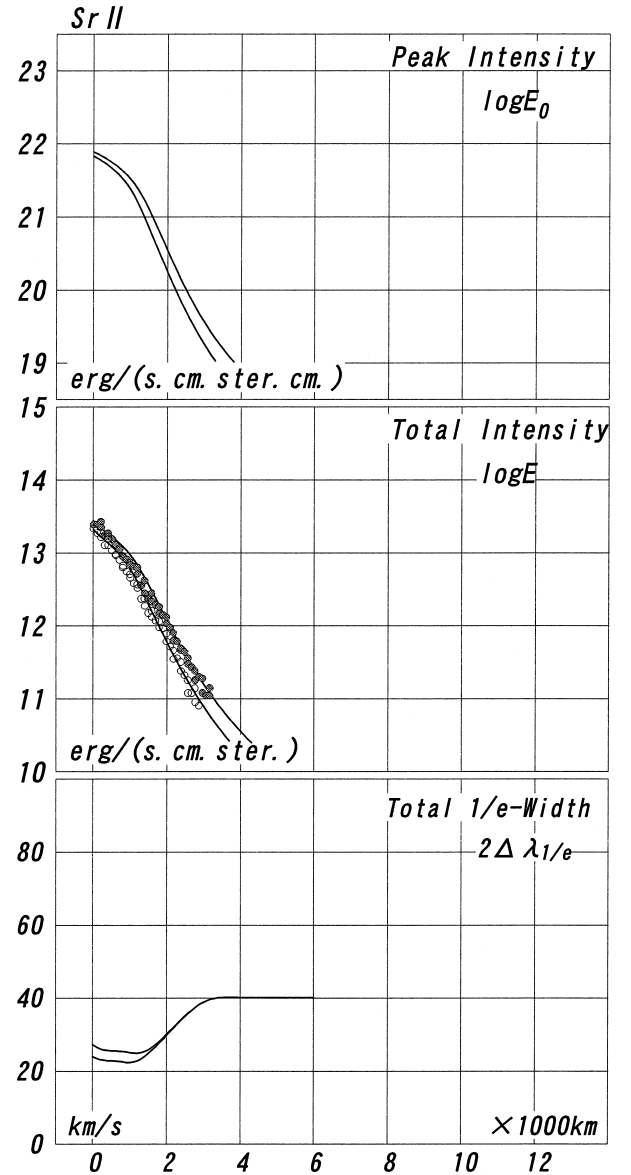


Fig. 6. Fitting of the calculated curves (solid) to the observed total intensities of Sr II 4078 and 4216 (circles). Each frame shows a curve or circles for Sr II 4078 upwards.

by the Boltzmann formula. The total number in the line of sight of hydrogen is calculated from equations (18) and (19), as

$$N_H(x) = \frac{N_1}{A_{Ca}} \left[1 + C_1 T_{CaII}^{3/2} \exp\left(-\frac{\chi_{CaII}}{kT_{CaII}}\right) / n_{es} \right] \quad (27)$$

If the temperatures of the active and less active region obtained earlier are used, the total number of hydrogen can be calculated for reasonable electron densities. Furthermore, the geometrical thickness of the spicules will be estimated with $n_H \sim n_{es}$. In table 2 the thicknesses at the top are listed assuming $\log n_H = 10$ which is a little denser than the probable density of the corona. They are comparable with the “diameters” obtained by Nishikawa (1988). 14 spicules are thinner than 1,000 km and 19 thinner than 2,000 km. With an assumption of constant thickness along the height (Lynch et al. 1973; Nishikawa, private communication), equation (27) gives the electron density at the bottom of the linear sections in figure 7. The bottom densities plotted in figure 10 are on the av-

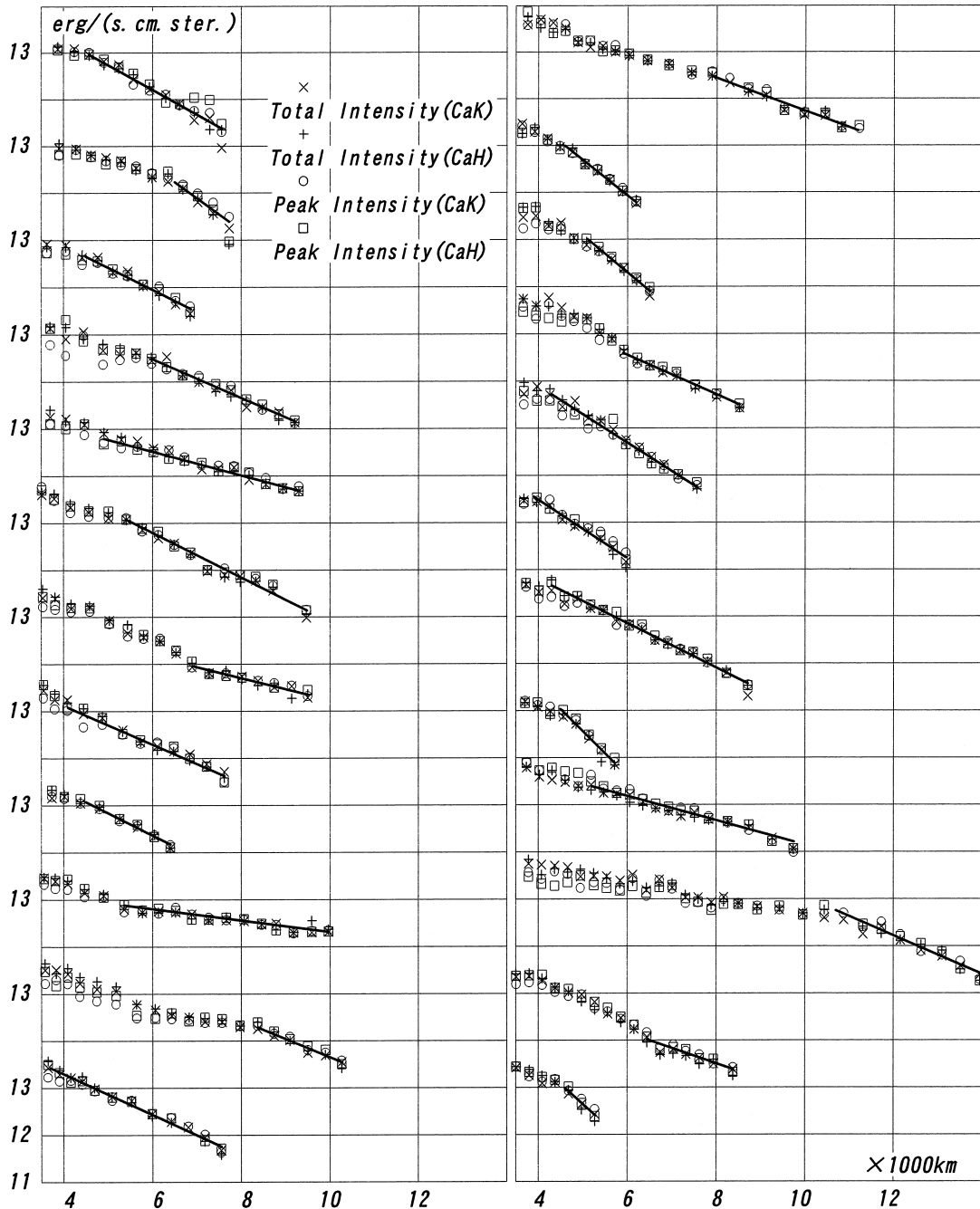


Fig. 7. Overlapped total and peak intensity eclipse curves of Ca II H and K spicules. Ordinates give the scale of the Ca II K total intensity. The solid lines give the emission gradients near the top of the spicules.

erage connected to the electron density of the mean spicule model in figure 3. This supports the assumption of the average top height density, $\log n_{es} = 10$. The large thickness of the short spicule no. 24 reduces to 870km if a revised top density, $\log n_{es} = 10.5$, is adopted. This still gives a reasonable bottom density of $\log n_{es} = 10.74$.

The electron density obtained from the red continuum of the 1970 eclipse data (Makita 1972) is plotted in figure 10. This assumed a thickness of 1000 km. If this increases to 2,500 km, the plot moves on to the mean model.

6. Results and Discussion

1) Chromospheric structures

The spicules start around 2,000 km and have a filling factor of 0.05 at 4,000 km. Their starting height might be low-

ered (e.g., Suemoto and Hiei 1962; Kanno et al. 1971) if the spicule density near the top is increased to $\log n_{es} = 11$. The spicule model presented by Beckers (1968) has in the pertinent heights considerably high electron densities and fairly small filling factors.

The Ca II H and K profiles observed at lower projection heights than 2,500 km during the 1958 eclipse are not always symmetric and are classified into 3 groups: 28 percent with double equal intensity peaks, 14 percent with two unequal intensity peaks, and the rest with a single peak. This indicates that they are mainly shaped not by self absorption but by macroscopic motion. Suemoto (1963) reports a line shift of 19 km/s.

The optical thicknesses of Ca II K and H ϵ become thin at higher projection heights than 4,500 km and 2,000 km, respec-

Table 2. Ca II H and K Spicule Data.

No.	height range (km)	grad- ient*	Log E_{top}	---- K/H ¥	Line K0/H0 ¥	Ratio# K/K0 -8.5	---- H/H0 -8.5	Line Width (km/s)	Log t_{top}	Log $N2_{top}$	Log $N1_{top}$	Thick- ness (km)	Log $N1_{bot.}$	Log $n1_{bot.}$	Log $nebot.$
bottom-top															
1	4580-7560	1.227	11.37	0.2	0.2	0.25	0.25	24.1	3.46	7.68	11.07	684	12.66	4.82	10.88
2	6500-7730	1.577	11.39	0.25	0.2	0.1	0.05	16.1	3.59	7.81	11.20	920	12.04	4.08	10.44
3	4500-6860	1.089	11.54	0.25	0.3	0.0	-0.5	12.8	3.58	7.80	11.19	897	12.31	4.35	10.60
4	5980-9220	0.951	11.15	0.25	0.4	-1	0.05	13.0	3.13	7.35	10.74	319	12.08	4.57	10.72
5	4920-9300	0.583	11.68	0.25	0.25	0.05	0.05	15.2	3.45	7.57	10.96	527	12.07	4.35	10.59
6	5400-9480	1.094	11.14	0.275	0.25	0.325	0.325	28.7	3.18	7.40	10.79	359	12.73	5.17	11.09
7	6900-9500	0.537	11.35	0.3	0.3	0.15	0.15	19.1	3.08	7.30	10.69	286	11.30	3.84	10.33
8	4100-7620	0.952	11.62	0.25	0.35	0.2	0.3	24.3	3.60	7.82	11.21	942	12.67	4.69	10.80
9	4500-6400	1.096	12.17	0.3	0.3	0.25	0.25	24.1	4.21	8.43	11.82	3855	12.73	4.14	10.49
10	5360-9970	0.284	12.31	0.25	0.25	0.25	0.25	24.1	3.76	7.98	11.38	1377	11.94	3.81	10.30
11	8370-10280	0.897	11.54	0.3	0.2	0.2	0.15	20.3	3.49	7.71	11.11	740	11.85	3.98	10.39
12	3700-7550	0.998	11.76	0.25	0.25	0.25	0.25	24.1	3.76	7.98	11.37	1364	13.04	4.91	10.94
13	7900-11180	0.813	11.31	0.25	0.3	0.3	0.25	25.6	3.22	7.44	10.64	252	11.80	4.39	10.62
14	4600-6180	1.742	11.79	0.25	0.2	0.2	0.15	20.3	4.03	8.25	11.45	1629	12.64	4.43	10.64
15	5100-6500	1.820	11.90	0.2	0.25	0.2	0.25	22.8	4.16	8.38	11.58	2193	12.68	4.34	10.60
16	5900-8520	0.968	11.49	0.25	0.3	0.05	0.1	16.1	3.48	7.70	10.89	454	11.99	4.34	10.60
17	4300-7560	1.396	11.74	0.25	0.25	0.25	0.25	24.1	3.89	8.10	11.30	1164	13.28	5.21	11.14
18	3900-5970	1.442	12.25	0.2	0.15	0.3	0.25	25.6	4.41	8.63	11.83	3890	13.12	4.53	10.70
19	4300-8710	1.089	11.57	0.3	0.3	0.25	0.25	24.1	3.61	7.83	11.02	614	13.11	5.32	11.23
20	4500-5720	2.163	11.88	0.25	0.25	0.3	0.3	27.1	4.22	8.43	11.63	2489	12.78	4.38	10.62
21	5200-9750	0.598	12.21	0.25	0.25	0.3	0.35	28.7	3.99	8.21	11.40	1472	12.59	4.42	10.64
22	10700-13950	0.944	11.42	0.25	0.25	0.0	0.0	13.6	3.40	7.61	10.81	377	12.14	4.57	10.72
23	6440-8380	0.752	11.38	0.25	0.3	0.2	0.25	22.8	3.26	7.48	10.67	276	11.31	3.87	10.33
24	4600-5280	1.991	12.40	0.25	0.2	0.3	0.25	25.6	4.70	8.92	12.12	7586	12.70	3.82	10.30
mean		1.125		0.253	0.260	0.191	0.195	21.8							

* multiply by 10^{-8} (cm)
in logarithms
¥ should be increased by 0.05 (see the text)

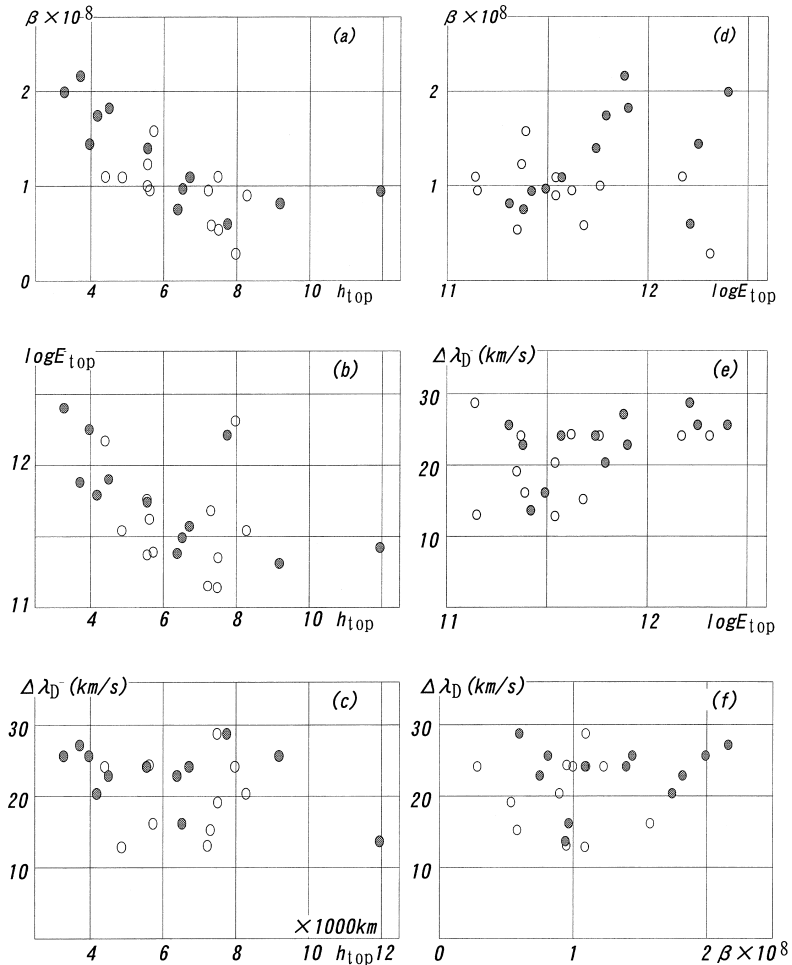


Fig. 8. Correlations of spicule parameters. The open and gray circles are from the less active and active regions, respectively. β : emission gradient, h_{top} : top height, E_{top} : total intensity at the top, and $\Delta\lambda_D$: Doppler width.

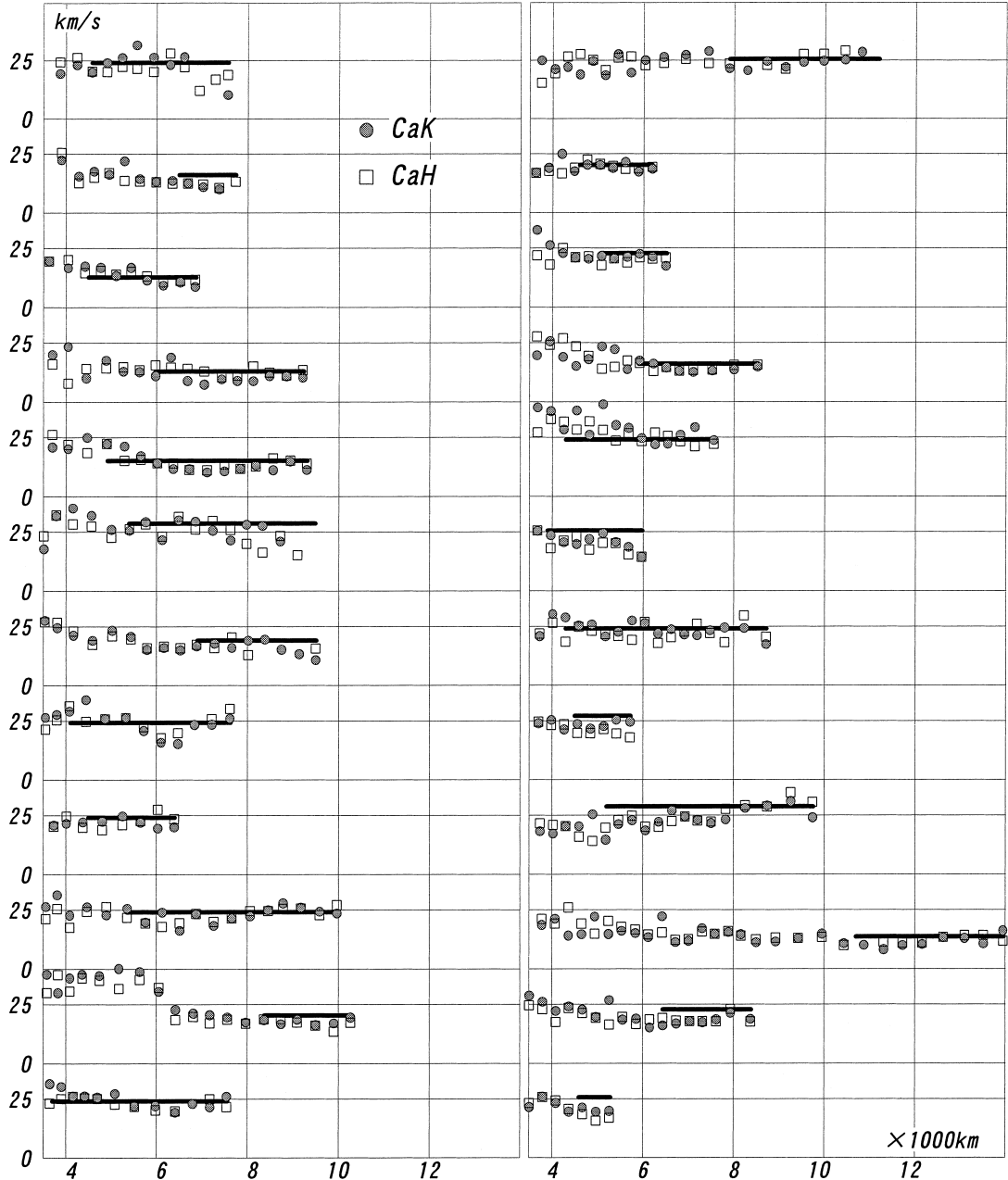


Fig. 9. Measured $1/e$ widths (symbols) and Doppler widths from the peak to total intensity ratio (solid lines).

tively, as shown in figure 11.

The radial or vertical optical thicknesses of Ca II K and H α lines can be estimated from the spicule densities given in table 1. The dotted curves in figure 11 suggest that the Ca II K and H α filtergrams see the chromosphere of $h = 4,000$ – $5,000$ km.

2) Excitation and ionization.

Figure 12 shows a summary of the source functions obtained from this study, the brightness temperature curves, and the photospheric intensity from Allen (1973). The H α emission and the Balmer continuum roughly correspond to half the photospheric intensity. The other emissions are weaker than this. The Ca II H and K, and H ϵ emissions from the active region are mainly enhanced by the increase of their source functions (gray circles, see section 3).

The ionization temperatures obtained are 6,000K for

Ca II and 5,140 K for Sr II. They correspond to UV radiation of 5,450 K and 5,100 K, respectively (see figure 3, Vernazza et al. 1981). Ca II is half doubly ionized at $h = 1,100$ km, Sr II remains singly ionized below $h = 4,000$ km, in contrast with the half-ionized hydrogen at $h = 2,700$ km (see figure 3).

The excitation and ionization obtained should be explained by more detailed analysis.

3) Individual spicules

The individual spicules have a Doppler width of 22 km/s on the average and a typical top density of $\log n_{es} = 10$, which are consistent with theoretical predictions (e.g., Kudoh and Shibata, 1999). Their thicknesses are comparable with Nishikawa's (1988) and most of them are thinner than 2,000 km. The spicules of the active and less active region are similar in their physical parameters except that the active spicules have a higher excitation temperature (see section 5).

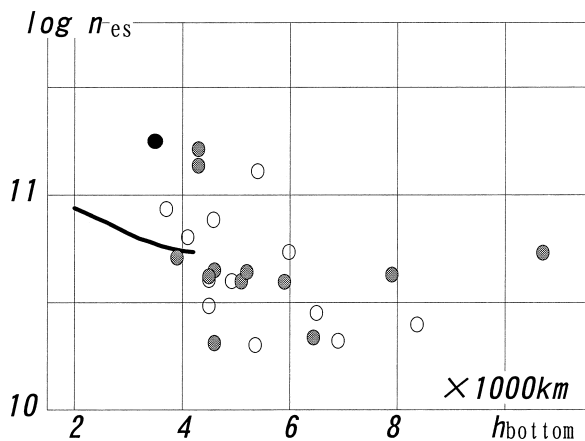


Fig. 10. Electron density of the spicules at their bottom heights (see the text). The open and gray circles are from the less active and active regions, respectively. The filled circle is from the 1970 eclipse (Makita 1972). The solid curve shows the average model in figure 3.

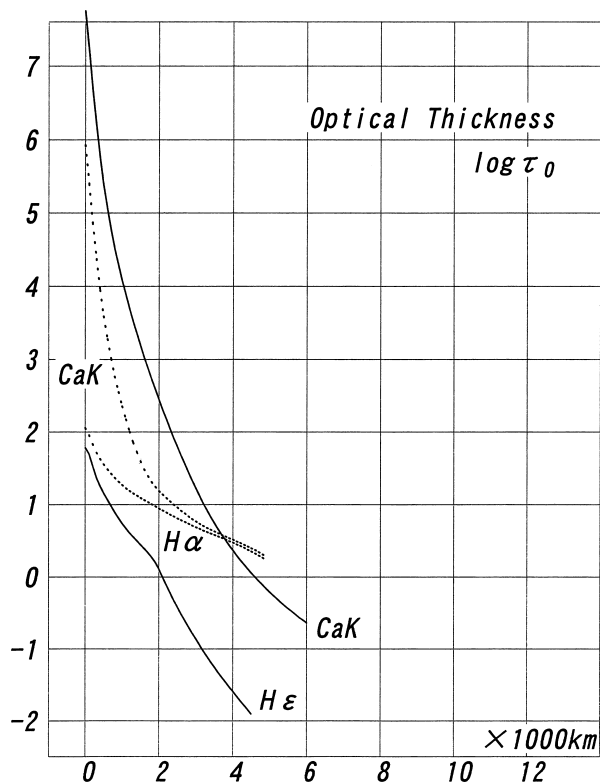


Fig. 11. Optical thickness of Ca II K, H ϵ , and H α emission. The solid curves are for the flash spectrum. The dotted curves are radial or vertical optical thicknesses estimated from the spicule densities given in table 1.

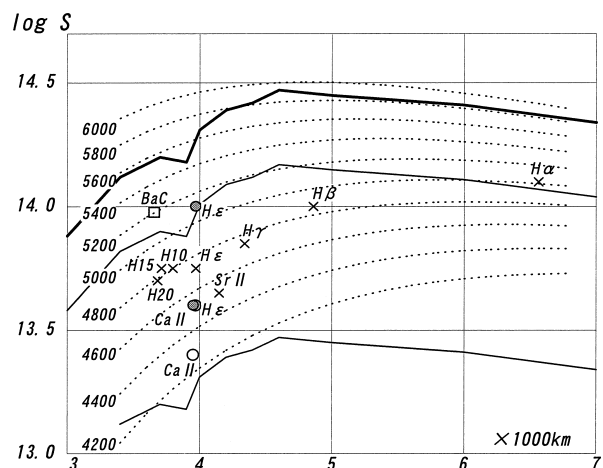


Fig. 12. Source functions of the emission lines. The open and gray circles are from the 1958 eclipse and for the less active and active regions, respectively. The crosses are from the 1962 eclipse. The square corresponds to the ionization temperature of the second level of hydrogen. The thick curve shows the photospheric brightness (Allen 1973). The thin curves are half and one tenth of the photospheric brightness. The dotted curves are brightnesses with the equivalent temperatures shown at their left ends.

Appendix 1. Ca II H and K, and H ϵ Emissions during the 1958 Total Solar Eclipse

The slitless flash spectrum of the 1958 total solar eclipse is rare in that it can provide line profiles (Suemoto and Hiei 1959, 1962). The peak intensity, total intensity, and $1/e$ -width of Ca II H and K, and H ϵ profiles have been measured with a scanning slot equivalent to 5 km/s (wavelength resolution) \times 0.65'' (spatial resolution). 24 spicules were selected for the measurement and half of them (Nos. 13–24) were below an active corona. The identity of the spicules was lost below 4,000–5,000 km and the measured positions there relied on the scale of the microphotometer. The sun's visible limb, the base of the chromosphere, was determined from eclipse curves of the continuum. For the active region the determination was made by eye instead of scanning. This seems to produce a little more scatter in analyses of the active region data. The absolute calibration of the intensity was made by Hiei (1963). Hirayama later found a color sensitivity difference between Ca II K and Ca II H wavelength regions (private communication, ref. Hirayama 1964). The Ca II K intensity should be increased by 0.05 in logarithms relative to the Ca II H and H ϵ intensities. The listed intensities in the following tables do not take into account this correction. Table A1 gives the projected heights of the moon's limb at the measured positions, table A2a–c gives the Ca II H data, table A3a–c gives the Ca II K data, and table A4a–c gives the H ϵ data. A graphic presentation of these data was made earlier (Makita 2000). Figure A1 shows comparison of the total intensities obtained by other observers (Cillie and Menzel 1935; Houtgast 1957; Vjaznityn 1956; Dunn et al. 1968) with our average eclipse curves calculated in the text for the less-active region. The eclipse curves for the active region may pass through the data points of the other observers better.

The author expresses his hearty thanks to Professors Ei-jiro Hiei and Tadashi Hirayama for their encouragements. Thanks are also due to Professor Tadashi Hirayama for critical and valuable comments and to Dr. David Brooks for checking the manuscript.

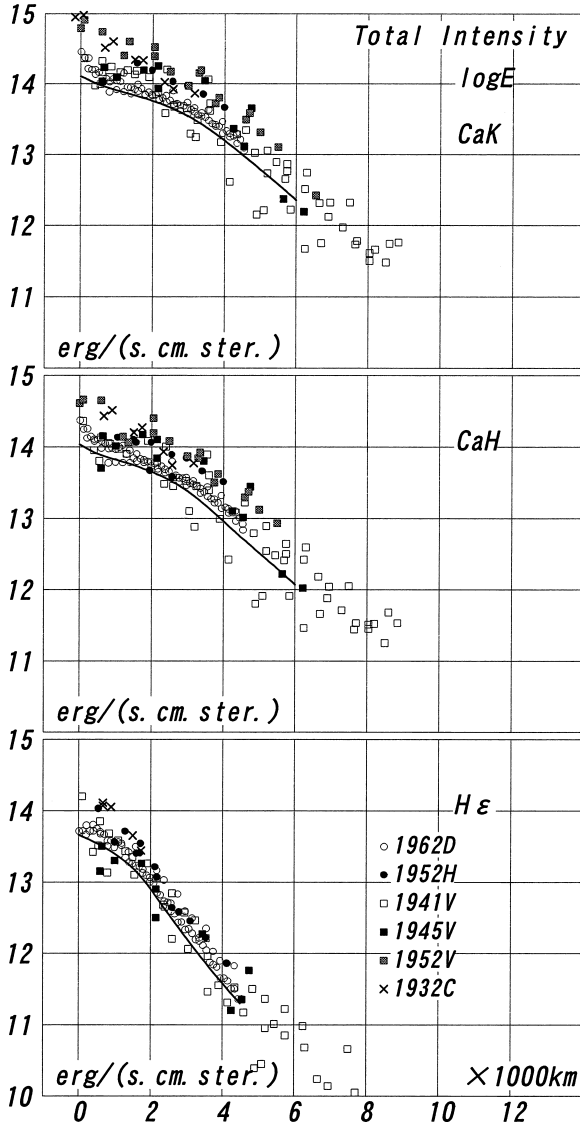


Fig. A1. Total intensities of Ca II H and K, and H ϵ . 1962D: Dunn et al. (1968), 1952H: Houtgast (1957), 1941V, 1945V and 1952V: Vjazanityn (1956), and 1932C: Cillie and Menzel (1935) are compared with the solid curves calculated in the text.

Appendix 2. Correlation between Ca II H and H ϵ Emissions

The Ca II H and H ϵ emissions tabulated in appendix 1 are combined in figure A2. They are close neighbors in the spectrum and measured by one and the same scanning. If different volumes of the chromosphere contribute to the emission, they may show separate behaviors. A good correlation of the two emissions in figure A2 is not contrary to the view that both the emissions are from the same volume. The slight dispersion seen in the active region diagram is due to the uncertainty of the height determination (see appendix 1).

Appendix 3. Emission Gradients of the Balmer Lines

Equation (7) in the text shows that the emission gradient of the Balmer lines is described by gradients of $n_e n_p$ and temperature. The latter gradient will be far smaller than the former in the high chromosphere where the density is low and the radiation is dominant. If this is the case, the emission gradient of the Balmer lines are related only to the gradient of $n_e n_p$. To increase the accuracy,

$$E_j / \left[C_3 (A_j / \lambda_j) C_1 j^2 T^{3/2} \exp \left(\frac{\chi_j - \chi_{ion}}{kT} \right) \right]$$

are plotted against the projected height as in figure A3 and logarithmic gradients of $n_e n_p$ are obtained from the averaged curve. The Balmer lines, H4–H33, during the 1962 eclipse (Dunn et al. 1968), with $\log E_j < 13.56$, are used. The constant parameters, A_j , λ_j , χ_j , are taken from *Astrophysical Quantities* by Allen (1973). The temperature is taken as 5,000 K which is effective only for low Balmer lines. The derived gradients are in table A5.

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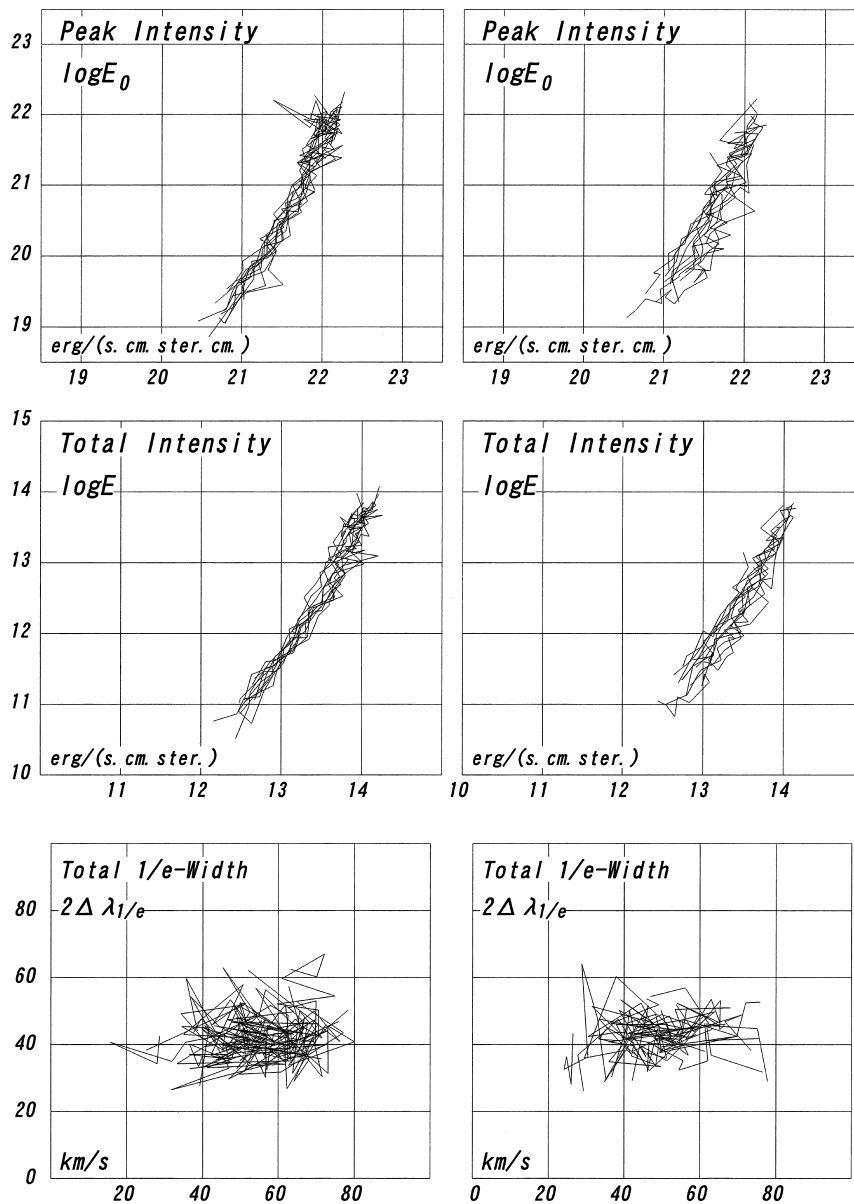


Fig. A2. Correlation between Ca II H (abscissa) and He ϵ (ordinate) emissions from the 1958 eclipse data.

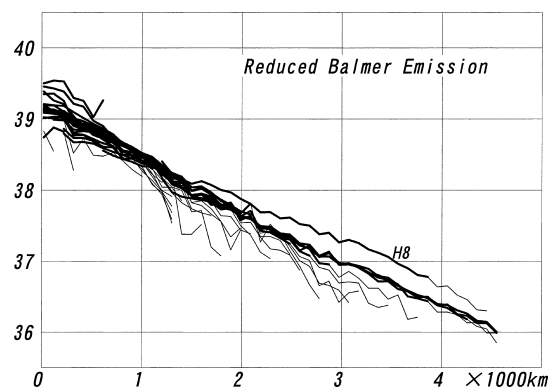


Fig. A3. Reduced Balmer line emissions to the integrated n_{enp} . H8 runs higher due to the blend with the He line. The part with the thin lines corresponds to the weaker intensity, $\log E < 11.8$, and they run lower, probably due to a calibration problem.

Table A2a. Total Intensity ($\log E - 10$) of Ca II H (erg/sec.cm.ser).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*
frame no.	71	70																							71
	1.144	0.661	0.877	0.928	1.422	1.465	0.791	0.969	2.115	1.123	2.302	1.409													70
69	1.150	1.289	1.192	1.597	1.358	1.266			2.050	1.494															69
68	1.519	1.630	1.349	1.764	1.484	1.243			2.103	1.700											1.007				68
67	1.664	1.826	1.137	1.432	1.855	1.471	1.378		2.234	1.890		1.132									1.250				67
66	1.827	2.159	1.432	1.541	1.843	1.569	1.563		2.301	2.095		1.477									1.651				66
65	2.066	2.070	1.577	1.753	1.974	1.736	1.517	1.336	2.376	2.014		1.466									1.774				65
64	2.327	2.234	1.782	1.882	2.067	2.061	1.726	1.569	2.322	2.141		1.543									1.882				64
63	2.456	2.403	2.012	2.072	2.154	2.271	1.950	1.703	2.341	2.176	1.335	1.780									2.010				63
62	2.535	2.456	2.102	2.271	2.310	2.471	2.191	1.896	2.482	2.204	1.621	1.880									2.160				62
61	2.736	2.556	2.324	2.332	2.376	2.632	2.305	1.925	2.089	2.489	2.300														61
60	2.819	2.671	2.433	2.468	2.574	2.789	2.540	2.101	2.294	2.466	2.384	2.018	2.234								1.797	2.514			60
59	2.930	2.805	2.601	2.552	2.679	2.948	2.632	2.319	2.425	2.629	2.483	2.155	2.307		1.185						2.008	2.440			59
58	3.111	2.899	2.614	2.663	2.890	3.042	2.943	2.641	2.668	2.867	2.853	2.446	2.468		1.449						2.312	2.568			58
57	3.171	3.120	2.859	2.905	2.906	3.137	2.982	2.810	2.784	2.969	2.965	2.569	2.553		1.578	1.453					2.375	2.547	0.997		57
56	3.210	3.102	2.998	2.926	3.154	3.350	3.134	2.936	2.938	3.151	3.055	2.757	2.666		1.867	1.752					2.417	2.654	1.342		56
55	3.361	3.336	3.080	3.016	3.208	3.376	3.302	3.127	3.028	3.209	3.244	2.993	2.769	1.767	1.725	1.965	1.919				2.472	2.685	1.242		55
54	3.442	3.358	3.134	3.043	3.244	3.463	3.355	3.310	3.136	3.230	3.134	2.975	2.808	1.985	1.923	2.069	2.019				2.521	2.623	1.417		54
53	3.546	3.283	3.282	3.192	3.427	3.600	3.398	3.298	3.175	3.368	3.346	3.135	2.944	2.220	2.130	2.192	2.266				2.620	2.744	1.464		53
52	3.628	3.419	3.372	3.232	3.502	3.594	3.653	3.357	3.224	3.440	3.330	3.326	2.984	2.358	2.347	2.412	2.416	1.831	2.195		2.671	2.832	1.437		52
51	3.769	3.536	3.499	3.337	3.612	3.699	3.639	3.482	3.358	3.766	3.502	3.343	3.228	2.687	2.563	2.665	2.630	2.111	2.415		2.728	3.081	1.702		51
50	3.735	3.553	3.482	3.510	3.742	3.794	3.724	3.594	3.502	3.603	3.525	3.341	3.175	2.736	2.751	2.843	2.916	2.424	2.528		2.795	3.020	1.981		50
49	3.770	3.632	3.607	3.520	3.659	3.755	3.625	3.668	3.519	3.636	3.687	3.396	3.267	2.877	2.797	3.085	3.023	2.596	2.702	1.591	2.942	2.987	2.130		49
48	3.867	3.801	3.724	3.579	3.837	4.207	4.040	3.788	3.799	4.026	3.990	3.666	3.507	3.131	2.980	3.167	3.155	2.746	2.830	1.657	3.012	3.123	2.339		48
47	3.917	3.742	3.644	3.530	3.750	3.806	3.870	3.670	3.650	3.792	3.671	3.579	3.340	3.001	3.146	3.163	3.350	2.938	2.882	2.217	3.060	3.069	2.396	2.027	47
46	3.814	3.776	3.737	3.615	3.874	3.934	3.990	3.724	3.784	3.998	3.805	3.712	3.493	3.295	3.500	3.330	3.582	3.095	3.131	2.461	3.117	3.186	2.571	2.280	46
45	3.910	3.821	3.825	3.582	3.848	3.850	4.093	3.885	3.930	3.917	3.806	3.581	3.538	3.343	3.481	3.338	3.547	3.227	3.142	2.675	3.270	3.296	2.793	2.694	45
44	4.024	3.881	3.731	3.785	3.997	4.093	4.008	3.993	3.881	4.010	4.011	3.967	3.766	3.526	3.615	3.493	3.718	3.308	3.477	2.639	3.475	3.375	2.878	2.869	44
43	3.857	3.825	3.801	3.748	3.841	3.965	4.048	3.769	3.908	3.992	4.045	3.936	3.696	3.364	3.550	3.546	3.678	3.497	3.345	2.851	3.391	3.318	3.064	3.002	43
42				3.829	3.838	3.989	4.026	3.955	3.856	4.082	4.000	4.043	3.804	3.729	3.776	3.810	3.910	3.767	3.415	2.918	3.560	3.391	3.172	3.128	42
41				3.948	4.255	4.200	4.175	3.969	3.887	3.987	3.986	3.939	3.661	3.720	3.756	3.652	3.805	3.615	3.115	3.433	3.265	3.160	3.190	41	
40				3.947	3.920	3.984	4.150	4.017	4.048	4.129	4.094	4.087	3.604	3.694	3.950	3.914	3.941	3.883	3.690	3.216	3.630	3.585	3.347	3.320	40
39							4.062	3.937	4.218	4.210	4.130	3.993	3.993	3.915	4.119	4.021	4.070	3.861	3.728	3.447	3.741	3.814	3.479	3.451	39
38											4.140	3.780	3.836	4.074	4.029	4.024	3.945	3.775	3.407	3.851	3.599	3.495	3.497	38	
37											4.082	3.947	3.992	4.044	3.955	4.100	3.713	3.858	3.560	3.886	3.689	3.571	3.678	37	
36												3.913	4.080	4.125	4.152	4.024	3.986	3.930	3.497	3.936	3.846	3.606	3.623	36	

*degree from the east end of the equator

Table A2b. Peak Intensity ($\log E_0 - 19$) of Ca II H (erg/sec.com.ster.cm).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*
frame no.	71	70					0.496		1.320	0.553															71
	0.523	0.169		0.375	0.871		0.496		1.313	0.862															70
69	1.033	0.850		0.713	1.029	0.586	0.553		1.310	0.936												0.504			69
68	1.084	1.137		0.823	1.272	0.753			1.346	1.086		0.366										0.765			68
67	0.946	1.334	0.683	1.017	1.384	0.800	0.760		1.490	1.251		0.318										1.115			67
66	0.981	1.661	0.977	1.148	1.306	0.853	0.802		1.586	1.443		0.630										1.301			66
65	1.371	1.601	1.159	1.266	1.476	0.902	0.846	0.438	1.618	1.381		0.636										1.443			65
64	1.604	1.713	1.261	1.340	1.532	1.272	1.103	0.776	1.577	1.517		0.691										1.595			64
63	1.739	1.871	1.462	1.509	1.567	1.406	1.343	0.943	1.577	1.549	0.718	1.083										1.708			63
62	1.903	1.818	1.587	1.688	1.710	1.719	1.548	1.203	1.058	1.058	0.718	1.083													62
61	2.021	1.976	1.755	1.802	1.765	1.783	1.674	1.252	1.284	1.813	1.640	1.173													61
60	1.994	2.027	1.790	1.875	1.943	1.978	1.809	1.339	1.544	1.775	1.527	1.363	1.497								0.964	2.107			60
59	2.107	2.024	1.908	1.891	1.875	2.150	1.926	1.475	1.666	1.886	1.591	1.451	1.543			0.663					1.119	1.905			59
58	2.316	2.279	1.932	2.062	2.299	2.154	2.242	1.815	1.945	2.048	2.015	1.713	1.652			0.870			0.477		1.472	2.096			58
57	2.316	2.428	2.179	2.515	2.201	2.264	2.257	2.015	2.093	2.234	2.086	1.803	1.780			1.068	0.841		0.738		1.543	2.021	0.317		57
56	2.412	2.298	2.246	2.331	2.303	2.405	2.312	2.008	2.137	2.434	2.195	1.949	1.910	0.901		1.330	1.014		1.043		1.602	2.131	0.599		56
55	2.572	2.583	2.369	2.398	2.424	2.603	2.488	2.283	2.274	2.409	2.407	2.156	1.979	1.128	1.033	1.437	1.139		1.189		1.763	2.130	0.613		55
54	2.619	2.573	2.417	2.318	2.451	2.579	2.490	2.518	2.286	2.445	2.224	2.121	1.958	1.362	1.234	1.476	1.245		1.227		1.767	2.010	0.804		54
53	2.736	2.572	2.595	2.468	2.581	2.744	2.603	2.504	2.384	2.557	2.527	2.327	2.191	1.572	1.439	1.633	1.459		1.368		1.847	2.183	0.850		53
52	2.753	2.618	2.639	2.547	2.704	2.704	2.742	2.467	2.401	2.578	2.394	2.504	2.192	1.694	1.678	1.804	1.650	1.252	1.533		1.915	2.257	0.809		52
51	2.973	2.828	2.737	2.588	2.748	2.801	2.912	2.706	2.563	2.906	2.525	2.486	2.399	1.950	1.889	2.004	1.808	1.536	1.766		2.012	2.540	1.070		51
50	2.827	2.687	2.834	2.816	2.824	2.850	2.850	2.713	2.546	2.785	2.625	2.583	2.352	2.012	2.064	2.264	2.104	1.776	1.748		2.113	2.497	1.323		50
49	2.872	2.840	2.854	2.740	2.710	2.842	2.834	2.770	2.671	2.847	2.695	2.492	2.567	2.218	2.059	2.450	2.160	1.946	2.043	0.939	2.146	2.374	1.498		49
48	3.017	2.958	2.955	2.750	2.943	3.245	3.253	2.762	2.931	3.160	3.017	2.742	2.824	2.438	2.261	2.485	2.282	2.110	2.085	1.144	2.287	2.517	1.686		48
47	2.993	2.895	2.801	2.840	2.718	2.813	2.924	2.718	2.689	2.990	2.725	2.631	2.569	2.367	2.335	2.406	2.462	2.139	2.201	1.435	2.389	2.423	1.810	1.386	47
46	2.970	3.031	2.932	2.793	3.145	3.190	3.126	2.766	2.912	3.150	2.876	2.903	2.714	2.547	2.728	2.484	2.585	2.338	2.376	1.777	2.574	2.478	1.898	1.665	46
45	2.994	3.033	2.966	2.753	3.191	2.977	3.105	2.847	3.043	3.028	2.898	2.724	2.694	2.582	2.720	2.538	2.581	2.575	2.369	1.960	2.605	2.540	2.085	2.017	45
44	2.918	3.048	2.770	2.954	3.243	3.181	3.172	3.125	2.946	3.111	3.170	2.999	2.888	2.641	2.720	2.618	2.786	2.477	2.682	1.877	2.684	2.744	2.107	2.155	44
43	2.849	2.888	2.964	2.912	3.096	3.234	3.153	2.855	2.961	3.038	3.178	3.098	2.773	2.568	2.598	2.631	2.772	2.642	2.494	2.120	2.626	2.603	2.388	2.222	43
42	2.387			2.981	2.928	3.227	3.126	2.994	2.860	3.214	2.909	3.153	2.969	2.832	2.902	2.959	3.088	3.019	2.658	2.112	2.786	2.528	2.381	2.327	42
41	2.843			3.054	3.222	3.193	3.187	2.896	2.854	3.093	3.036	2.978	2.800	2.775	2.804	2.955	2.954	2.782	2.708	2.317	2.523	2.572	2.366	2.444	41
40	2.904			2.938	3.074	3.035	3.167	2.935	3.165	3.182	3.104	3.189	2.981	2.787	3.087	3.016	3.036	3.043	2.834	2.392	2.796	2.723	2.538	2.491	40
39								3.121	2.903	3.274	3.251	3.213	3.134	2.983	3.224	3.160	3.177	2.990	2.877	2.690	2.947	3.126	2.619	2.662	39
38												3.159	2.864	2.787	3.108	3.016	3.036	2.998	2.973	2.500	3.055	2.728	2.528	2.631	38
37												3.122	3.022	3.152	3.156	3.033	3.120	2.688	2.808	2.715	3.024	2.800	2.677	2.773	37
36													2.870	3.270	3.052	3.120	3.078	3.152	3.086	2.567	3.110	3.030	2.788	2.835	36

*degree from the east end of the equator

Table A2c. Total $1/e$ -Width of Ca II H (km/sec).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.	
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*	
frame no.	71	70					31.1		53.8	34.3															71	
				26.8	27.9				45.4	26.8															70	
69	33.4	21.0		24.8	32.1	46.6			51.0	31.8											23.2				69	
68	23.8	24.0		30.0	25.4	32.4			50.7	34.5											28.1				68	
67	44.1	24.8	23.3	22.3	23.5	39.0	25.7		48.9	37.5		58.2									28.0				67	
66	56.0	24.8	22.3	20.8	27.1	51.7	40.5		48.9	37.1		55.4									21.8				66	
65	40.1	25.9	21.3	25.8	22.2	59.3	31.6	60.1	42.8	33.7		55.1									26.6				65	
64	42.7	26.6	26.2	28.1	22.5	52.2	34.4	51.4	39.7	36.9		42.3									21.6				64	
63	44.6	27.1	29.9	29.1	28.2	62.9	31.5		46.5	34.0	43.0	42.3									22.4				63	
62	40.1	33.7	27.7	31.2	28.3	46.2	32.8	39.6	46.7	38.6	39.4	49.6	45.6												62	
61	40.6	29.9	30.9	27.0	31.2	56.5	31.5	35.4	54.8	35.6	36.7														61	
60	52.5	25.2	29.0	29.5	30.4	52.5	39.2	44.7	45.1	38.4	64.1	47.2								60.9	19.6				60	
59	48.5	51.0	40.1	28.6	44.7	45.3	41.7	52.3	41.9	43.3	70.8	40.4	47.3			31.1					69.0	26.5			59	
58	48.4	34.3	38.9	28.0	36.9	54.8	34.4	51.5	37.5	52.4	63.2	43.4	54.0			31.0			40.8		56.9	25.7			58	
57	52.0	40.3	36.5	15.7	43.8	56.6	45.8	49.3	39.6	48.4	70.1	45.4	47.5			26.9	44.2			59.0	58.1	26.0	35.1		57	
56	50.4	54.7	45.7	31.9	52.3	64.1	53.1	62.5	49.0	35.0	72.1	49.9	45.4	38.6		26.7	42.2			36.3	55.1	24.6	45.7		56	
55	51.3	45.0	42.3	35.5	50.0	46.5	53.6	50.5	40.4	50.7	61.0	50.5	53.0	37.2	41.8	26.2	46.3			44.0	44.2	27.8	35.2		55	
54	55.1	50.2	47.8	45.9	47.8	62.0	55.6	42.6	52.5	42.2	74.9	51.8	52.0	36.8	40.8	29.2	51.0			51.9	45.1	31.7	35.3		54	
53	49.8	55.8	36.6	45.3	46.7	53.1	47.8	47.3	50.8	48.8	59.7	50.0	41.4	39.5	41.9	25.9	53.9			45.4	48.6	29.2	35.4		53	
52	57.4	50.6	41.6	39.7	46.1	58.6	52.1	54.1	51.3	56.6	70.6	46.2	50.2	41.1	37.6	32.6	46.6	28.6		40.4	44.6	31.2	35.8		52	
51	53.4	49.0	47.6	50.1	51.3	44.7	44.7	43.0	40.0	51.0	70.6	59.0	55.1	44.8	41.0	34.7	53.2	30.6		35.5	39.8	24.8	38.0		51	
50	63.2	64.0	36.4	41.0	50.8	70.5	56.8	56.9	53.6	56.5	71.5	48.1	53.0	38.0	35.4	29.3	47.0	40.2		53.3	39.9	24.1	36.8		50	
49	43.6	57.2	46.3	61.8	65.5	66.5	49.0	48.8	55.6	53.1	65.6	64.3	38.6	33.2	42.7	27.8	56.1	40.5	38.2	35.6	49.2	30.4	33.1		49	
48	55.0	56.7	46.3	55.3	52.0	62.2	52.5	59.1	51.4	57.0	63.8	54.0	30.4	35.4	42.0	38.5	63.3	34.8	41.6	38.9	45.4	28.8	39.0		48	
47	59.8	62.8	54.3	46.8	67.7	64.4	61.9	62.5	62.9	57.7	69.9	66.5	44.6	33.5	49.8	46.7	56.3	45.8	43.0	42.6	39.1	33.0	32.6	33.6		47
46	62.4	66.2	56.6	56.5	31.7	55.3	63.4	62.4	33.2	58.8	66.1	47.7	49.3	38.6	36.0	53.6	63.0	42.4	46.0	39.2	27.9	35.2	38.2	31.2		46
45	64.3	67.7	56.0	56.8	47.0	61.4	69.6	67.2	57.9	63.2	68.3	70.5	54.6	43.8	43.7	48.1	65.3	36.0	49.2	39.4	31.5	40.2	42.2	36.7		45
44	71.6	53.5	62.1	56.7	40.6	64.8	57.2	58.6	62.6	64.2	62.7	68.5	49.3	42.5	54.6	55.1	54.0	51.0	36.6	46.7	40.5	29.5	46.5	41.0		44
43	69.1	69.2		51.1	53.0	49.5	60.1	56.7	61.5	66.9	68.3	56.1	49.9	49.0	63.9	55.5	61.0	50.3	52.9	46.2	41.5	37.9	34.7	48.4		43
42	64.1			61.7	60.1	60.7	66.3	62.4	66.6	65.5	80.0	64.5	48.2	51.3	55.6	53.5	49.4	38.7	43.6	48.9	42.8	51.6	45.6	51.3		42
41	68.0			62.3	72.7	77.1	69.6	69.4	66.1	72.9	73.0	66.6	55.6	60.6	57.3	68.0	56.3	45.9	48.8	52.5	54.8	38.0	48.8	46.2		41
40	62.7			72.5	49.2	65.1	78.3	76.2	52.9	65.3	73.1	66.1	54.5	66.7	60.0	56.5	54.5	44.6	56.9	51.1	54.2	41.9	48.2	48.4		40
39								52.0	71.4	75.9	62.7	62.8	62.1	65.2	65.1	55.4	55.6	51.4	59.0	40.6	45.1	33.4	50.0	49.6		39
38												69.9	61.3	65.4	61.7	71.0	69.7	60.6	49.4	60.0	44.6	56.2	57.3	57.0		38
37												69.7	61.6	58.5	63.0	66.6	72.1	75.1	70.0	48.4	49.2	53.0	59.0	52.0		37
36													75.5	53.8	76.5	73.5	75.8	77.8	49.6	61.2	46.5	49.8	48.5	44.8		36

*degree from the east end of the equator

Table A3a. Total Intensity ($\log E - 10$) of CaII K (erg/sec.cm.ster).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.		
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*		
frame no.	71	70																							71		
	0.959	1.247		1.128	1.679	0.985	1.287			2.314	1.543														70		
				1.355	1.735		1.527			2.301	1.684																
69	1.596	1.569		1.455	1.820	1.560	1.509			2.294	1.746															69	
68	1.554	1.805		1.456	1.917	1.773	1.642			2.488	2.018											1.381				68	
67	1.903	2.099	1.439	1.820	2.182	1.868	1.732			2.472	2.103	1.411										1.576				67	
66	2.087	2.242	1.639	1.863	2.162	1.894	1.761			2.541	2.255	1.635										1.780				66	
65	2.209	2.333	1.912	1.994	2.138	1.973	1.810	1.713		2.562	2.294	1.644										1.905				65	
64	2.417	2.531	2.032	2.143	2.358	2.358	1.935	1.871		2.559	2.402	1.748										2.119				64	
63	2.735	2.685	2.337	2.527	2.540	2.567	2.226	2.082		2.690	2.407	2.066										2.405				63	
62	2.819	2.754	2.314	2.442	2.587	2.677	2.487	2.152	2.094	2.732	2.495	2.156										2.249				62	
61	2.953	2.803	2.636	2.612	2.739	2.856	2.570	2.257	2.341	2.751	2.537	2.168	2.336									2.554				61	
60	3.084	2.938	2.634	2.582	2.773	3.069	2.651	2.346	2.530	2.727	2.654	2.285	2.479								2.061	2.598				60	
59	3.078	2.951	2.895	2.665	2.898	3.101	2.925	2.599	2.709	2.792	2.765	2.442			1.425							2.291	2.691			59	
58	3.436	3.348	2.919	3.060	3.091	3.223	3.209	2.864	2.940	3.081	3.056	2.731	2.738		1.659				1.309			2.505	2.763			58	
57	3.420	3.292	3.113	2.905	3.208	3.329	3.177	2.946	3.048	3.143	3.097	2.813	2.806		1.866	1.757			1.807			2.630	2.833	1.381		57	
56	3.381	3.302	3.189	3.134	3.251	3.521	3.392	3.242	3.192	3.335	3.220	2.977	2.928	1.761	2.100	1.996			2.033			2.685	2.903	1.502		56	
55	3.502	3.452	3.338	3.225	3.311	3.603	3.412	3.244	3.159	3.391	3.424	3.146	2.995	2.004	1.806	2.178	2.230		2.200			2.836	3.033	1.627		55	
54	3.586	3.576	3.339	3.332	3.459	3.737	3.630	3.444	3.255	3.463	3.508	3.214	3.118	2.265	2.152	2.324	2.383		2.267			2.750	2.920	1.701		54	
53	3.580	3.605	3.460	3.418	3.505	3.655	3.666	3.396	3.318	3.494	3.505	3.287	3.099	2.444	2.397	2.402	2.583		2.436			2.857	3.026	1.808		53	
52	3.778	3.603	3.446	3.469	3.454	3.701	3.760	3.557	3.492	3.613	3.563	3.433	3.204	2.590	2.571	2.614	2.747	2.157	2.491			2.924	3.029	1.739		52	
51	3.726	3.763	3.510	3.577	3.783	3.757	3.847	3.682	3.599	3.760	3.833	3.594	3.513	2.853	2.766	2.897	2.963	2.466	2.734			3.049	3.236	2.070		51	
50	3.840	3.848	3.556	3.592	3.778	3.897	3.822	3.732	3.633	3.854	3.739	3.610	3.629	2.962	2.879	3.015	3.147	2.649	2.826			3.200	3.407	2.240		50	
49	3.845	3.883	3.671	3.691	3.886	3.900	3.817	3.846	3.812	3.876	3.826	3.672	3.700	3.149	3.033	3.331	3.215	2.815	2.886	1.852		3.233	3.178	2.473		49	
48	3.959	3.853	3.852	3.864	3.901	4.143	4.155	4.009	3.833	4.124	3.974	3.856	3.591	3.292	3.352	3.357	3.578	2.916	3.129	2.161		3.248	3.508	2.564		48	
47	3.953	3.806	3.756	3.667	3.814	4.018	3.890	3.903	3.813	3.963	3.753	3.743	3.612	3.460	3.284	3.555	3.453	3.168	3.057	3.168	2.402	3.418	3.388	2.812	2.429	47	
46	4.045	3.843	3.840	3.902	3.996	4.226	4.087	3.913	3.963	4.107	4.078	4.082	3.690	3.531	3.496	3.768	3.703	3.366	3.374	2.718		3.393	3.477	2.970	2.640	46	
45	4.088	3.922	3.820	3.828	3.938	4.110	4.141	3.882	4.097	4.145	3.996	3.960	3.883	3.607	3.470	3.603	3.886	3.446	3.265	2.866		3.477	3.486	3.012	2.865	45	
44	4.184	4.035	4.004	3.942	4.015	4.166	4.219	3.975	4.150	4.122	3.988	3.805	3.846	3.529	3.606	3.738	3.727	3.462	3.554	2.985		3.527	3.473	3.112	3.114	44	
43	4.012	3.929	3.955	3.970	3.911	4.039	4.097	4.078	4.086	4.125	4.162	3.994	3.851	3.660	3.843	3.655	3.925	3.757	3.512	3.082		3.593	3.671	3.267	3.088	43	
42	4.083			3.970	3.973	4.140	4.095	3.980	4.071	4.091	3.905	3.936	3.823	3.803	3.910	3.924	4.000	3.832	3.677	3.194		3.785	3.702	3.365	3.259	42	
41	3.875			3.997	4.058	4.206	4.153	4.213	4.190	4.173	4.186	4.003	3.970	3.943	3.933	3.931	4.164	3.748	3.796	3.304		3.758	3.726	3.348	3.422	41	
40	3.796			4.131	3.988	3.959	4.305	4.138	4.071	4.304	4.189	4.079	3.960	3.815	3.993	3.877	4.105	3.758	3.885	3.338		3.760	3.748	3.618	3.564	40	
39							4.232	4.082	4.141	4.242	4.202	4.175	3.982	4.033	4.098	4.108	3.897	3.959	3.400			3.813	3.748	3.601	3.777	39	
38							4.131				4.131	4.062	4.030	4.185	4.006	4.095	3.927	3.986	3.640			3.996	3.864	3.659	3.676	38	
37							4.079				4.079	3.996	4.026	4.085	4.055	4.213	4.061	3.889	3.624			4.064	3.976	3.693	3.810	37	
36												4.149	4.100	4.222	4.122	4.246	4.246	3.939	4.051			3.784	4.011	3.984	3.675	3.799	36

*degree from the east end of the equator

Table A3b. Peak Intensity ($\log E_0 - 19$) of Ca II K (erg/sec.cm.ster.cm).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	27.5	traced position no.
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*	
frame no.																									frame no.	
71																									71	
70	0.555	0.899		0.719	1.231	0.331	0.672			1.602	0.882														70	
				0.905	1.174		0.896			1.584	1.051															
69	0.959	1.201		1.010	1.410	0.839	0.947			1.534	1.111															69
68	1.000	1.393		1.164	1.511	0.991	0.999			1.715	1.393	0.556									0.795					68
67	1.125	1.584	1.106	1.514	1.660	1.083	1.065			1.717	1.454	0.598									1.178					67
66	1.354	1.712	1.216	1.512	1.674	1.214	1.158			1.823	1.625	0.882									1.306					66
65	1.441	1.826	1.523	1.727	1.680	1.176	1.165	0.804		1.818	1.620	0.897									1.381					65
64	1.571	1.971	1.570	1.769	1.854	1.501	1.283	1.137		1.872	1.689	1.012									1.744					64
63	1.935	1.994	1.744	1.866	2.005	1.718	1.636	1.323		1.939	1.688	0.885	1.388								2.009					63
62	2.042	2.089	1.801	1.988	1.980	1.929	1.901	1.526	1.412	2.083	1.756	1.400	1.400								1.911					62
61	2.255	2.170	2.030	2.111	2.081	2.001	1.895	1.644	1.623	1.960	1.769	1.436	1.634								2.199					61
60	2.273	2.328	1.986	2.070	2.048	2.287	1.944	1.616	1.816	2.036	1.796	1.621	1.772								1.192	2.172				60
59	2.348	2.207	2.265	1.973	2.162	2.292	2.207	1.813	1.924	1.990	1.788	1.703	1.707			0.880					1.445	2.168				59
58	2.607	2.582	2.235	2.581	2.326	2.317	2.473	2.023	2.180	2.318	2.086	1.989	1.944			1.157			0.802		1.657	2.268				58
57	2.537	2.555	2.436	2.162	2.518	2.398	2.463	1.967	2.301	2.313	2.151	1.988	2.015			1.362	1.035				1.812	2.350	0.726			57
56	2.516	2.530	2.447	2.385	2.554	2.659	2.515	2.322	2.400	2.468	2.247	2.188	2.158	1.119		1.602	1.174				1.948	2.409	0.804			56
55	2.663	2.646	2.566	2.516	2.566	2.956	2.573	2.353	2.441	2.490	2.515	2.332	2.345	1.392	1.198	1.673	1.458				2.119	2.490	0.995			55
54	2.649	2.789	2.593	2.548	2.742	2.916	2.660	2.584	2.648	2.577	2.605	2.415	2.331	1.586	1.454	1.778	1.638				2.134	2.312	1.075			54
53	2.806	2.784	2.688	2.695	2.654	2.770	2.955	2.563	2.492	2.615	2.523	2.399	2.445	1.802	1.694	1.819	1.830				2.072	2.424	1.207			53
52	2.899	2.805	2.685	2.713	2.632	2.856	2.948	2.629	2.706	2.658	2.545	2.479	2.417	1.919	1.899	2.025	1.977	1.556	1.749		2.142	2.409	1.127			52
51	2.805	2.902	2.784	2.804	2.952	2.862	2.899	2.788	2.720	2.938	2.856	2.670	2.800	2.135	2.051	2.363	2.105	1.797	2.011		2.329	2.715	1.468			51
50	3.026	3.046	2.780	2.749	3.044	2.988	3.038	2.840	2.843	3.013	2.724	2.832	2.794	2.291	2.154	2.319	2.284	2.001	2.080		2.523	2.898	1.628			50
49	2.902	3.016	2.869	2.827	3.019	2.959	2.896	3.037	2.929	2.883	2.826	2.848	2.881	2.389	2.326	2.571	2.241	2.069	2.068	1.098	2.500	2.563	1.752			49
48	3.083	3.044	3.059	3.046	3.037	3.286	3.287	3.085	2.930	3.322	2.971	3.023	2.781	2.586	2.590	2.716	2.600	2.170	2.380	1.352	2.540	2.934	1.911			48
47	3.085	2.963	2.970	2.811	2.974	3.141	3.019	3.043	2.911	2.965	2.673	2.813	2.925	2.723	2.515	2.831	2.517	2.359	2.483	1.661	2.833	2.695	1.983	1.741		47
46	3.183	3.020	2.998	3.096	3.149	3.300	3.261	2.963	3.143	3.221	3.158	3.192	2.799	2.787	2.642	3.090	2.869	2.685	2.544	2.001	2.594	2.839	2.276	1.958		46
45	3.125	3.008	2.913	3.035	3.074	3.169	3.233	3.031	3.177	3.281	3.056	2.927	2.934	2.793	2.538	2.772	2.875	2.658	2.470	2.110	2.736	2.845	2.242	2.135		45
44	3.274	2.991	3.183	3.120	3.106	3.250	3.285	3.031	3.208	3.211	2.953	2.804	3.000	2.630	2.724	3.025	2.757	2.595	2.668	2.285	2.837	2.733	2.329	2.376		44
43	3.182	3.110	3.250	3.110	2.979	3.214	3.143	3.207	3.183	3.207	3.214	3.032	3.065	2.778	2.942	2.695	2.972	2.894	2.628	2.319	2.931	3.031	2.480	2.326		43
42	3.076			2.997	3.124	3.243	3.182	3.067	3.196	3.044	2.883	2.977	2.881	2.812	2.937	3.013	3.003	2.931	2.881	2.412	3.107	3.086	2.532	2.435		42
41	3.080			3.005	3.051	3.241	3.165	3.307	3.173	3.205	3.191	3.006	3.140	2.960	2.931	3.011	3.306	2.822	2.966	2.493	3.022	2.928	2.495	2.654		41
40	2.755			3.148	3.155	3.020	3.296	3.152	3.050	3.266	3.205	3.037	3.070	2.899	3.076	3.085	2.766	2.622	2.895	3.071	2.895	3.071	2.890	2.672		40
39								3.198	3.154	3.105	3.217	3.283	3.244	2.936	3.095	3.205	3.252	2.959	3.038	2.486	2.994	2.867	2.753	2.948		39
38												3.149	3.044	3.032	3.217	3.065	2.945	2.937	3.141	2.797	3.110	2.892	2.792	2.720		38
37												3.018	2.964	3.012	3.016	3.084	3.236	3.254	2.915	2.685	3.308	3.198	2.768	2.821		37
36													3.193	3.170	3.283	3.132	3.201	2.886	3.224	2.775	3.107	2.867	2.856	2.910		36

*degree from the east end of the equator

Table A4a. Total Intensity ($\log E - 10$) of He ϵ (erg/sec.cm.ster).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.	
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*	
frame no.	63																								frame no.	
	0.879					0.905			1.019																63	
	1.079					0.731			1.078																62	
	1.109					1.193	1.063																		61	
	1.417	1.062	0.518																						60	
	1.682	1.512	0.974	1.039	1.108	1.499	1.425			1.245	0.863	0.763									1.049				59	
	1.850	1.539	1.174	1.135	1.397	1.644	1.561	1.279	1.366	1.496	1.376	0.866														58
	1.984	1.773	1.596	1.267	1.624	1.790	1.743	1.603	1.470	1.615	1.604	1.071	1.011									0.990				57
	2.208	2.062	1.746	1.402	1.731	1.925	1.895	1.621	1.463	1.756	1.679	1.322										0.828				56
	2.382	2.190	1.966	1.658	1.883	2.001	2.032	1.939	1.580	1.954	1.956	1.517										1.106				55
	2.508	2.350	2.027	1.804	2.078	2.252	2.168	1.941	1.860	2.112	2.037	1.663	1.137													54
	2.694	2.512	2.316	1.962	2.314	2.401	2.366	2.142	1.994	2.250	2.217	1.760	1.439													53
	2.823	2.690	2.530	2.082	2.323	2.445	2.469	2.175	2.172	2.383	2.350	1.950	1.495													52
	2.968	2.890	2.660	2.217	2.588	2.655	2.667	2.368	2.289	2.531	2.443	2.112	1.627									1.027				51
	3.131	3.008	2.859	2.386	2.694	2.818	2.809	2.520	2.426	2.714	2.658	2.286	1.793	1.347	1.537	1.268	1.528					1.431				50
	3.296	3.057	2.954	2.620	2.784	2.955	2.992	2.646	2.524	2.846	2.757	2.390	1.880	1.543	1.684	1.565	1.644					1.265				49
	3.478	3.446	3.141	2.763	3.054	3.096	3.181	2.848	2.853	2.983	3.016	2.516	1.973	1.706	1.815	1.858	1.890	1.343			1.599	1.738				48
	3.482	3.368	3.289	2.893	3.074	3.164	3.134	2.972	3.019	3.227	3.083	2.709	2.118	1.890	2.097	1.970	2.133	1.651			1.592	1.630				47
	3.506	3.551	3.444	3.018	3.234	3.350	3.400	3.033	3.166	3.243	3.293	2.975	2.261	2.167	2.298	2.308	2.408	1.972	1.632			1.553	1.480			46
	3.721	3.535	3.571	3.219	3.405	3.439	3.374	3.192	3.199	3.411	3.354	3.055	2.363	2.341	2.550	2.405	2.650	2.199	1.684	1.446	1.767	1.676	1.118	1.299		45
	3.851	3.703	3.644	3.256	3.497	3.783	3.558	3.503	3.251	3.487	3.662	3.294	2.560	2.522	2.647	2.582	2.763	2.405	1.921	1.421	1.937	1.733	1.559	1.672		44
	3.632	3.556	3.535	3.400	3.471	3.681	3.718	3.511	3.626	3.494	3.508	3.398	2.721	2.652	2.813	2.748	2.938	2.590	2.028	1.618	2.158	1.945	1.802	2.041		43
	3.818			3.504	3.616	3.667	3.695	3.540	3.637	3.558	3.674	3.554	2.847	2.934	3.054	2.917	3.136	2.805	2.234	1.715	2.168	1.936	1.863	1.956		42
				3.605	3.670	3.754	3.809	3.558	3.688	3.712	3.546	3.578	2.998	3.030	3.183	3.094	3.235	3.071	2.482	2.046	2.330	2.051	1.880	2.131		41
				3.770	3.570	3.600	3.854	3.776	3.633	3.748	3.740	3.721	3.136	3.094	3.338	3.264	3.472	3.192	2.671	2.268	2.521	2.137	2.100	2.293		40
								3.768	3.980	4.081	3.964	3.780	3.358	3.378	3.648	3.499	3.666	3.389	2.931	2.576	2.748	2.442	2.343	2.533		39
												3.864	3.313	3.445	3.678	3.572	3.557	3.293	3.025	2.649	2.631	2.662	2.253	2.641		38
												3.777	3.505	3.607	3.762	3.739	3.800	3.491	3.204	2.935	3.012	2.747	2.608	2.867		37
													3.498	3.653	3.843	3.774	3.774	3.807	3.271	3.150	3.292	2.783	2.652	2.935		36

*degree from the east end of the equator

Table A4b. Peak Intensity ($\log E_0 - 19$) of He (erg/sec.cm.ster.cm).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*
frame no.	63																								frame no.
	0.185					0.131				0.322															63
	0.288					0.052				0.276															62
	0.487					0.414	0.537																		61
	0.649	0.385	0.048																						60
	0.862	0.675	0.326	0.310	0.434	0.677	0.651		0.338	0.513	-0.14	0.083	0.138								0.357				59
	1.032	0.782	0.592	0.486	0.584	0.842	0.827	0.477	0.610	0.803	0.526	0.210	0.233								0.530				58
	1.239	1.048	0.893	0.596	0.922	0.945	1.092	0.807	0.682	0.946	0.814	0.247	0.400								0.425				57
	1.497	1.297	1.130	0.808	1.034	1.150	1.172	0.928	0.807	1.060	0.874	0.579									0.334				56
	1.701	1.547	1.372	0.990	1.152	1.287	1.357	1.140	0.941	1.134	1.119	0.736	0.335								0.331				55
	1.812	1.677	1.486	1.142	1.379	1.507	1.411	1.216	1.218	1.454	1.242	0.925	0.529												54
	2.025	1.808	1.706	1.291	1.637	1.616	1.669	1.356	1.314	1.551	1.405	0.977	0.710												53
	2.104	1.979	1.906	1.351	1.624	1.716	1.780	1.430	1.500	1.683	1.591	1.218	0.739												52
	2.262	2.130	1.978	1.562	1.913	1.892	1.923	1.617	1.546	1.901	1.749	1.391	0.921				0.349				0.698				51
	2.449	2.315	2.231	1.772	1.981	2.105	2.081	1.825	1.746	1.965	1.900	1.548	1.139	0.641	0.796	0.766	0.681				0.823				50
	2.628	2.336	2.375	1.974	2.075	2.190	2.351	1.893	1.802	2.178	1.939	1.661	1.187	0.818	1.037	0.922	0.795				0.626				49
	2.723	2.763	2.466	2.144	2.368	2.388	2.563	2.097	2.137	2.371	2.421	1.733	1.297	0.998	1.178	1.105	1.119	0.616			0.824	0.805			48
	2.757	2.671	2.565	2.319	2.419	2.474	2.300	2.207	2.316	2.541	2.313	2.005	1.399	1.191	1.453	1.413	1.398	1.036	0.652		0.818	0.878			47
	2.816	2.912	2.801	2.322	2.640	2.735	2.699	2.330	2.496	2.495	2.545	2.264	1.513	1.473	1.666	1.605	1.582	1.335	0.820	0.469	0.794	0.935			46
	2.971	2.764	2.790	2.510	2.805	2.761	2.569	2.413	2.375	2.673	2.627	2.323	1.660	1.677	1.816	1.676	1.885	1.491	0.909	0.743	1.017	1.043	0.606	0.712	45
	3.198	2.936	2.933	2.679	2.723	3.101	2.871	2.745	2.432	2.800	2.989	2.557	1.834	1.784	1.865	1.922	1.969	1.660	1.174	0.762	1.232	1.103	0.787	0.928	44
	2.931	2.720	2.873	2.693	2.809	3.093	2.999	2.657	2.934	2.722	2.781	2.673	2.050	1.957	1.985	2.047	2.123	1.902	1.291	0.964	1.383	1.115	1.045	1.243	43
	3.207												2.095	2.216	2.297	2.129	2.445	2.036	1.488	0.977	1.469	1.191	1.091	1.318	42
	2.812												2.377	2.287	2.451	2.273	2.425	2.395	1.734	1.329	1.604	1.315	1.165	1.475	41
	2.918												2.419	2.361	2.586	2.421	2.694	2.458	1.944	1.511	1.796	1.460	1.352	1.673	40
								2.723	3.271	3.320	3.175	2.909	2.657	2.613	2.818	2.727	2.830	2.757	2.183	1.860	1.970	1.638	1.610	1.814	39
												2.995	2.513	2.590	2.881	2.737	2.793	2.554	2.344	1.874	1.888	1.946	1.583	1.989	38
												3.053	2.761	2.822	3.009	2.847	2.917	2.633	2.427	2.192	2.288	1.995	1.867	2.092	37
													2.740	2.854	3.159	2.988	2.975	3.231	2.567	2.461	2.576	1.988	1.972	2.165	36

*degree from the east end of the equator

Table A4c. Total $1/e$ -Width of $H\epsilon$ (km/sec).

traced position no.	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	traced position no.
position angle*	48.7	48.0	47.5	45.0	44.7	44.0	43.6	42.1	42.7	41.5	40.9	40.0	33.4	32.8	31.9	31.6	31.3	31.0	29.5	30.3	28.9	27.9	27.3	27.5	position angle*
frame no.																									frame no.
63																									63
62																									62
61	39.4									44.1															61
60	41.5	38.4					27.6			59.8															60
59	50.1	51.8	37.0	42.5	45.0	63.0	45.0			43.9											43.3				59
58	51.7	49.2	33.4	34.3	54.2	48.2	45.4	52.4	43.5	42.0	62.5														58
57	44.0	41.4	40.1	40.7	40.2	56.4	38.1	51.7	47.7	42.2	60.2	35.2									28.0				57
56	44.7	48.6	35.9	35.0	38.4	49.8	42.1	43.7	37.5	43.0	67.0														56
55	39.3	38.3	36.7	40.7	48.2	47.4	40.9	52.4	33.4	58.0	59.7	51.1													55
54	39.1	41.6	34.6	43.0	36.2	43.6	47.2	45.9	40.5	36.8	54.5	46.0													54
53	35.2	40.5	32.8	38.8	36.0	50.3	40.2	53.4	41.8	39.1	52.1	43.5													53
52	42.6	42.0	33.5	44.8	40.1	41.0	41.0	43.2	37.1	47.4	50.3	46.1													52
51	38.9	48.1	35.6	36.7	37.9	46.7	45.2	45.6	45.1	31.8	42.4	45.0									36.4				51
50	38.8	43.0	37.0	33.1	38.8	42.6	43.2	42.8	39.4	42.1	42.0	44.8	41.8	40.4	45.2	26.2	54.4				32.6				50
49	35.4	45.8	37.0	34.1	41.6	45.2	32.8	46.9	42.0		50.0	39.1	43.7	41.8	40.5	36.9	57.0				36.3				49
48	42.5	36.7	36.3	34.8	37.8	39.8	38.2	41.0	39.0	32.4	31.4	56.9	42.4	44.8	36.6	47.6	44.6				48.2	64.0			48
47	43.2	41.6	41.3	29.7	37.2	40.4	51.9	43.6	35.1	37.2	43.9	46.7	52.4	41.6	33.4	32.2	41.5	33.7			53.4	47.8			47
46	38.4	36.6	35.8	40.0	26.5	33.0	34.4	35.8	34.3	43.2	40.8	46.7	41.9	41.2	34.9	39.5	53.5	39.7			46.5	29.5			46
45	40.5	46.6	41.7	40.7	33.0	37.9	47.4	45.8	51.7	43.4	43.3	46.2	40.0	36.3	45.2	41.8	46.5	40.8	51.9	42.2	51.4	38.3	29.2	30.3	45
44	30.3	43.6	34.5	30.2	45.7	41.3	38.9	40.8	48.7	37.4	33.8	42.8	42.4	38.0	47.2	41.6	52.4	43.1	46.3	40.6	41.7	36.0	51.6	46.5	44
43	42.0	41.8	43.0	43.4	33.9	31.2	36.4	42.7	36.2	48.1	36.1	37.2	38.3	41.8	49.8	44.2	48.4	42.5	44.6	31.8	50.2	60.4	28.8	49.0	43
42	28.9																								42
41	35.4																								41
40	57.5																								40
39																									39
38																									38
37																									37
36																									36

*degree from the east end of the equator

Table A5. Mean Emission Gradients of Balmer Lines.

height (km)	emission gradient $\times 10^{-8}$ (cm ⁻¹)
500	1.589
1000	1.934
1500	1.900
2000	1.785
2500	1.600
3000	1.255